

# Chiral U(1) flavor models and flavored Higgs doublets for the top $A_{FB}$ and the $W_{jj}$

Chaehyun Yu (KIAS)

Based on arXiv:1108.0350 [hep-ph]; 1108.4005 [hep-ph]  
with P. Ko and Yuji Omura (KIAS)

The 1<sup>st</sup> KIAS Phenomenology workshop,  
KIAS, Seoul, Korea, Nov 17-19, 2011

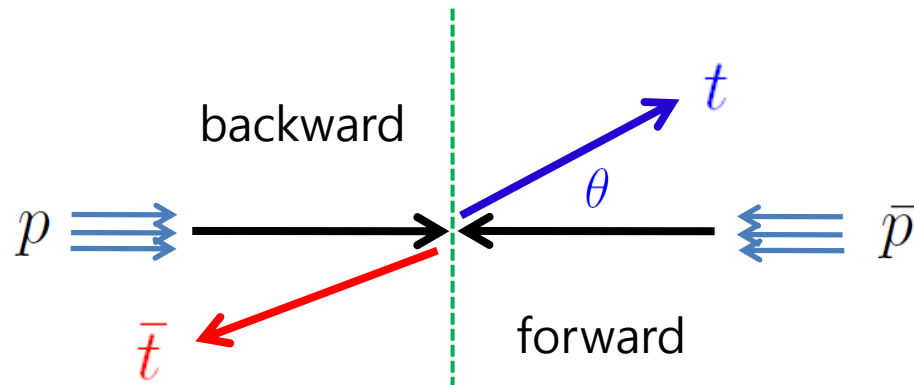
# Contents

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- Top forward-backward asymmetry at Tevatron
- $W + jj$  excess at Tevatron
- Chiral flavor  $U(1)$  model
- Phenomenology
- Conclusions

# Top $A_{FB}$ at Tevatron

- $A_{FB}$  in top pair production : one of the most interesting quantities.
  - charge asymmetry when CP conservation is assumed.



$$A_{FB} = \frac{N(t; \cos \theta > 0) - N(t; \cos \theta < 0)}{N(t; \cos \theta > 0) + N(t; \cos \theta < 0)}$$

- Symmetric under charge conjugation at LO in the SM.
- At NLO, the interference of processes that differ under charge conjugation leads a small forward-backward asymmetry.

$$A_{SM}^{t\bar{t}} = 0.058 \pm 0.009 \quad (\text{MCFM})$$

# Top $A_{\text{FB}}$ at CDF

CDF, 1101.0034

- CDF(2011) : lepton+jets channel ( $5.3 \text{ fb}^{-1}$ )

$$A^{\text{P}\bar{\text{P}}} = 0.150 \pm 0.055 \text{ (stat+sys)} \quad \sim 2\sigma$$

$$A^{\text{t}\bar{\text{t}}} = 0.158 \pm 0.075 \text{ (stat+sys)} \quad \sim 1.32\sigma$$

- $\Delta y = y_t - y_{\bar{t}}$

	CDF	SM	
$A^{\text{t}\bar{\text{t}}}( \Delta y  < 1.0)$	$0.026 \pm 0.118$	$0.039 \pm 0.006$	
$A^{\text{t}\bar{\text{t}}}( \Delta y  > 1.0)$	$0.611 \pm 0.256$	$0.123 \pm 0.008$	$\sim 2\sigma$
$A^{\text{t}\bar{\text{t}}}(m_{\text{t}\bar{\text{t}}} < 450\text{GeV})$	$-0.116 \pm 0.153$	$0.040 \pm 0.006$	
$A^{\text{t}\bar{\text{t}}}(m_{\text{t}\bar{\text{t}}} > 450\text{GeV})$	$0.475 \pm 0.114$	$0.088 \pm 0.013$	$\sim 3.4\sigma$

- $m_{\text{t}\bar{\text{t}}}$

SM

(See the talk by H.Kim and S.Choi.)

# Top $A_{FB}$ at CDF

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- CDF (2011) : dilepton channel

✓ Observed (including background)

$$A_{FB}^{obs} = 0.138 \pm 0.054 (stat.)$$

From Hyunsoo Kim's talk, Tev2011

✓ Observed (background subtracted)

$$A_{FB}^{sub} = 0.205 \pm 0.073 (stat.) \pm 0.021 (bkg\ shape)$$

✓ Detector effect corrected (true)

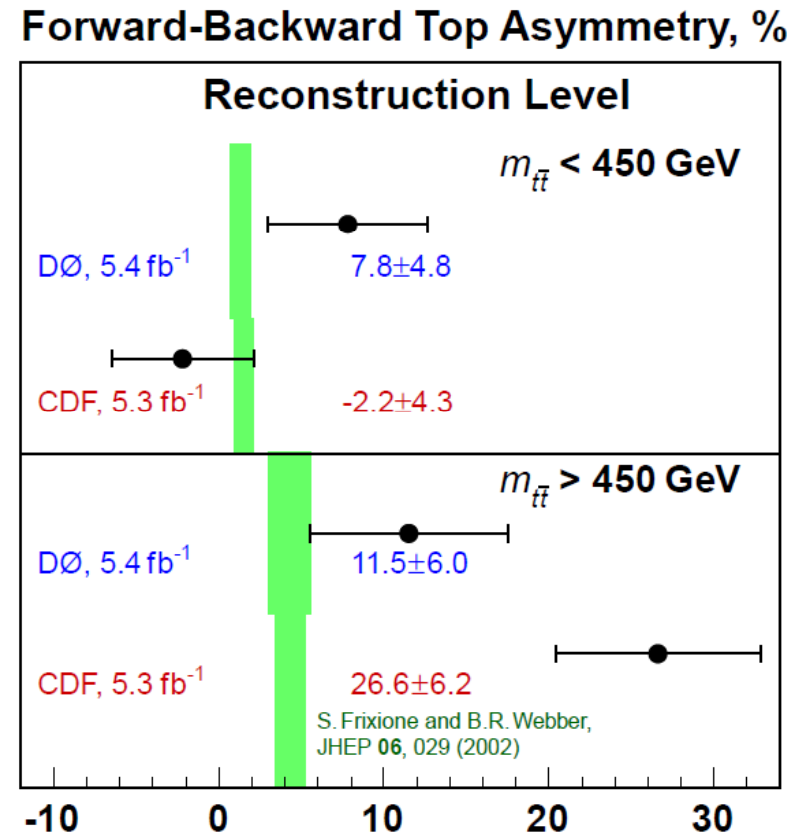
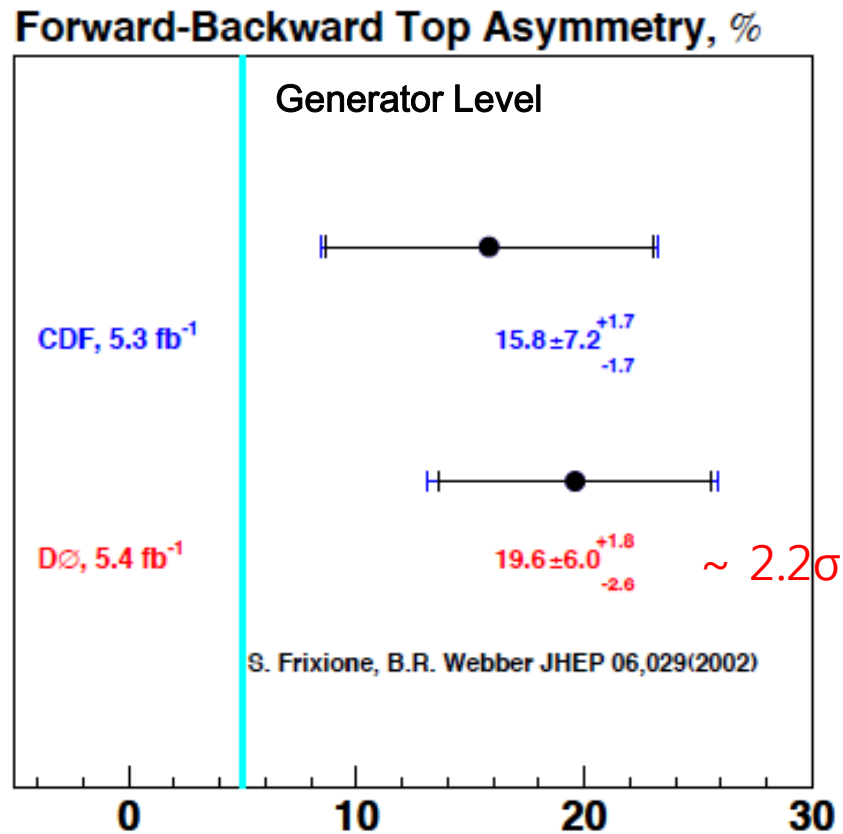
$$A_{FB}^{true} = 0.417 \pm 0.148 (stat.) \pm 0.053 (syst.)$$

$2.7\sigma$  from SM prediction  $0.05 \pm 0.01$

# Top $A_{FB}$ at D0

From Regina Demina's talk, EPS-HEP 2011

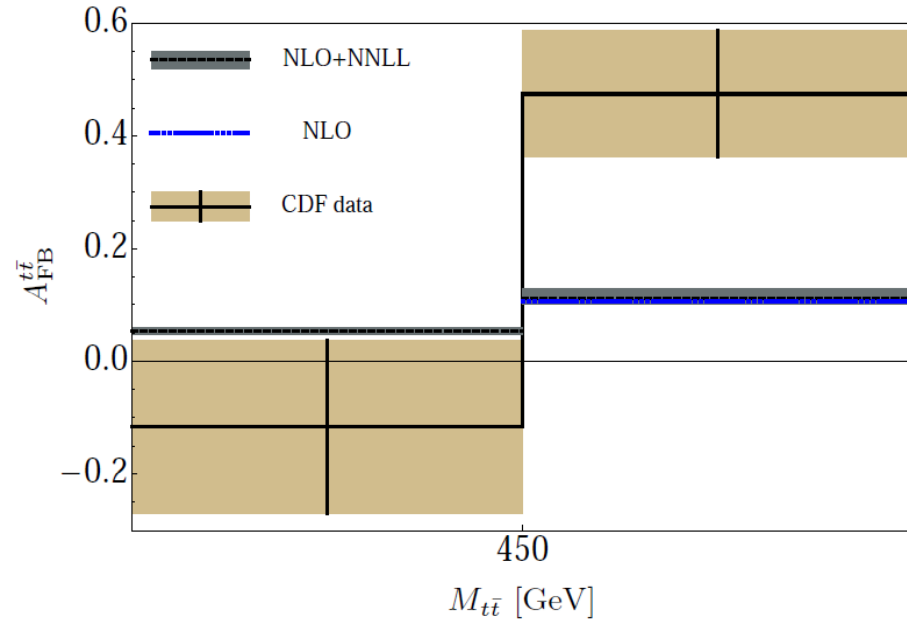
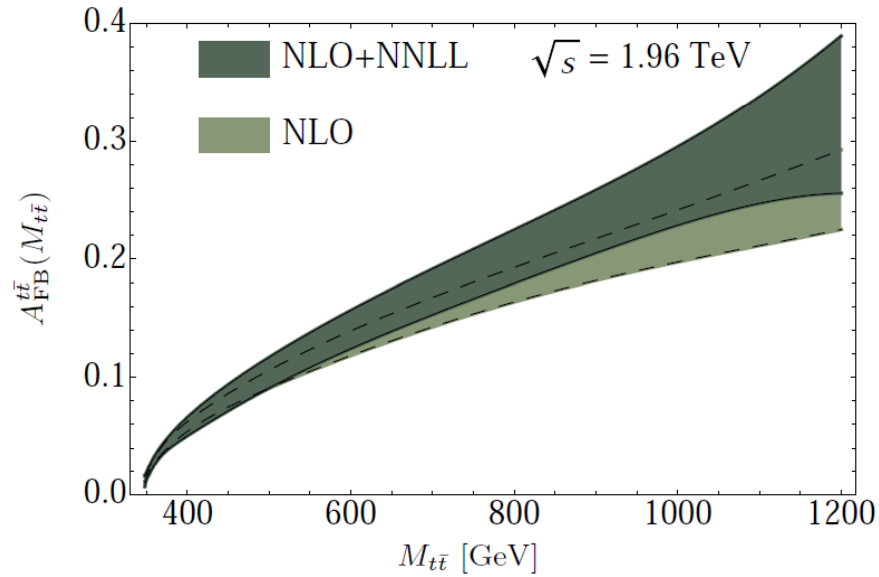
- Using  $5.4 \text{ fb}^{-1}$  of data in the lepton+jets channel



- consistent with CDF in the generator level.
- no convincing evidence for the dependence on  $m_{t\bar{t}}$ .

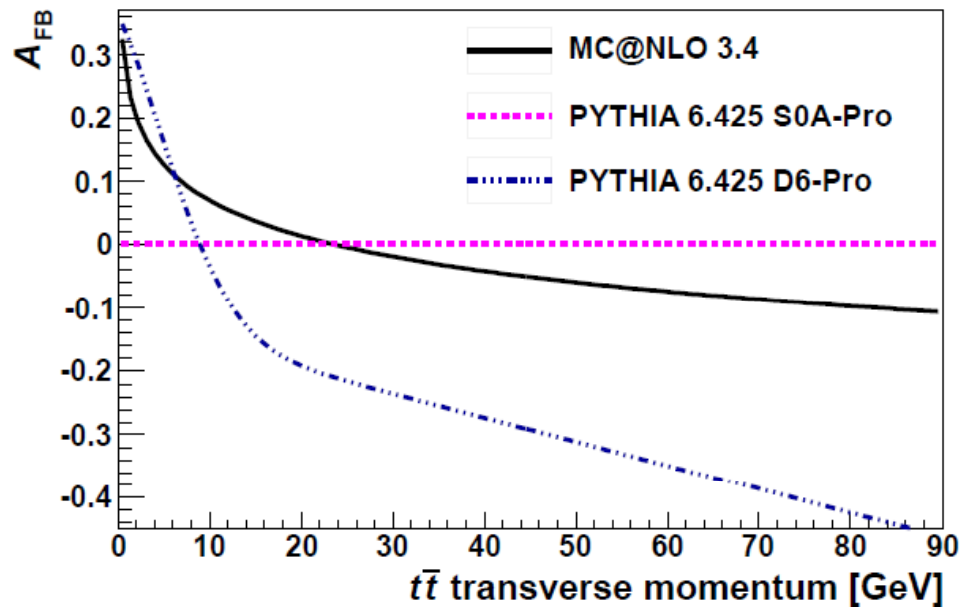
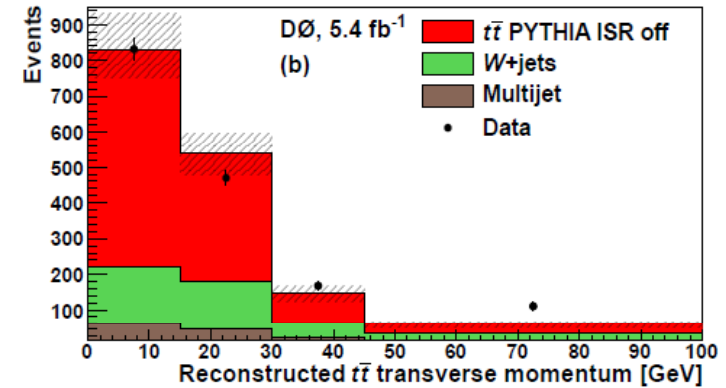
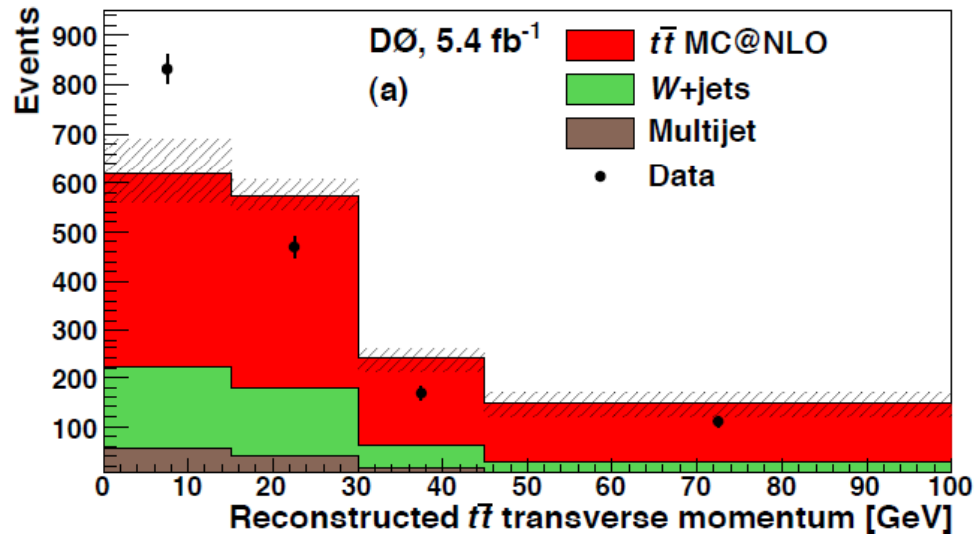
# SM prediction at NLO+NNLL

Ahrens, Ferroglia, Neubert, Pecjak, Yang, PRD84



(In units of %)	$A_{\text{FB}}^{t\bar{t}}$	$A_{\text{FB}}^{t\bar{t}} (M_{t\bar{t}} < 450 \text{ GeV})$	$A_{\text{FB}}^{t\bar{t}} (M_{t\bar{t}} > 450 \text{ GeV})$
NLO	$7.14^{+0.67}_{-0.54}$	$5.3^{+0.4}_{-0.4}$	$10.4^{+1.0}_{-0.6}$
<b>NLO+NNLL</b>	<b><math>7.16^{+1.05}_{-0.68}</math></b>	<b><math>5.2^{+0.8}_{-0.6}</math></b>	<b><math>10.8^{+1.7}_{-0.9}</math></b>
CDF	$15.8 \pm 7.5$	$-11.6 \pm 15.3$	$47.5 \pm 11.2$

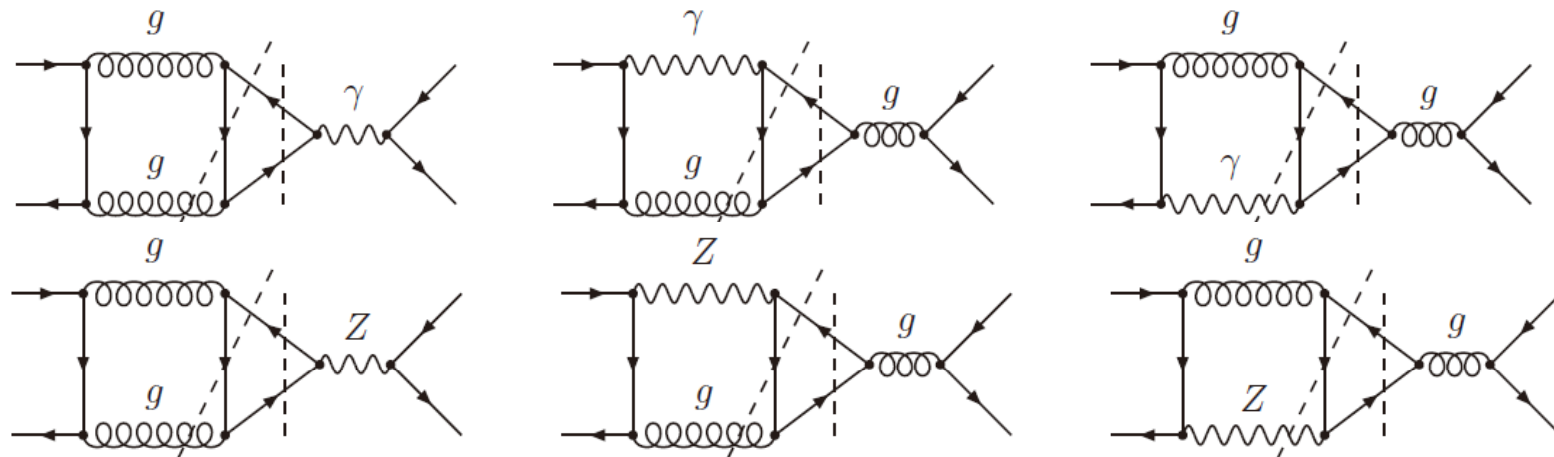
# $P_T^{t\bar{t}}$ dependence



- $D\bar{0}$  disagrees with MC@NLO.
- changes sign at  $\sim 20 \text{ GeV}$ .
- the asymmetry would be enhanced if the data are lost at high  $p_T$ .
- disagreement between MC generators.



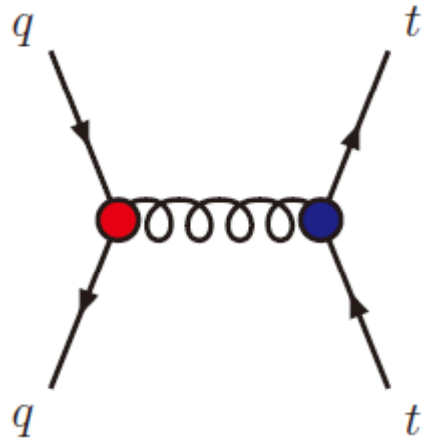
# EW contributions



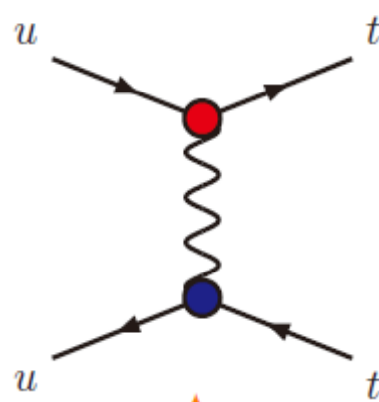
- reanalyze electromagnetic as well as weak corrections.
  - enhancement of  $A_{FB}$  by about a factor 1.1.
- restrict the  $t\bar{t}$  system to a transverse momentum  $< 20$  GeV.
  - enhancement of  $A_{FB}$  by factors between 1.3 and 1.5.

(In units of %)	$A_{FB}^{t\bar{t}}$	$A_{FB}^{t\bar{t}} (M_{t\bar{t}} < 450 \text{ GeV})$	$A_{FB}^{t\bar{t}} (M_{t\bar{t}} > 450 \text{ GeV})$
SM	$8.7 \pm 1.0$	$6.2 \pm 0.4$	$12.8 \pm 1.1$
MCFM	$5.8 \pm 0.9$	$4.0 \pm 0.6$	$8.8 \pm 1.3$
CDF	$15.8 \pm 7.5$	$-11.6 \pm 15.3$	$47.5 \pm 11.2$

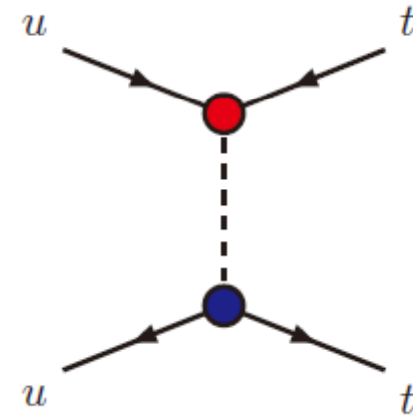
# New models



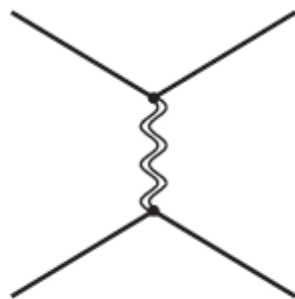
s-channel: coloured resonance  $\mathcal{G}_\mu$



t-channel:  $Z', W', \phi$

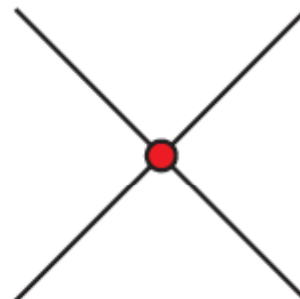


u-channel: exotic scalars



(new) heavy VB

Integrate

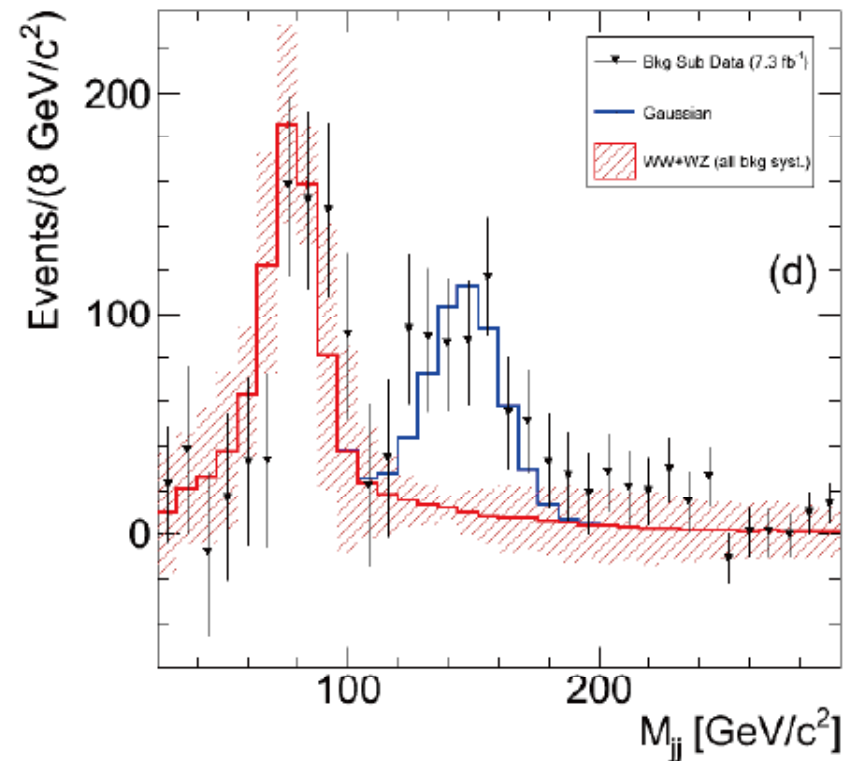
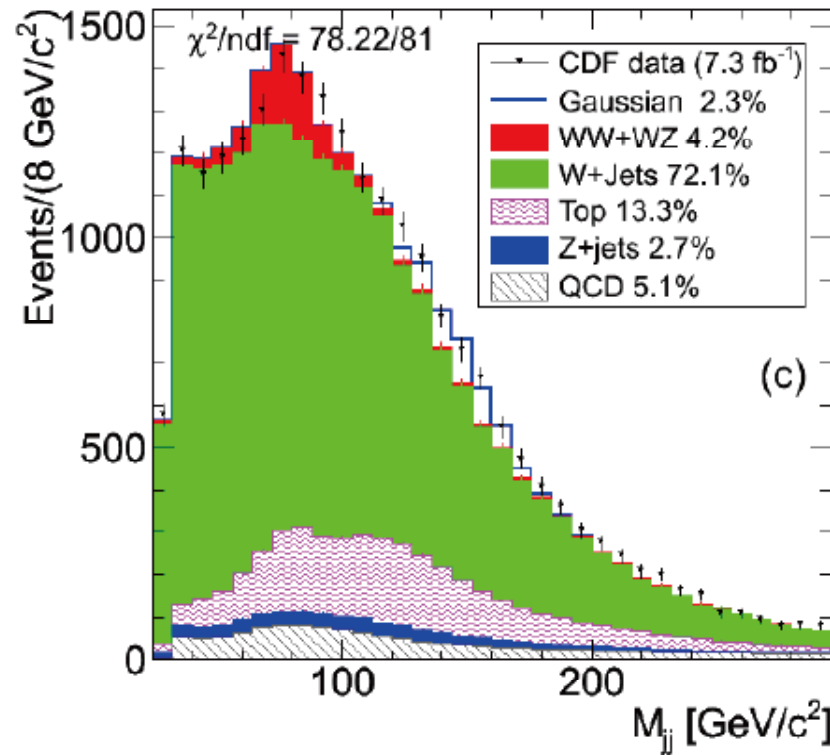


4-fermion interaction

- flavor dependent.
- challenging to construct a realistic model.
  - anomaly free,
  - renormalizable,
  - realistic Yukawa couplings.

# Wjj excess at CDF

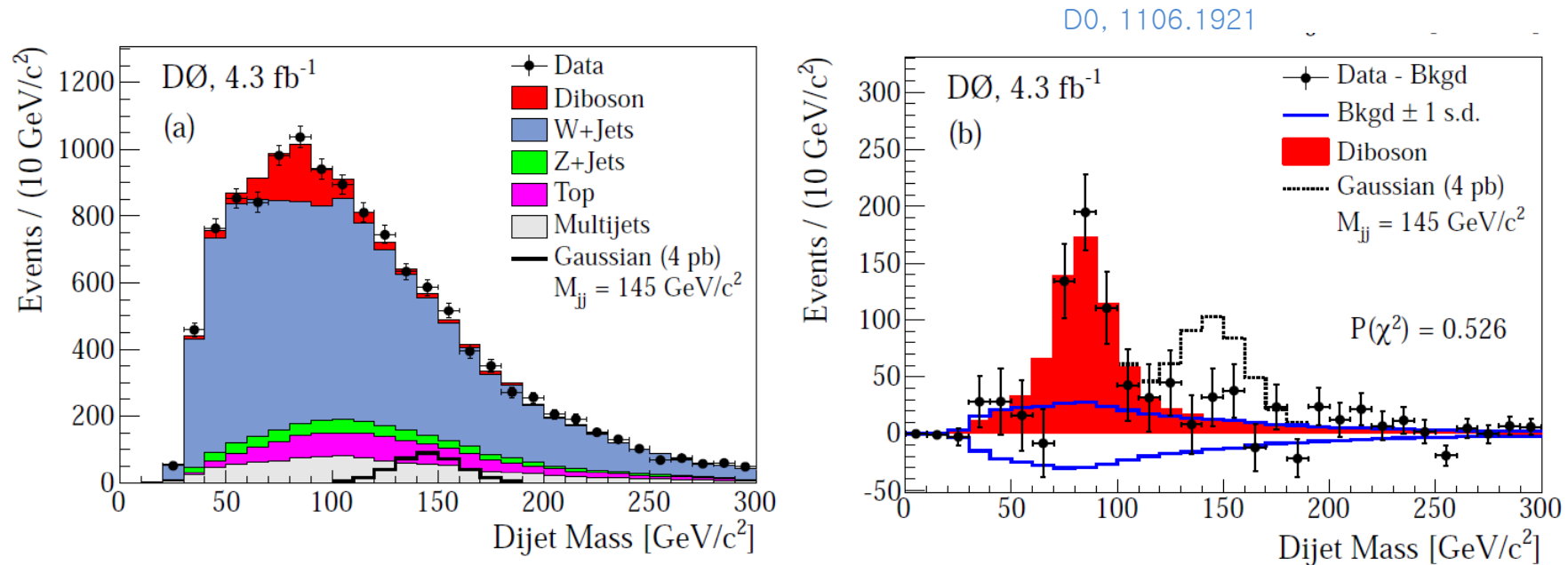
From Punzi's talk, Blois 2011



- 4.1 sigma deviation with  $7.3 \text{ fb}^{-1}$ .
- assume an additional Gaussian peak.
- $\sigma(p\bar{p} \rightarrow WX) \times Br(X \rightarrow jj) \sim 4 \text{ pb}$  with  $m_X \approx 145 \text{ GeV}$ .

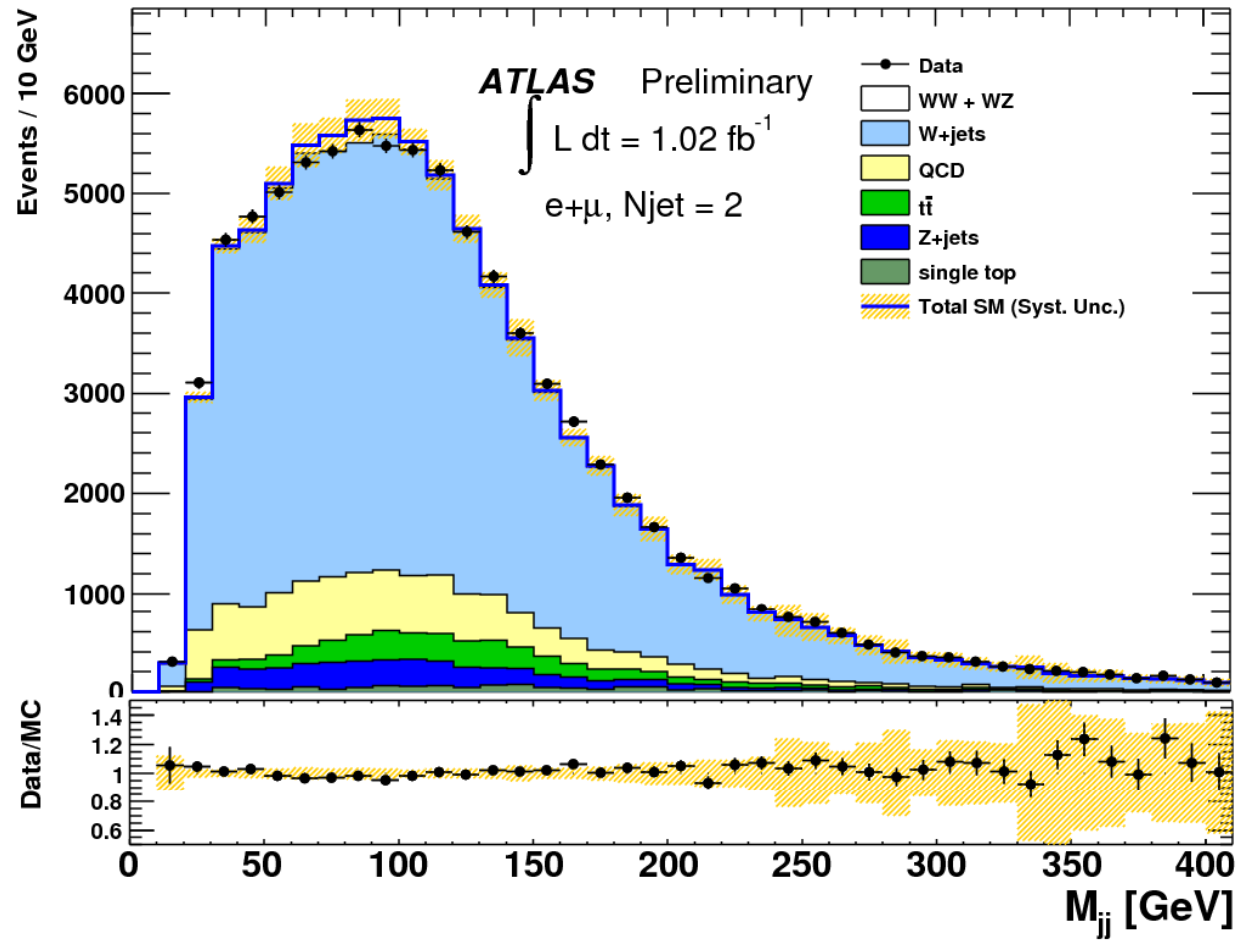
(See the talks by Eichten, Song, H.Kim.)

# Wjj at D0

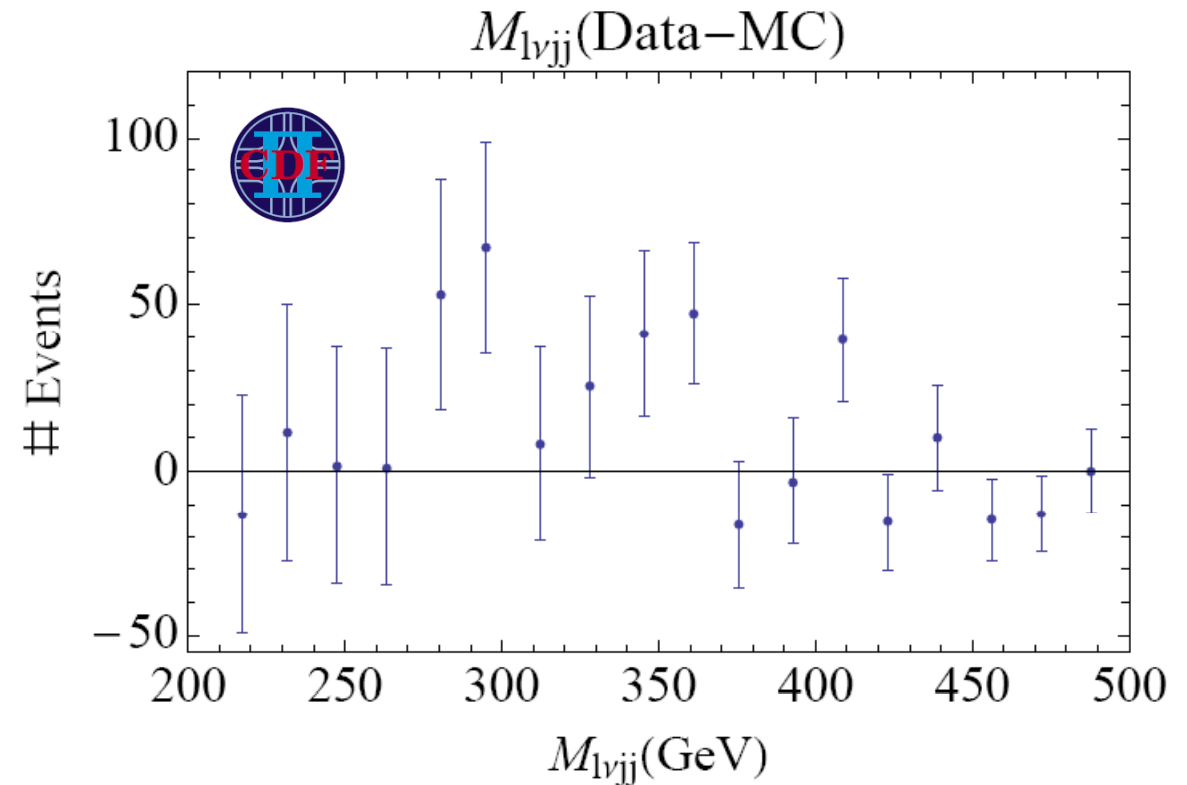
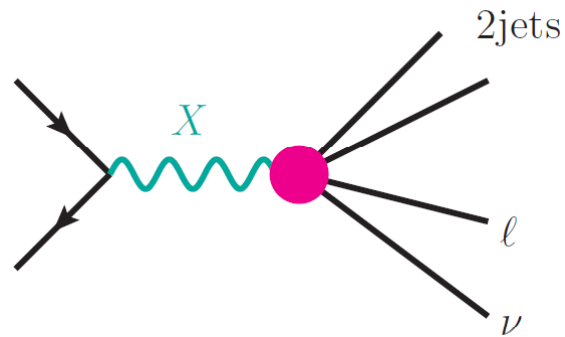


- no evidence for anomalous, resonant production of dijets.
- upper limits for the dijet production : 1.9 pb for  $m_{jj} = 145 \text{ GeV}$  (95% C.L.).
- reject the hypothesis of a cross section of 4 pb at the level of 4.3 s.d.

# Wjj at ATLAS



# s-channel resonance?



- difficult to derive any conclusion on an s-channel resonance.
- a broad excess in the region around  $M_{l\nu jj} \sim 270$  GeV?

# Flavor dependent $U(1)'$ Model

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- many studies for a relatively light  $Z'$  gauge boson with mass  $\sim 150$  GeV.
- the  $Z'$  is associated with some  $U(1)'$  gauge symmetry.
- better be leptophobic to avoid the LEP II and Drell-Yan bounds.
- approximately lighter than 200 GeV from the dijet production in the UA2 experiments.
- difficult to assign flavor-dependent charges to down-type quarks due to the strong constraints from FCNC experiments. (See the talk by Ligeti.)
- Yukawa interactions : additional Higgs fields.
- a flavor-dependent leptophobic  $U(1)'$  : anomalous.
  - introduce additional fermions to cancel the gauge anomalies.
- Both  $Z'$  and Higgs fields affect the top  $A_{FB}$  and  $W_{jj}$  production.

# Flavor dependent $U(1)'$ Model

- Charge assignment : SM fermions

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)'$
$Q_1$	3	2	1/6	$q_L$
$Q_2$	3	2	1/6	$q_L$
$Q_3$	3	2	1/6	$q_L$
$\overline{D}_1$	$\overline{3}$	1	1/3	$-q_L$
$\overline{D}_2$	$\overline{3}$	1	1/3	$-q_L$
$\overline{D}_3$	$\overline{3}$	1	1/3	$-q_L$
$\overline{U}_1$	$\overline{3}$	1	-2/3	$u_1$
$\overline{U}_2$	$\overline{3}$	1	-2/3	$u_2$
$\overline{U}_3$	$\overline{3}$	1	-2/3	$u_3$
$H$	1	2	1/2	0

Left-handed quarks and right-handed down-type quarks have universal couplings.

Flavor-dependent

Higgs



# Flavor dependent U(1)' Model

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- Charge assignment : Higgs fields

	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)'$
$H_1$	1	2	1/2	$-q_L - u_1$
$H_2$	1	2	1/2	$-q_L - u_2$
$H_3$	1	2	1/2	$-q_L - u_3$
$\Phi$	1	1	1	$-q_\Phi$

- introduce three Higgs doublets charged under U(1)' in addition to H uncharged under U(1)'.

$$\begin{aligned}
 V_y = & y_{i1}^u H_1 \bar{U}_1 Q_i + y_{i2}^u H_2 \bar{U}_2 Q_i + y_{i3}^u H_3 \bar{U}_3 Q_i \\
 & + y_{ij}^d \bar{D}_j Q_i i\tau_2 H^\dagger \\
 & + y_{ij}^e \bar{E}_j L_i i\tau_2 H^\dagger + y_{ij}^n H \bar{N}_j L_i.
 \end{aligned}$$

- The U(1)' is spontaneously broken by U(1)' charged complex scalar  $\Phi$ .

# Flavor dependent $U(1)'$ Model

- Anomaly cancelation requires extra fermions I:  $SU(2)$  doublets

	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)'$
$Q'$	3	2	1/6	$-(q_1 + q_2 + q_3)$
$D'_R$	3	1	-1/3	$-(d_1 + d_2 + d_3)$
$U'_R$	3	1	2/3	$-(u_1 + u_2 + u_3)$
$L'$	1	2	-1/2	0
$E'$	1	1	-1	0
$l_{L1}$	1	2	-1/2	$Q_L$
$l_{R1}$	1	2	-1/2	$Q_R$
$l_{L2}$	1	2	-1/2	$-Q_L$
$l_{R2}$	1	2	-1/2	$-Q_R$

one extra generation

$SU(2)_L^2 \cdot U(1)'$

vector-like pairs

$U(1)'^2 \cdot U(1)$

a candidate for CDM

(See the talk by Omura.)

# Flavor dependent U(1)' Model

- 2 Higgs doublet model :  $(u_1, u_2, u_3) = (0, 0, 1)$

	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)'$
$H$	1	2	1/2	0
$H_3$	1	2	1/2	1
$\Phi$	1	1	1	$q_\Phi$

$$V_y = y_{i1}^u \overline{Q}_i \widetilde{H} U_{R1} + y_{i2}^u \overline{Q}_i \widetilde{H} U_{Rj} + y_{i3}^u \overline{Q}_i \widetilde{H}_3 U_{Rj} \\ + y_{ij}^d \overline{Q}_i H D_{Rj} + y_{ij}^e \overline{L}_i H \overline{E}_j + y_{ij}^n \overline{L}_i \widetilde{H} N_j.$$

$$V_h = Y_{ij}^u \overline{\hat{U}}_{Li} \hat{U}_{Rj} \hat{h}_0 + Y_{ij}^d \overline{\hat{D}}_{Li} \hat{D}_{Rj} \hat{h}_0,$$

$$Y_{ij}^u = \frac{m_i^u \cos \alpha}{v \cos \beta} \delta_{ij} + \frac{2m_i^u}{v \sin 2\beta} (g_R^u)_{ij} \sin(\alpha - \beta),$$

$$Y_{ij}^d = \frac{m_i^d \cos \alpha}{v \cos \beta} \delta_{ij},$$

}  $\propto$  the fermion mass

# Flavor dependent $U(1)'$ Model

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- **3 Higgs doublet model:**  $(u_1, u_2, u_3) = (-q, 0, q)$

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)'$
$H_1$	1	2	1/2	$q$
$H_2$	1	2	1/2	0
$H_3$	1	2	1/2	$-q$
$\Phi$	1	1	0	$-1$

$$\begin{aligned} \mathcal{L}_Y = & y_{i1}^u H_1 \bar{U}_1 Q_i + y_{i2}^u H_2 \bar{U}_2 Q_i + y_{i3}^u H_3 \bar{U}_3 Q_i \\ & + y_{ij}^d H_2^\dagger \bar{D}_j Q_i + y_{ij}^e H_2^\dagger \bar{E}_j L_i + y_{ij}^n H_2 \bar{N}_j L_i. \end{aligned}$$

# Flavor dependent U(1)' Model

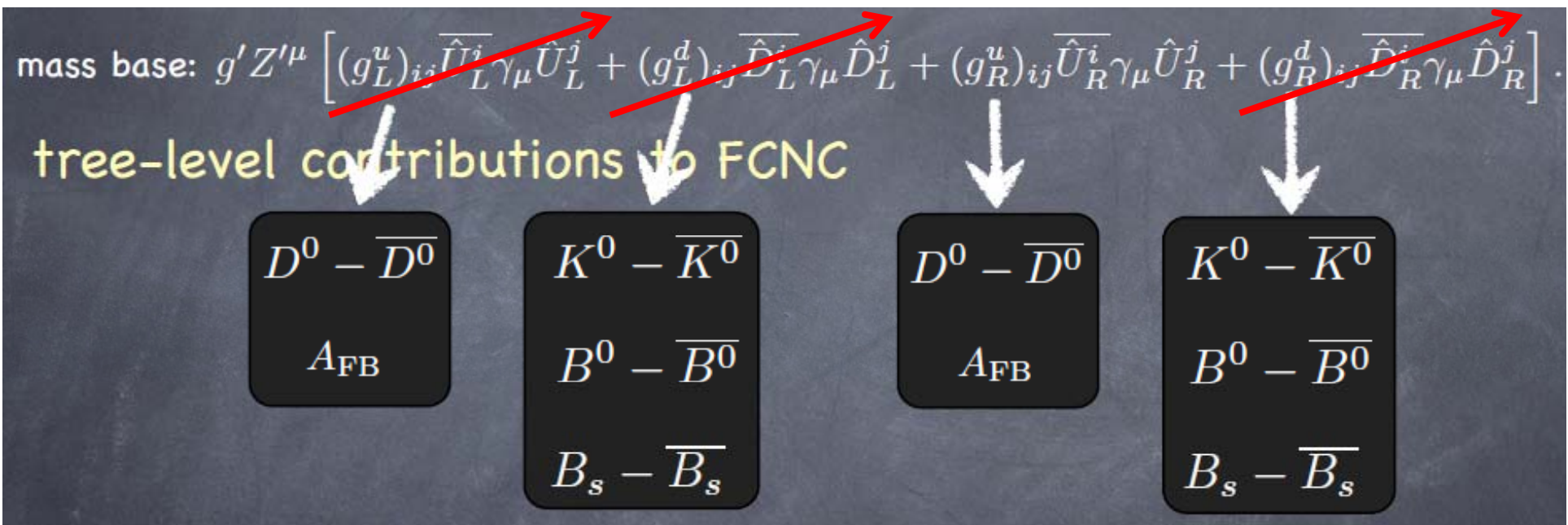
- Gauge coupling in the mass base
- Z' interacts only with the right-handed up-type quarks

$$g' Z'^{\mu} \sum_{i,j=1,2,3} (g_R^u)_{ij} \overline{U}_R^i \gamma_{\mu} U_R^j$$

- The 3 X 3 coupling matrix  $g_R^u$  is defined by

$$(g_R^u)_{ij} = (U_R^u)_{ik} U_R^{\dagger}{}_{kj}$$

biunitary matrix diagonalizing the up-type quark mass matrix



# Flavor dependent U(1)' Model

- Yukawa coupling in the mass base (2HDM)

- lightest Higgs h:  $V_h = Y_{ij}^u \overline{\hat{U}}_{Li} \hat{U}_{Rj} h + Y_{ij}^d \overline{\hat{D}}_{Li} \hat{D}_{Rj} h + Y_{ij}^e \overline{\hat{E}}_{Li} \hat{E}_{Rj} h + h.c.,$

$$Y_{ij}^u = \frac{m_i^u \cos \alpha}{v \cos \beta} \cos \alpha_\Phi \delta_{ij} + \frac{2m_i^u}{v \sin 2\beta} (g_R^u)_{ij} \sin(\alpha - \beta) \cos \alpha_\Phi,$$

$$Y_{ij}^d = \frac{m_i^d \cos \alpha}{v \cos \beta} \cos \alpha_\Phi \delta_{ij},$$

$$Y_{ij}^e = \frac{m_i^l \cos \alpha}{v \cos \beta} \cos \alpha_\Phi \delta_{ij},$$

- lightest charged Higgs h<sup>±</sup>:  $V_{h^\pm} = -Y_{ij}^{u-} \overline{\hat{D}}_{Li} \hat{U}_{Rj} h^- + Y_{ij}^{d+} \overline{\hat{U}}_{Li} \hat{D}_{Rj} h^+ + h.c.,$

$$Y_{ij}^{u-} = \sum_l (V_{\text{CKM}})_{li}^* \left\{ \frac{\sqrt{2}m_l^u \tan \beta}{v} \delta_{lj} - \frac{2\sqrt{2}m_l^u}{v \sin 2\beta} (g_R^u)_{lj} \right\},$$

$$Y_{ij}^{d+} = (V_{\text{CKM}})_{ij} \frac{\sqrt{2}m_j^d \tan \beta}{v},$$

- lightest pseudoscalar Higgs a:  $V_a = -iY_{ij}^{au} \overline{\hat{U}}_{Li} \hat{U}_{Rj} a + iY_{ij}^{ad} \overline{\hat{D}}_{Li} \hat{D}_{Rj} a + iY_{ij}^{ae} \overline{\hat{E}}_{Li} \hat{E}_{Rj} a + h.c.,$

$$Y_{ij}^{au} = \frac{m_i^u \tan \beta}{v} \delta_{ij} - \frac{2m_i^u}{v \sin 2\beta} (g_R^u)_{ij},$$

$$Y_{ij}^{ad} = \frac{m_i^d \tan \beta}{v} \delta_{ij},$$

$$Y_{ij}^{ae} = \frac{m_i^l \tan \beta}{v} \delta_{ij}.$$

# Top-antitop-quark pair production at Tevatron

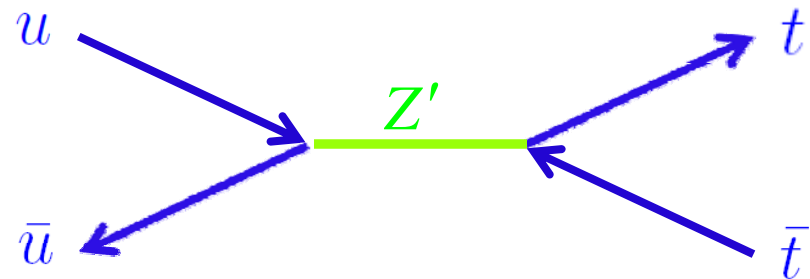
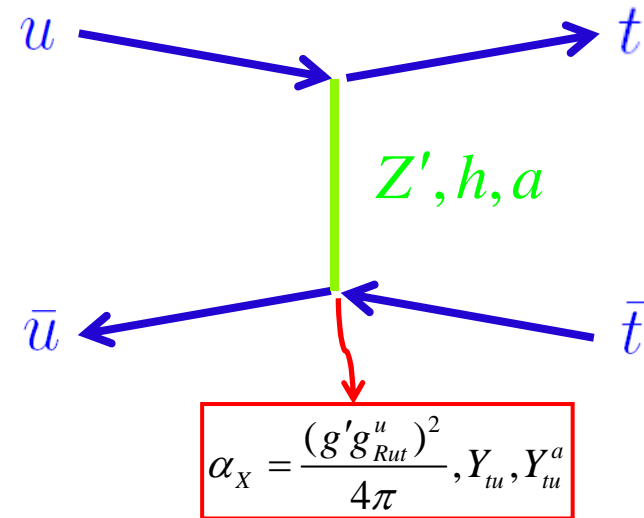
## 1. Z' dominant scenario

cf. Jung, Murayama, Pierce, Wells, PRD81(2010)

## 2. Higgs dominant scenario

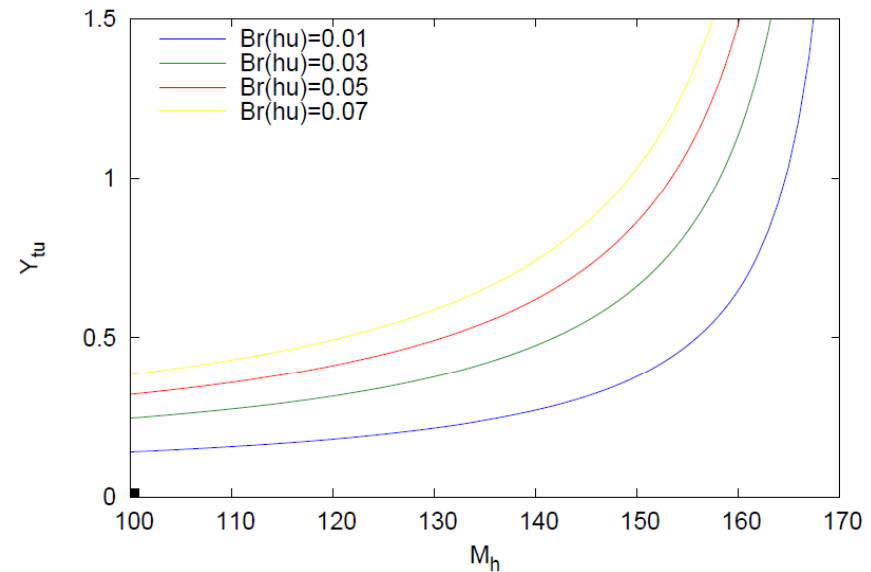
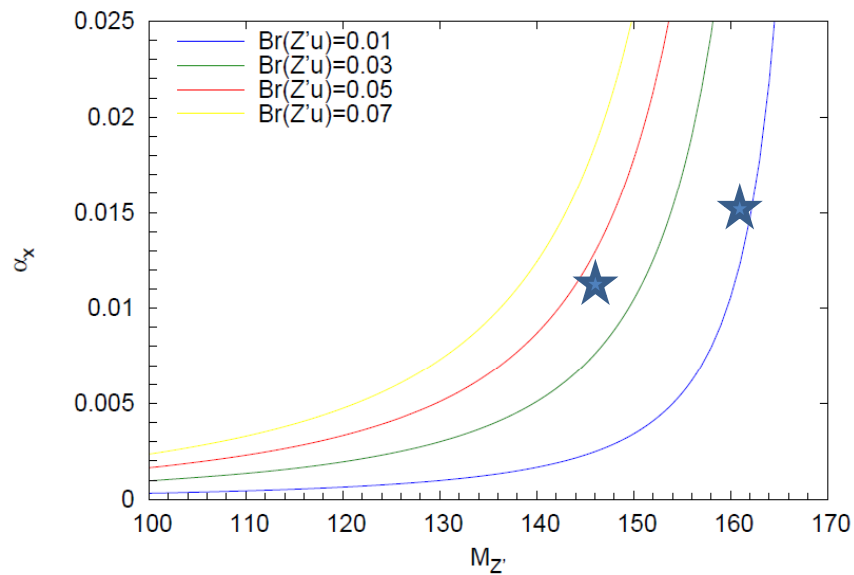
cf. Babu, Frank, Rai, 1104.4782

## 3. Mixed scenario



# Top quark decay

- decay into  $W+b$  in SM :  $\text{Br}(t \rightarrow Wb) \sim 100\%$ .
- If the top quark decays to  $Z' + u$  or  $h + u$ ,  $\text{Br}(t \rightarrow Wb)$  might significantly be changed.

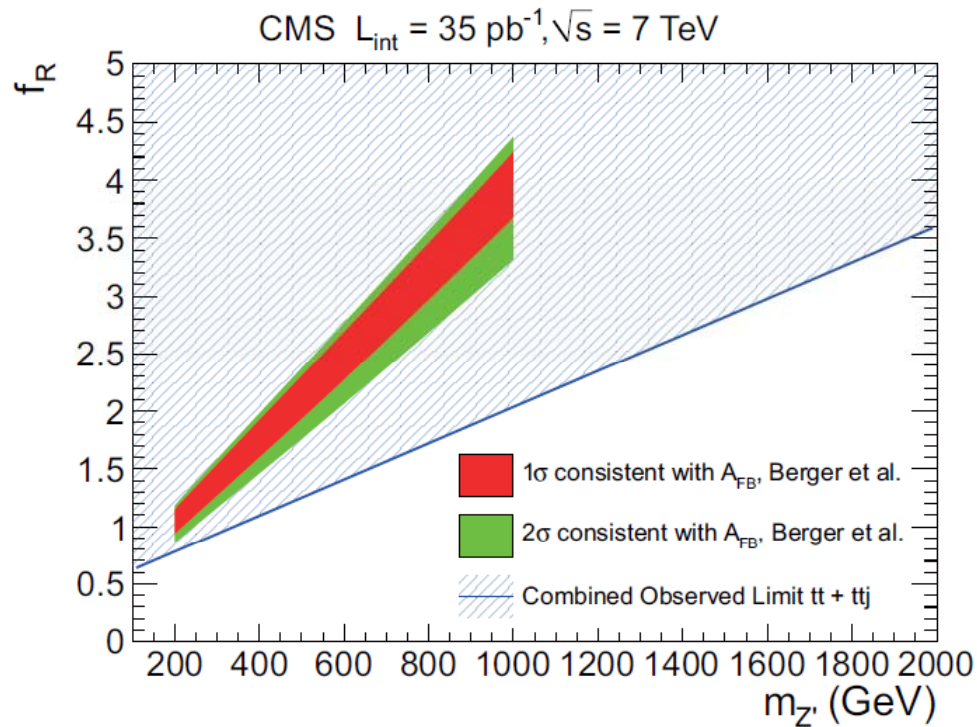


- assume  $\text{Br}(t \rightarrow Z'u) < 5\%$  and  $m_h > m_t$ .
- $m_{Z'} = 145$  or  $160$  GeV and  $m_h = 180$  GeV.

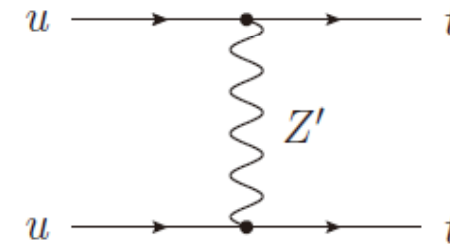


# Same sign top quark pair production at LHC

$$\mathcal{L} = g_W \bar{u} \gamma^\mu (f_L P_L + f_R P_R) t Z'_\mu + \text{h.c.},$$



CMS, 1106.2142



Top FCNC

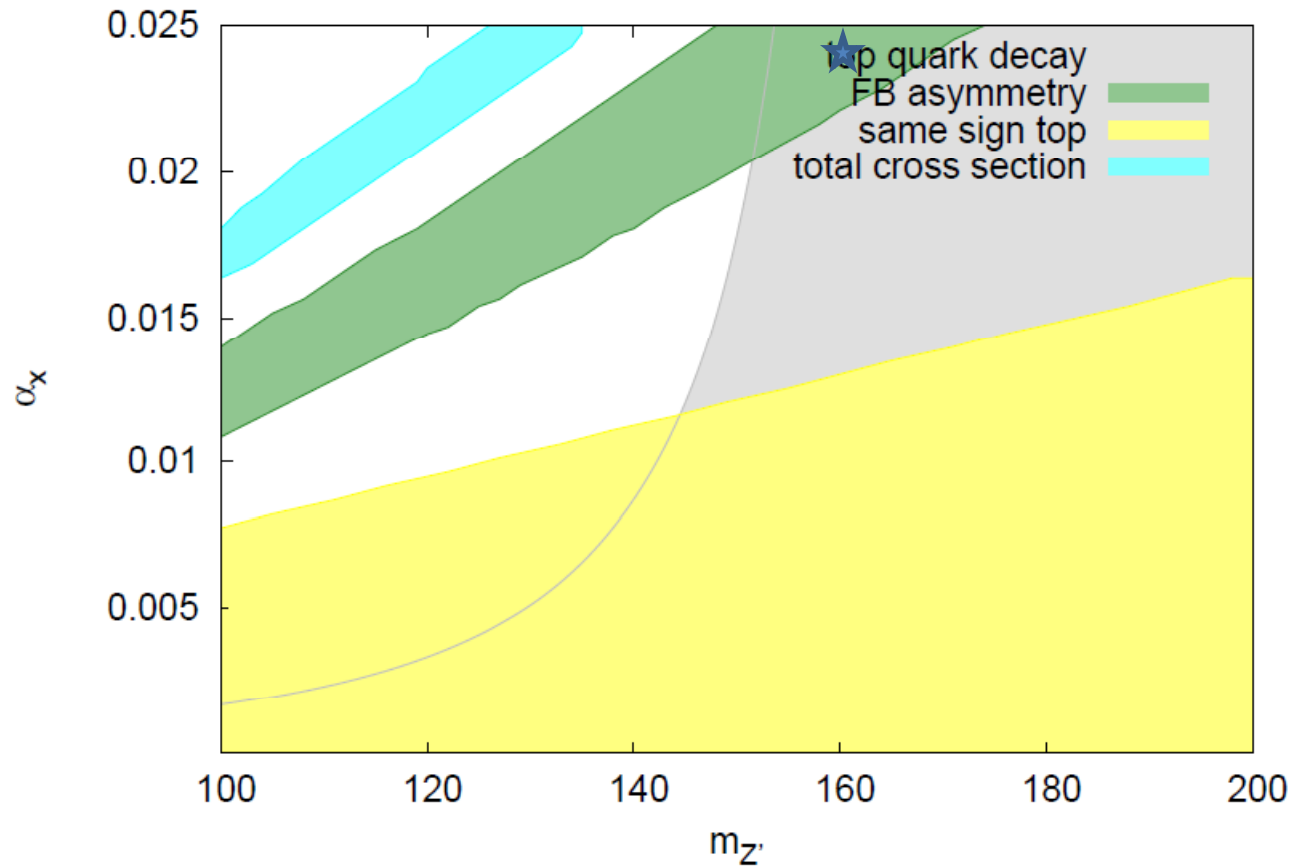
$$pp \rightarrow tt(j) < 17 \text{ pb at } 95 \text{ CL.}$$

Can avoid this constraint?

- Interference between  $Z'$  and scalar and pseudoscalar Higgs bosons.

# Favored region

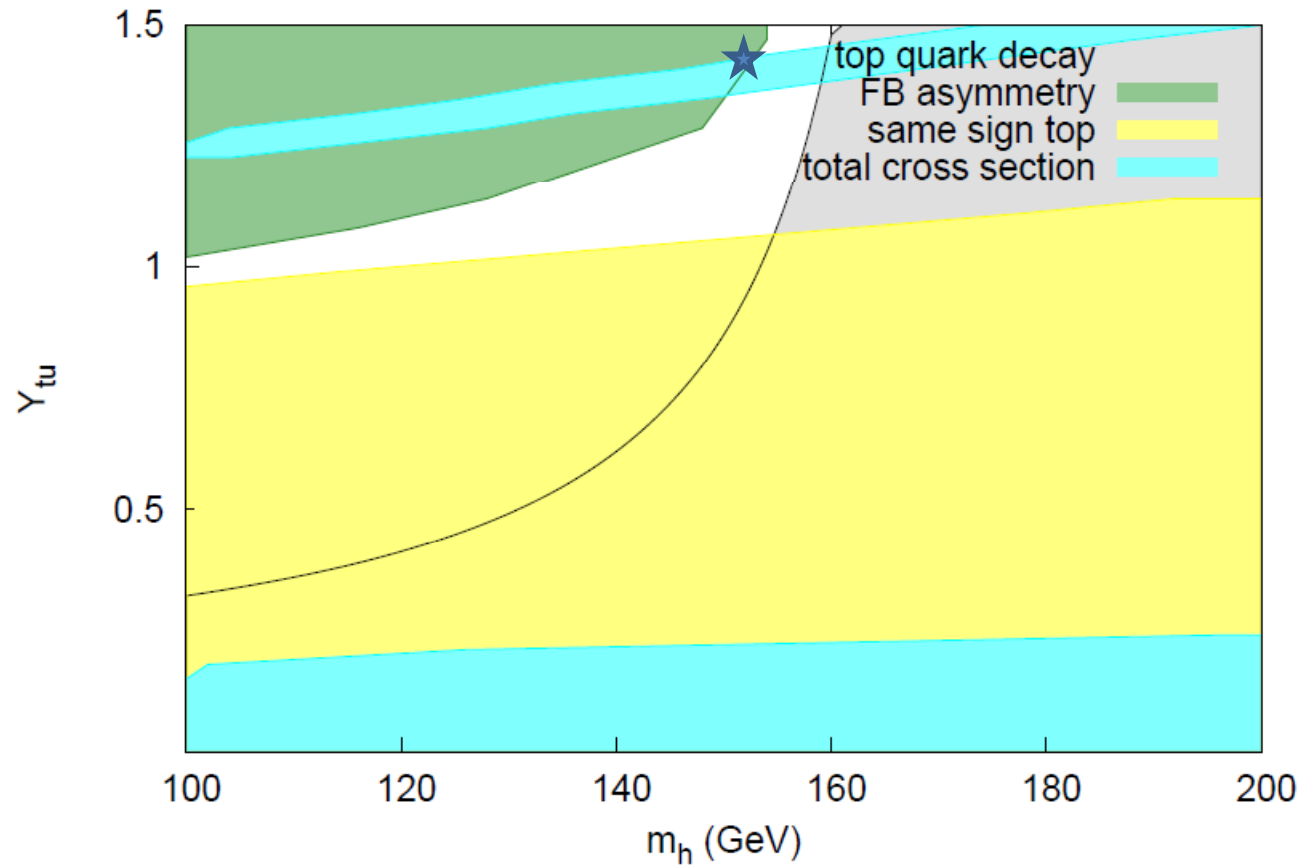
Z' dominant case



★ = Jung, Murayama, Pierce, Wells' model [PRD81,015004 \(2010\)](#)

# Favored region

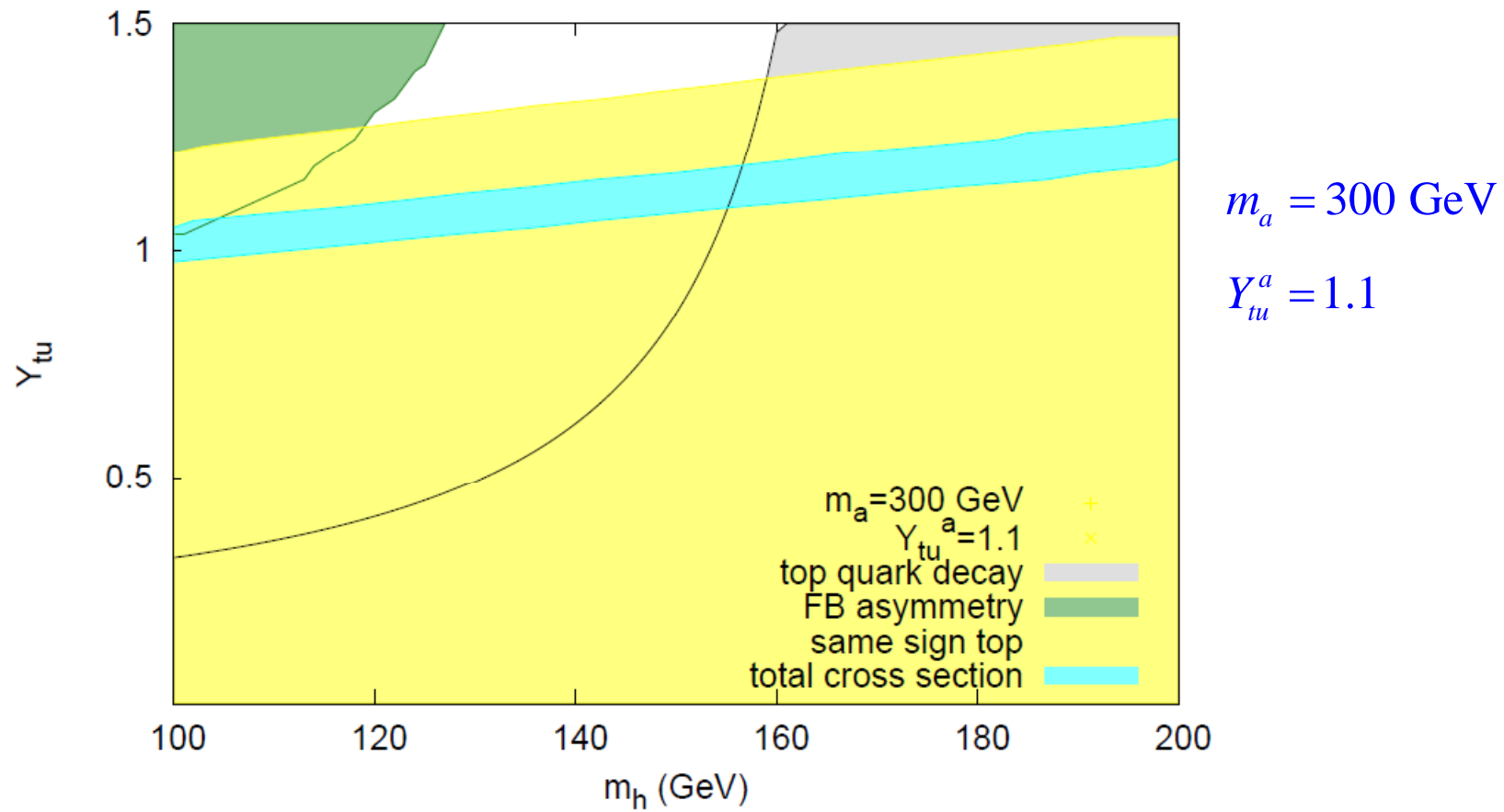
Scalar Higgs (h) dominant case



★ = Babu, Frank, Rai's model 1104.4782

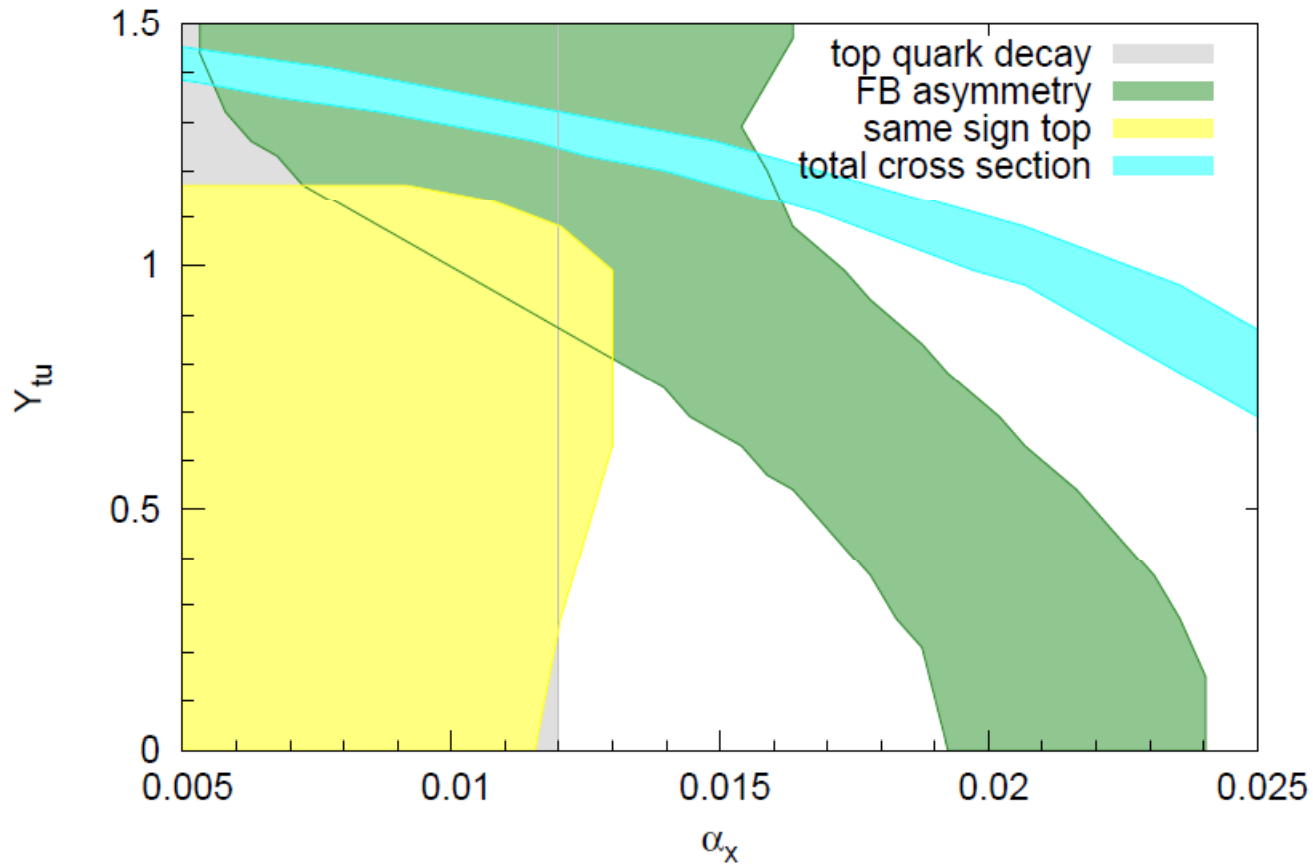
# Favored region

scalar (h) + pseudoscalar (a) Higgs case



# Favored region

Z'+h case



$$m_{Z'} = 145 \text{ GeV}$$

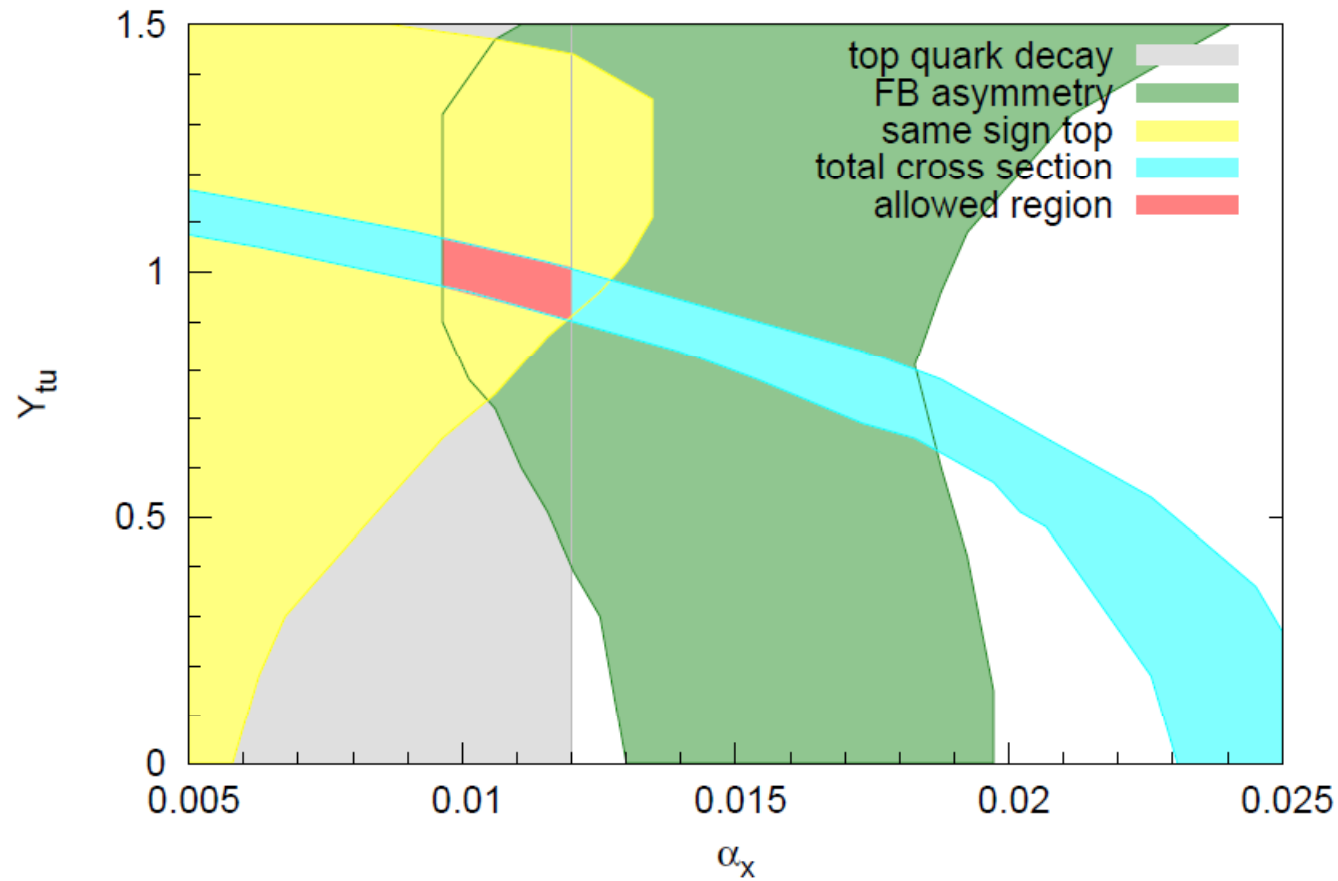
$$m_h = 180 \text{ GeV}$$

$$m_a = 300 \text{ GeV}$$

$$Y_{tu}^a = 1.1$$

# Favored region

Z'+h+a case



$$m_{Z'} = 145 \text{ GeV}$$

$$m_h = 180 \text{ GeV}$$

$$m_a = 300 \text{ GeV}$$

$$Y_{tu}^a = 1.1$$

$$A_{\text{FB}} = 0.084 \sim 0.12$$

## $A_{\text{FB}}$ versus $\sigma_{\text{tt}}$

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$$m_{Z'} = 145 \text{ GeV}$$

$$180 \text{ GeV} < m_h < 1 \text{ TeV}$$

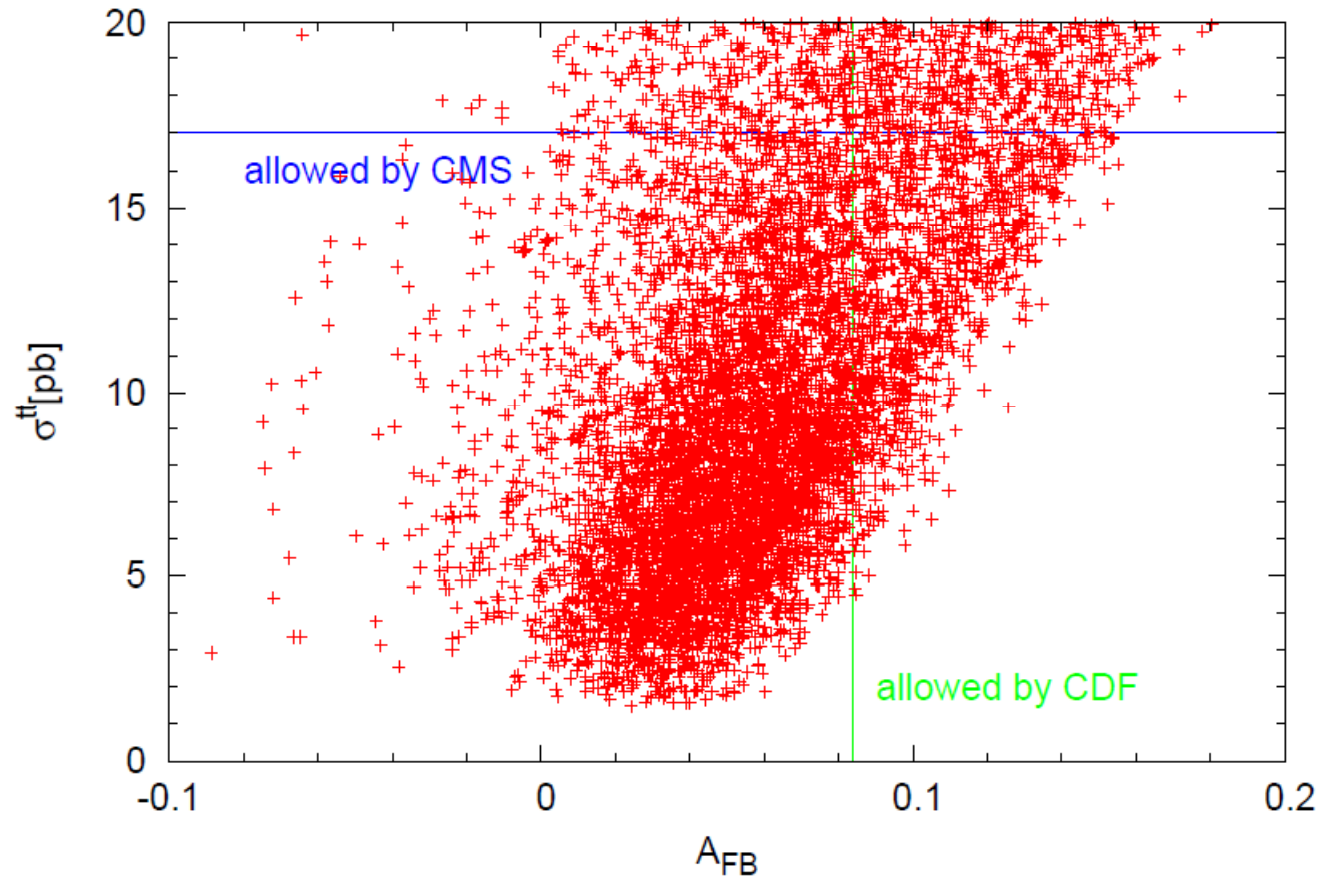
$$180 \text{ GeV} < m_a < 1 \text{ TeV}$$

$$0.005 < \alpha_x < 0.025$$

$$0.5 < Y_{tu} < 1.5$$

$$0.5 < Y_{tu}^a < 1.5$$

# $A_{FB}$ versus $\sigma_{tt}$



$$m_{Z'} = 145 \text{ GeV}$$

$$180 \text{ GeV} < m_h < 1 \text{ TeV}$$

$$180 \text{ GeV} < m_a < 1 \text{ TeV}$$

$$0.005 < \alpha_X < 0.025$$

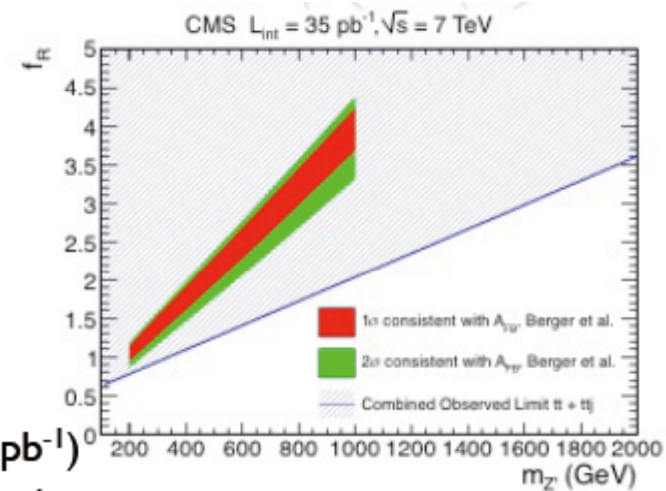
$$0.5 < Y_{tu} < 1.5$$

$$0.5 < Y_{tu}^a < 1.5$$



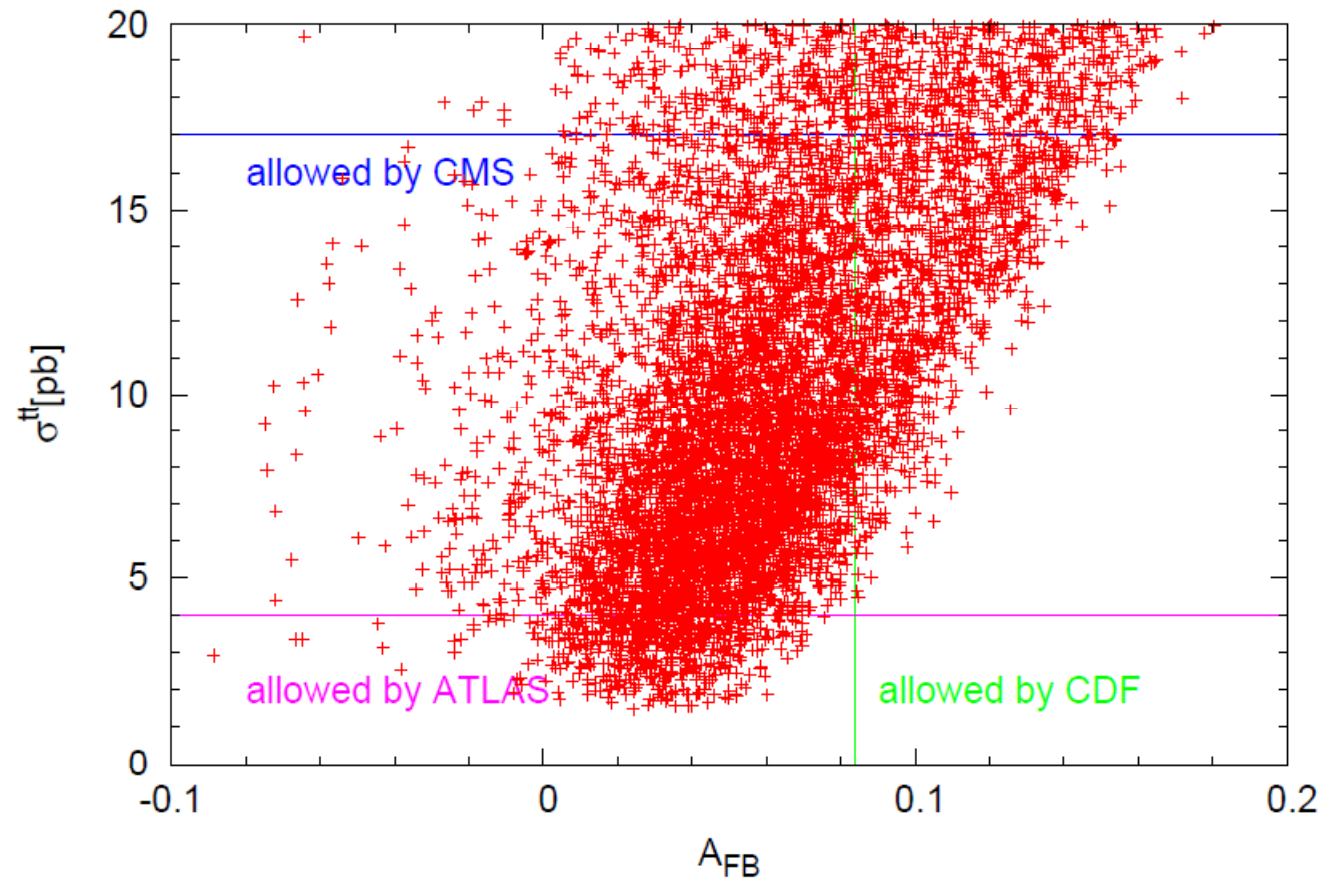
# Same sign top pair production at ATLAS

- ATLAS and CMS have searched for flavor-changing neutral currents in the top sector
  - CMS EXO-11-065, JHEP 1108 (2011) 005
  - ATLAS CONF-2011-169
- Limits are set on cross section of  $m$  positive-sign top pairs
  - CMS, ee,mumu, emu:  $\sigma(Z' \rightarrow ttX) < 17.0 \text{ pb} \text{ (} 35 \text{ pb}^{-1}\text{)}$
  - ATLAS, mumu:  $\sigma(Z' \rightarrow ttX) < 2.9 - 4.0 \text{ pb} \text{ (} 1.6 \text{ fb}^{-1}\text{)}$



From Blekman's talk, TOP2011

# $A_{FB}$ versus $\sigma_{tt}$



## $A_{\text{FB}}$ with SM NLO contribution

- In the SM,

$$A_{\text{FB}}^{\text{SM}} = \frac{\sigma_{\text{LO},\text{F}} + \sigma_{\text{NLO},\text{F}} - \sigma_{\text{LO},\text{B}} - \sigma_{\text{NLO},\text{B}}}{\sigma_{\text{LO}} + \sigma_{\text{NLO}}} = \frac{\cancel{\Delta\sigma_{\text{LO}}} + \Delta\sigma_{\text{NLO}}}{\sigma_{\text{LO}} + \sigma_{\text{NLO}}} \sim 8.7\%.$$

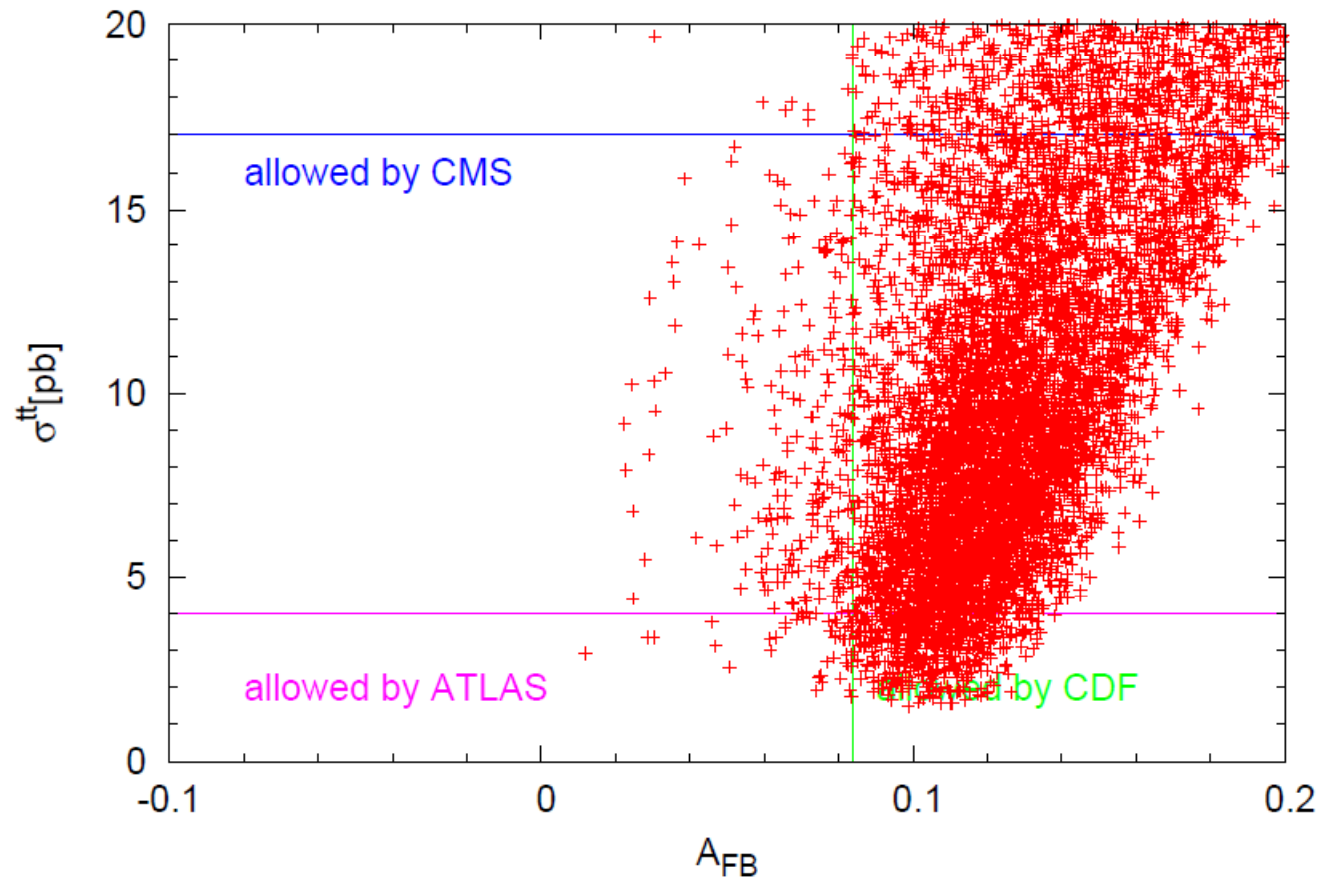
- In our calculation,

$$A_{\text{FB}}^{\text{New}} = \frac{K(\Delta\sigma_{\text{LO}} + \Delta\sigma_{\text{NEW}})}{K(\sigma_{\text{LO}} + \sigma_{\text{NEW}})} = \frac{\cancel{\Delta\sigma_{\text{LO}}} + \Delta\sigma_{\text{NEW}}}{\sigma_{\text{LO}} + \sigma_{\text{NEW}}} \sim 12\%.$$

- Consider both contributions of NLO and New physics,

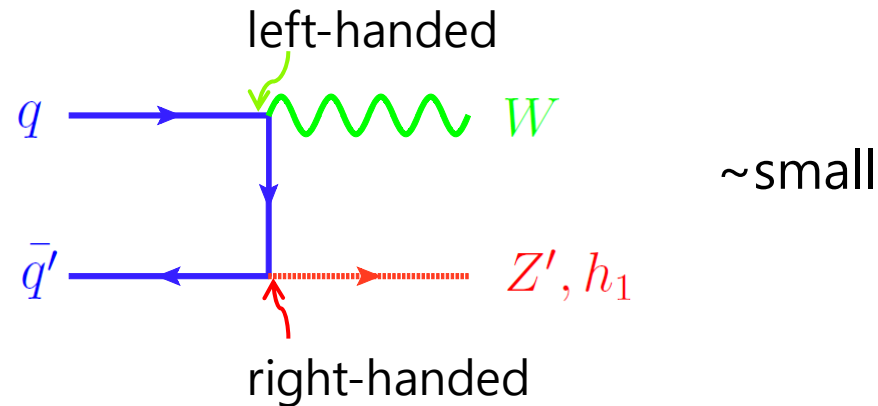
$$A_{\text{FB}} = A_{\text{FB}}^{\text{SM}} + A_{\text{FB}}^{\text{New}} / K \sim 18\%.$$

# $A_{\text{FB}}$ versus $\sigma_{t\bar{t}}$

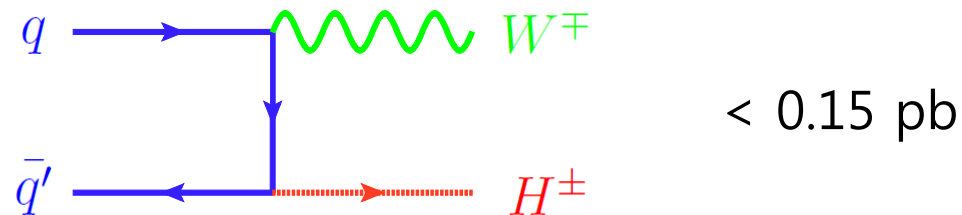


# W+jj anomaly and flavor dependent U(1)' model

- Since the Z' boson and Higgs boson dominantly couples to the right-handed top quark, their contributions to the W+jj production are small.

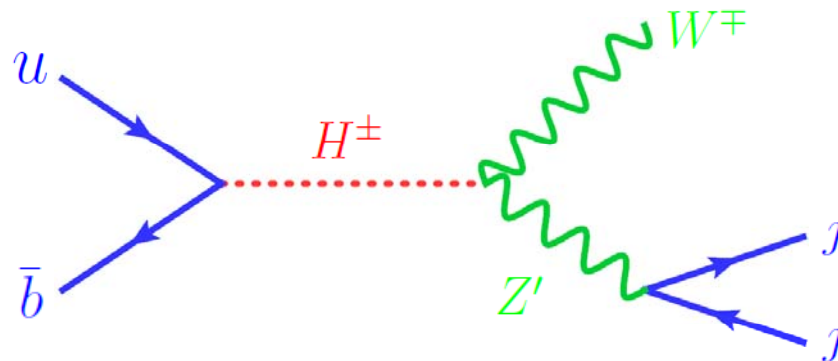


- Charged Higgs boson may contribute to the W+jj anomaly?



# W+jj anomaly and flavor dependent U(1)' model

$$\mathcal{L} = -g' m_W \sin 2\beta H^+ W^{-\mu} Z'_\mu + h.c..$$



$$m_{Z'} = 145 \text{ GeV}$$

$$m_{h^+} = 270 \text{ GeV}$$

$$\sigma(Wjj) \lesssim O(10) \text{ pb} \times \sin^2 2\beta$$

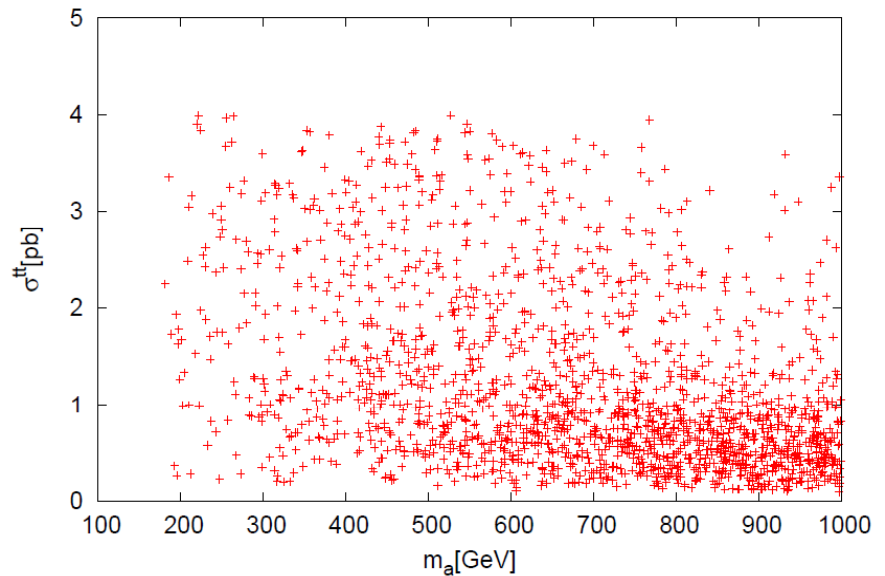
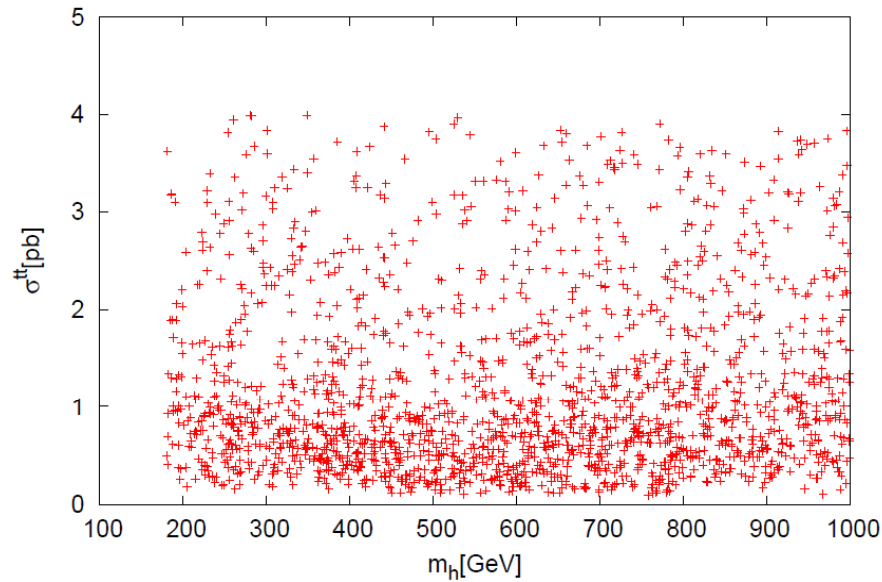
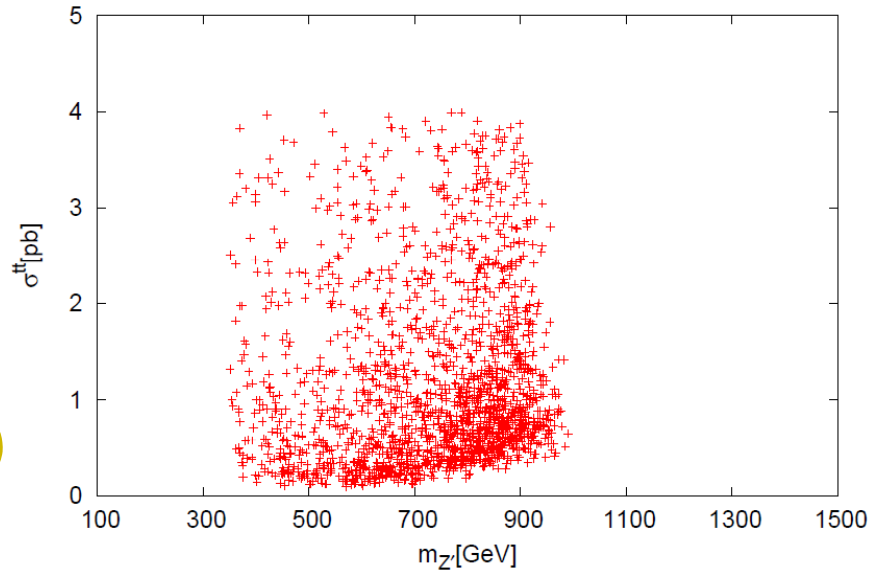
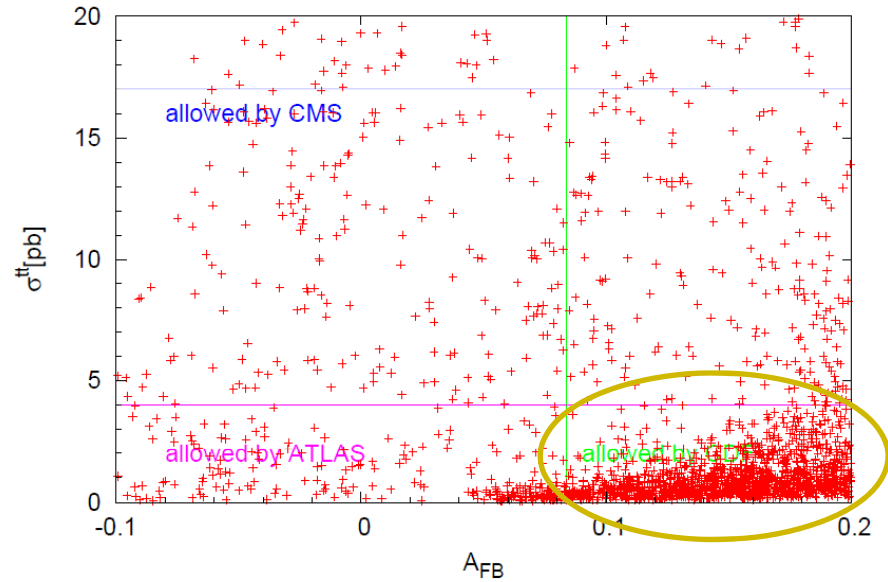
## More general case

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$$140 \text{ GeV} < m_{Z'} < 1.5 \text{ TeV}$$

# More general case

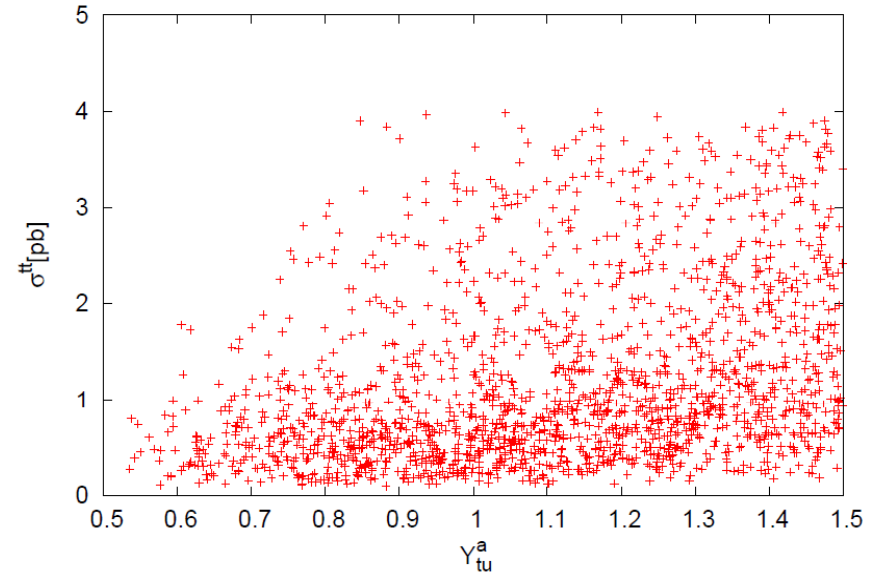
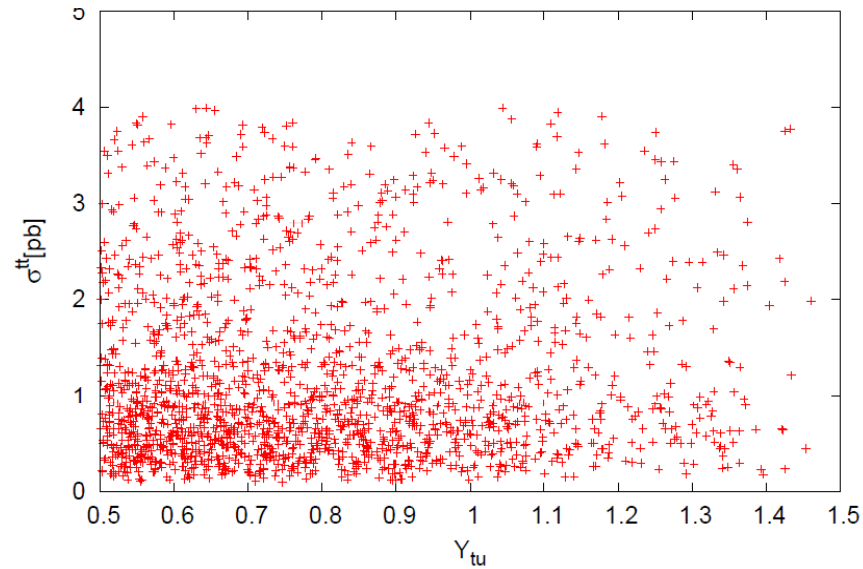
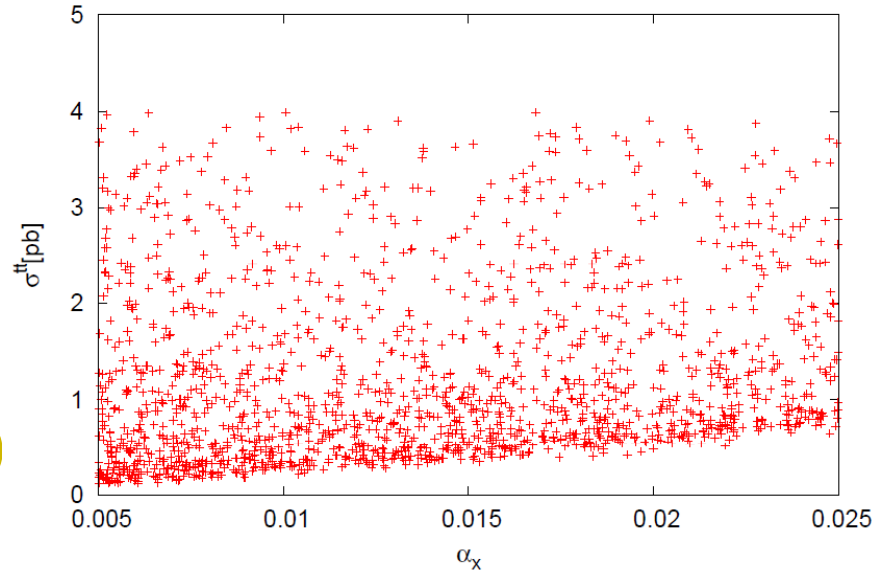
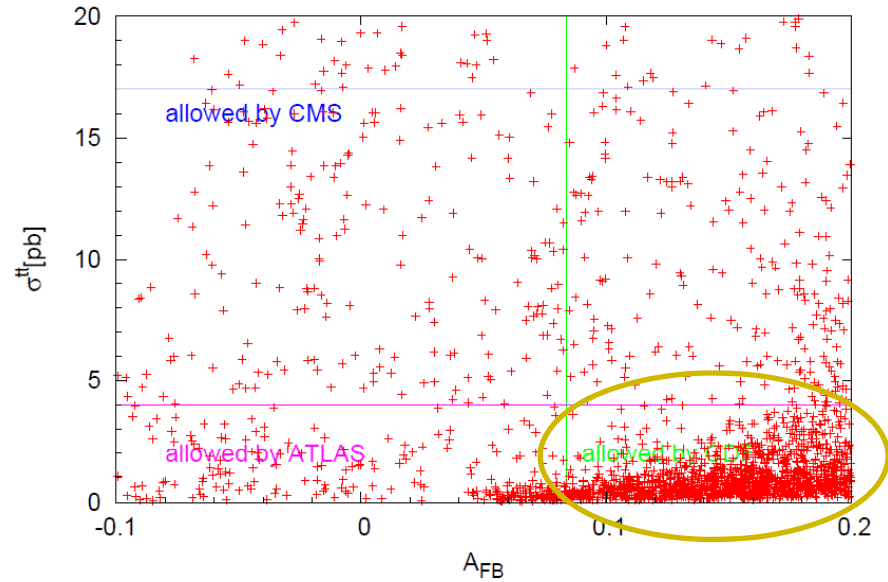
$140 \text{ GeV} < m_{Z'} < 1.5 \text{ TeV}$





# More general case

$140 \text{ GeV} < m_{Z'} < 1.5 \text{ TeV}$



# Conclusions

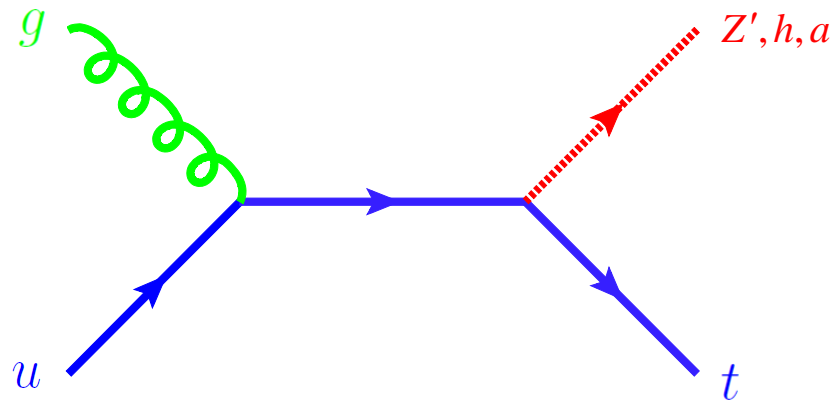
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- construct a complete  $U(1)'$  model where the right-handed up-type quarks in the standard model are charged.
- require extra Higgs charged under  $U(1)'$  for a realistic model.
- requires extra chiral fermions for anomaly cancellation  $\rightarrow$  CDM.
- Interferences between  $Z'$ ,  $h$ , and  $a$  reduce the rate for the same sign top pair production.
- The CDF  $W+jj$  excess may be resolved by the  $WZ'$  production through a charged Higgs mediation.

# Backup slides

# Single top quark production

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- **D0** [D0, 1105.2788](#)

$$\sigma(p\bar{p} \rightarrow tbq) = 2.90 \pm 0.59 \text{ pb}$$

In the SM,

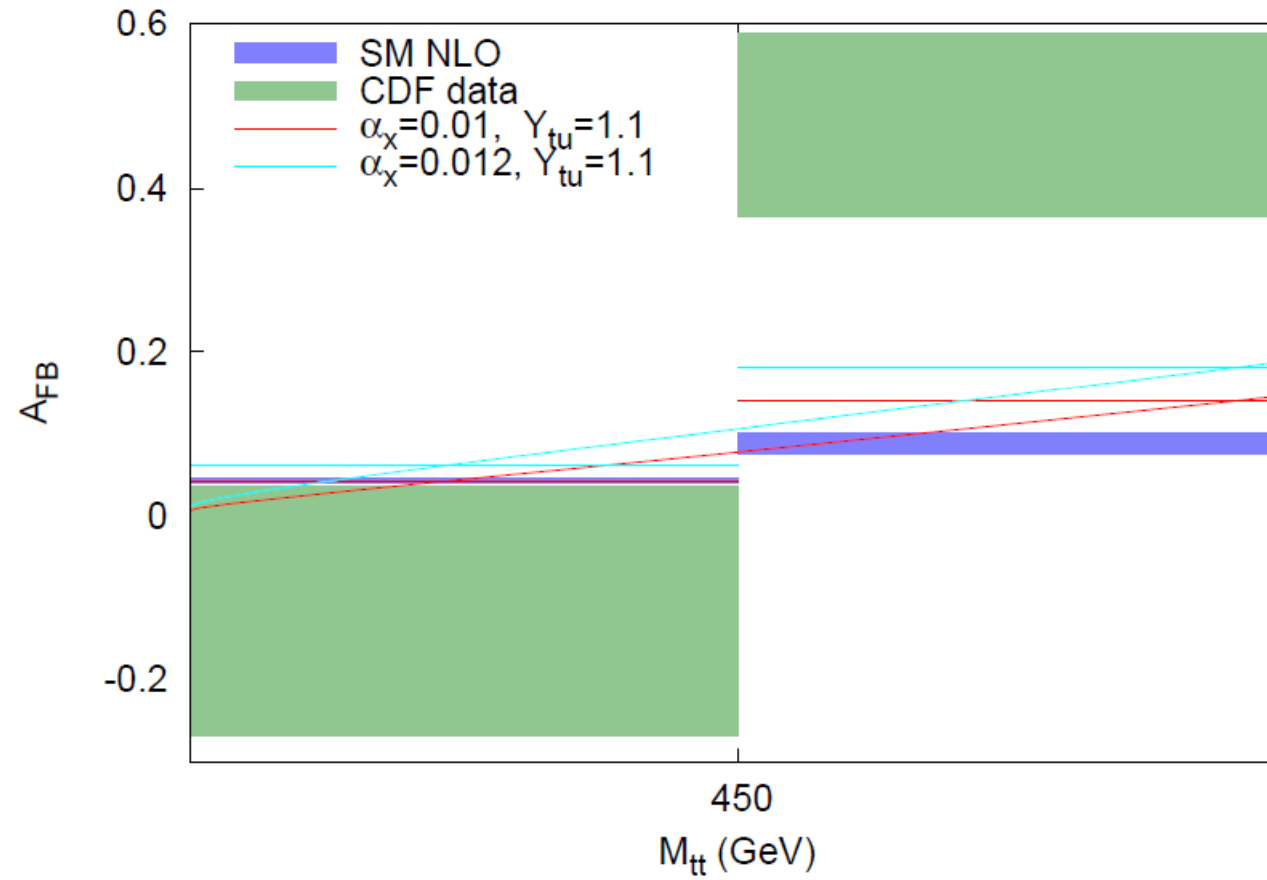
$$\sigma(p\bar{p} \rightarrow tbq) = 2.26 \pm 0.12 \text{ pb}$$

- **CMS** [CMS, 1106.3052](#)

$$\sigma(pp \rightarrow tbq) = 83.6 \pm 29.8 \pm 3.3 \text{ pb}$$

$$\sigma(pp \rightarrow tbq) = 64.3^{+2.1+1.5}_{-0.7-1.7} \text{ pb}$$

# Forward-backward asymmetry



# Flavor dependent U(1)' Model

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- Anomaly cancelation requires extra fermions II: SU(3)<sub>c</sub> triplets

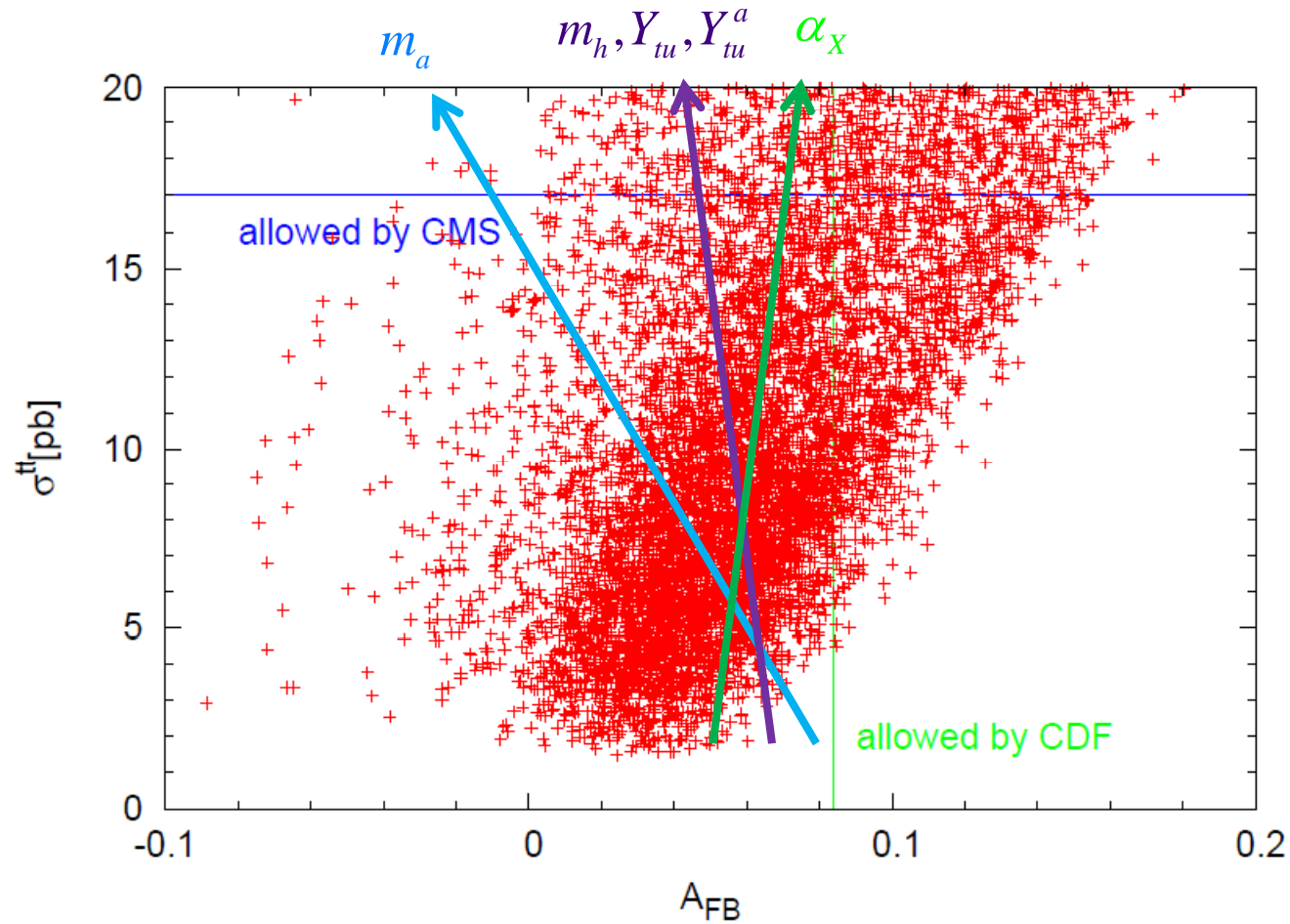
	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)'$
$q_{L1}$	3	1	-1/3	$Q_L$
$q_{R1}$	3	1	-1/3	$Q_R$
$q_{L2}$	3	1	-1/3	$-Q_L$
$q_{R2}$	3	1	-1/3	$-Q_R$

- introduce the singlet scalar X to the SM in order to allow the decay of the extra colored particles.

$$V_m = \lambda_i X^\dagger \overline{D_{Ri} q_{L1}} + \lambda_i \overline{X} D_{Ri} q_{L2}$$

a candidate for CDM

# $A_{FB}$ versus $\sigma_{tt}$



$$m_{Z'} = 145 \text{ GeV}$$

$$180 \text{ GeV} < m_h < 1 \text{ TeV}$$

$$180 \text{ GeV} < m_a < 1 \text{ TeV}$$

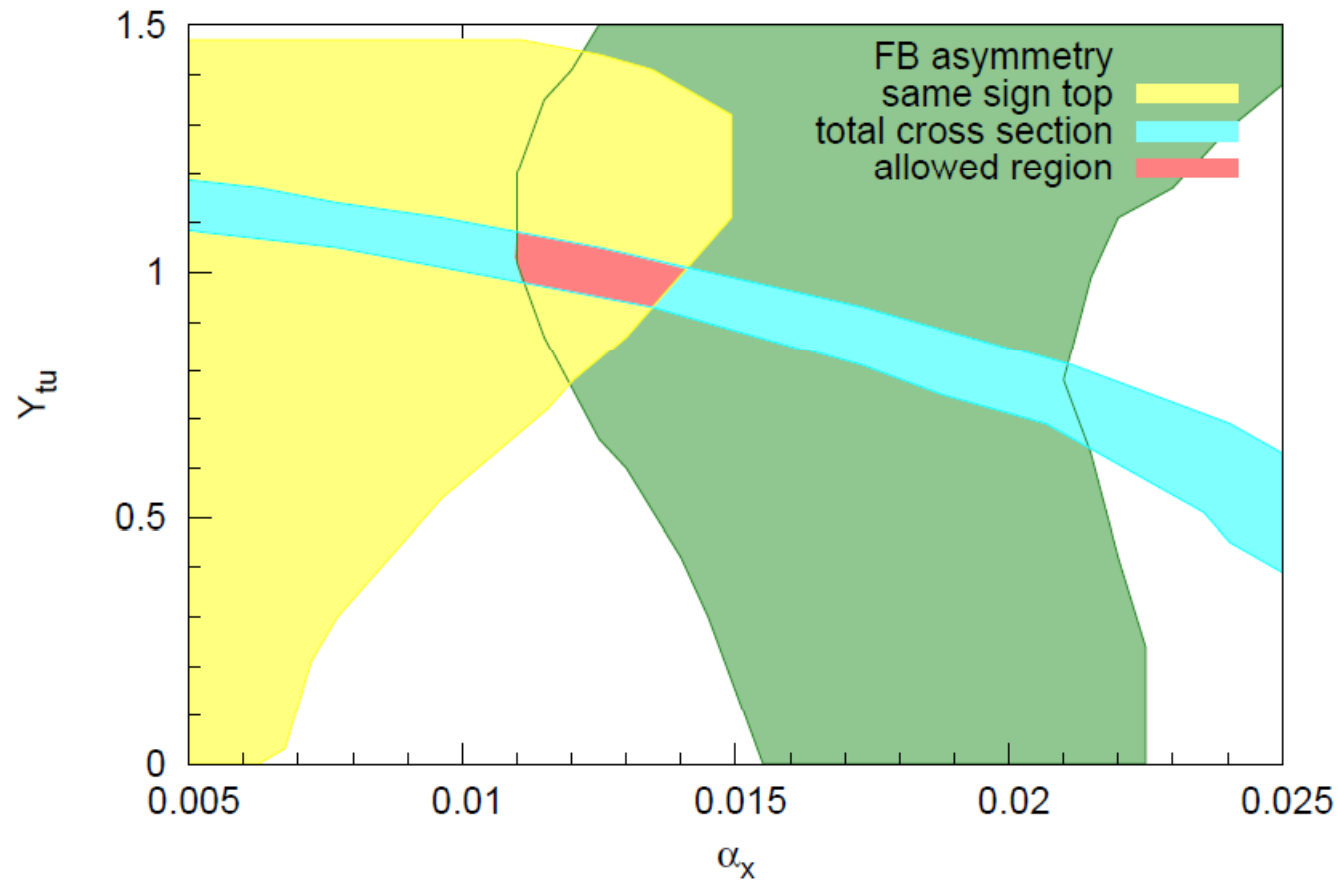
$$0.005 < \alpha_X < 0.025$$

$$0.5 < Y_{tu} < 1.5$$

$$0.5 < Y_{tu}^a < 1.5$$

# Favored region

Z'+h+a case



$m_{Z'} = 160 \text{ GeV}$

$m_h = 180 \text{ GeV}$

$m_a = 300 \text{ GeV}$

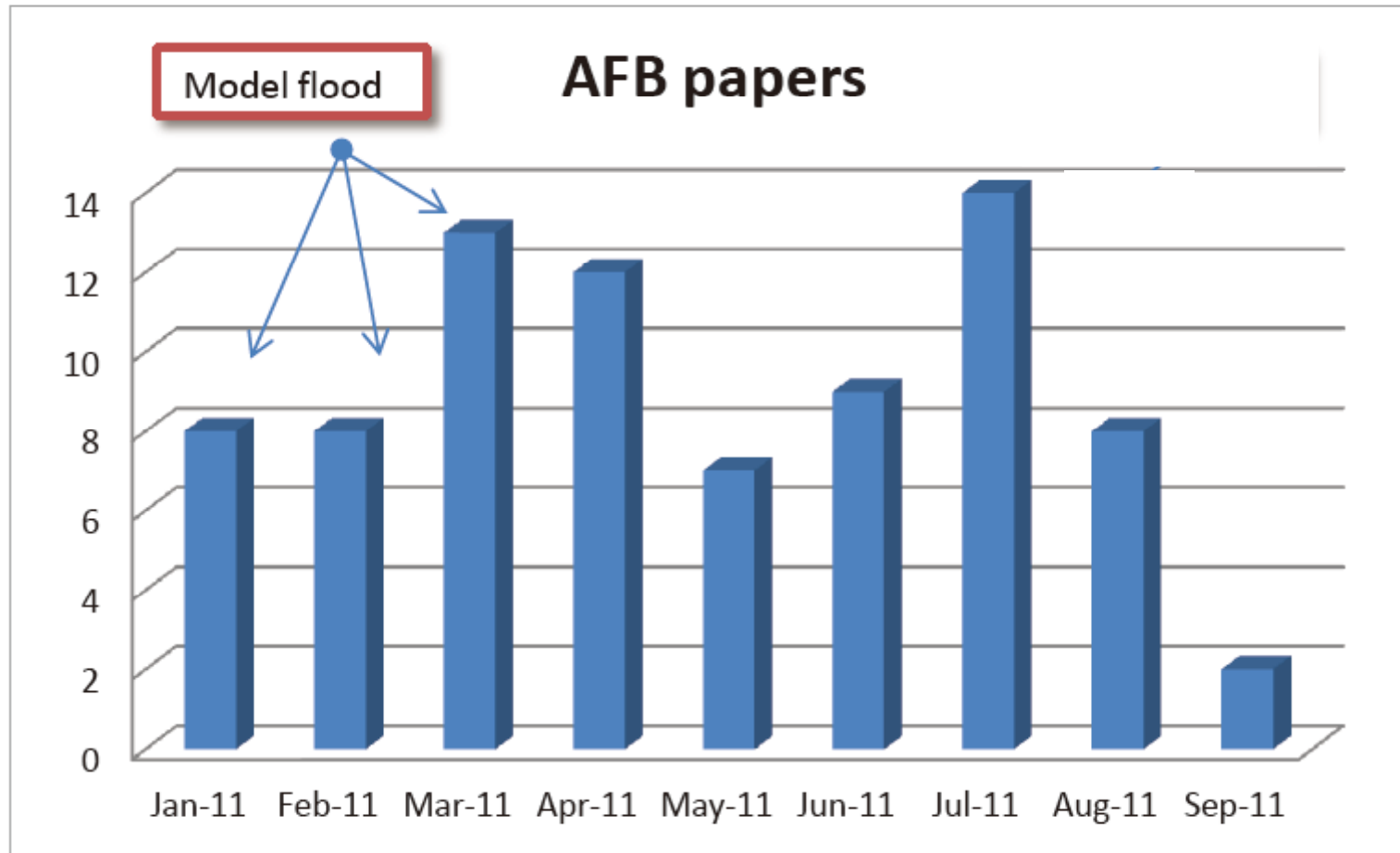
$Y_{tu}^a = 1.1$

$A_{FB} = 0.084 \sim 0.12$



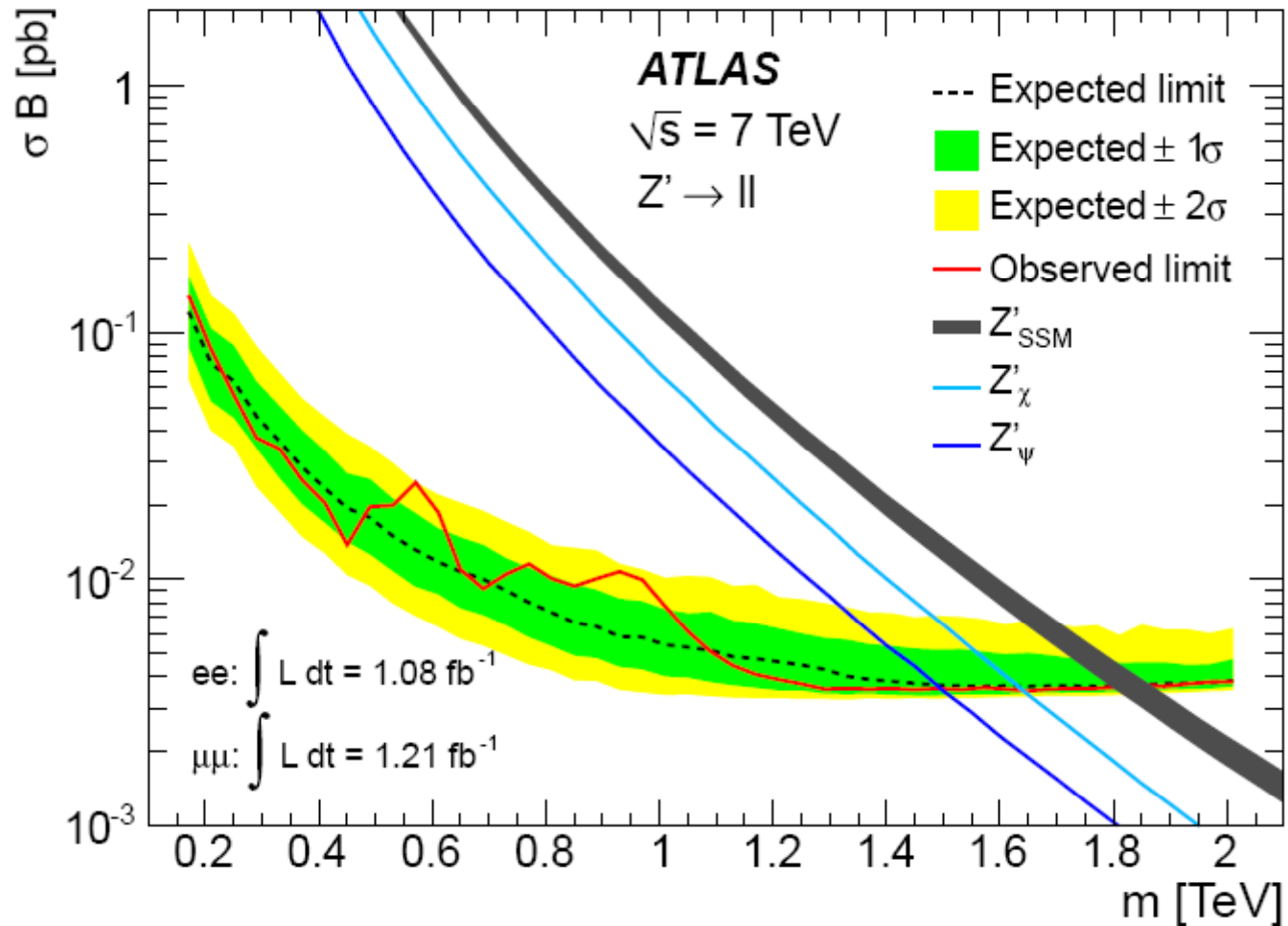
# New models

From Saavedra's talk, TOP2011



# Dilepton production at ATLAS

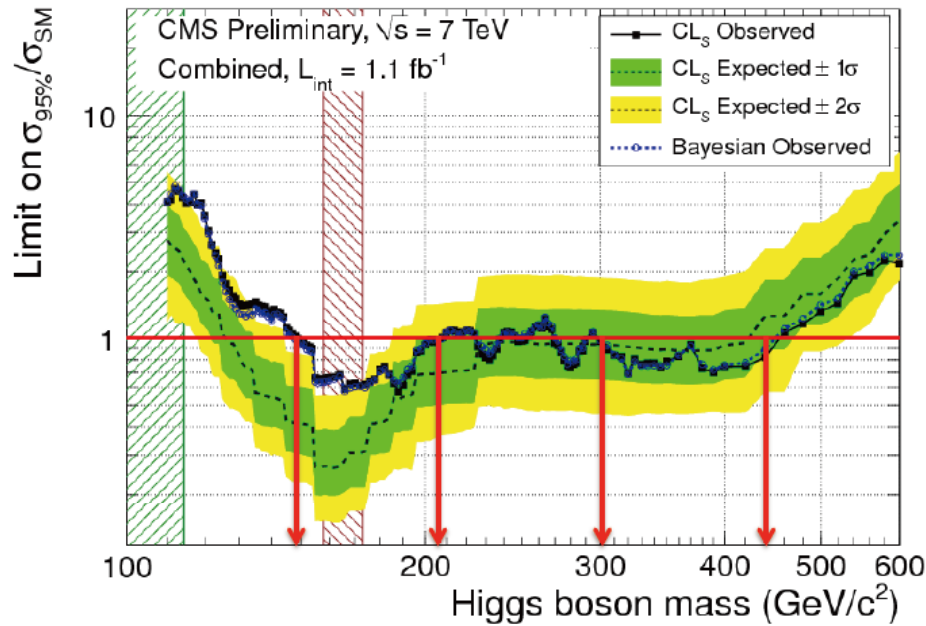
ATLAS, 1108.1582



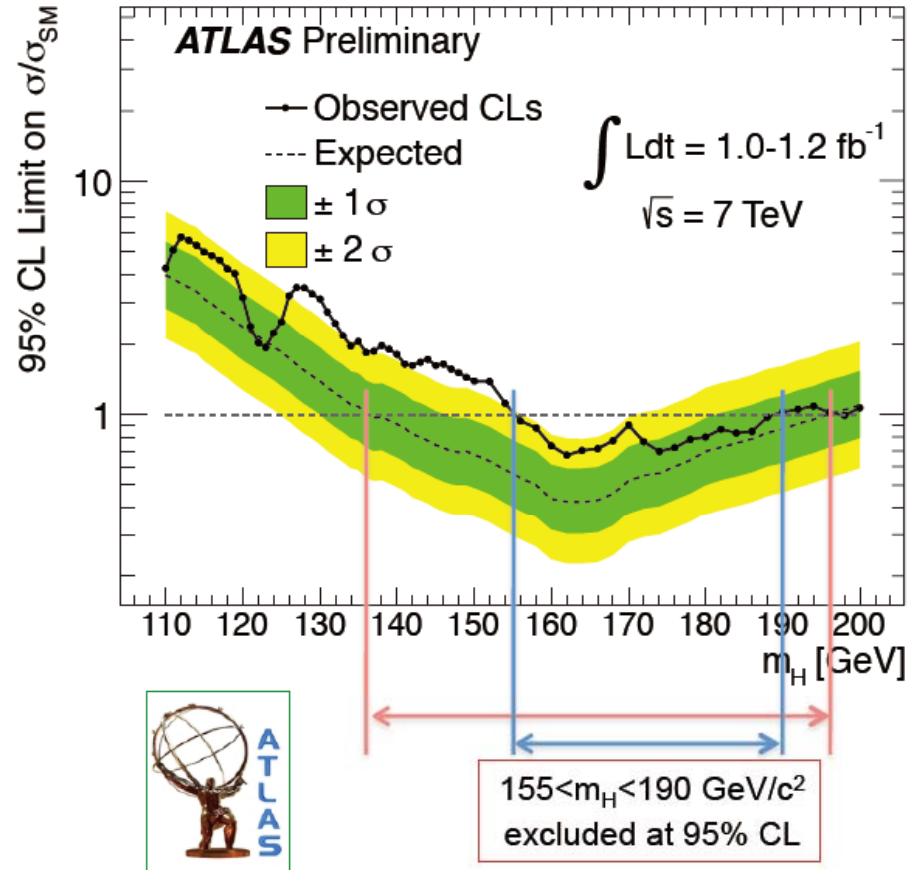
- excludes the Sequential SM  $Z'$  lighter than 1.83 TeV.

# Higgs search at LHC

From Korytov and Cranmer's talks, EPS-HEP 2011



**Excluded (GeV)**  
**[149-206] ... [300-440]**



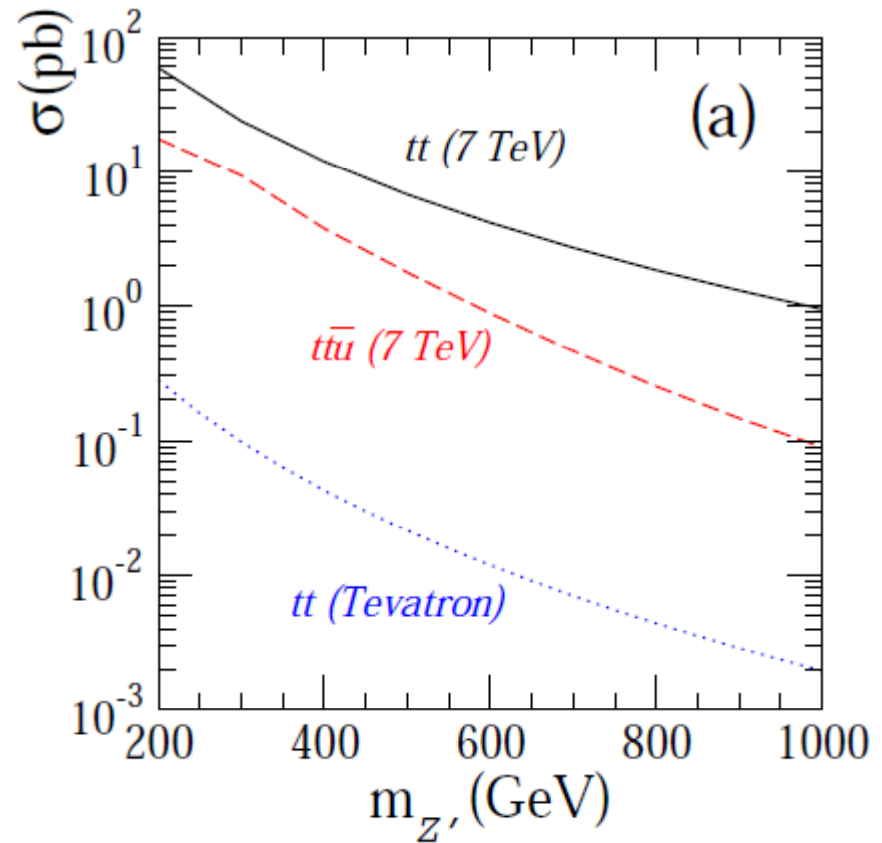
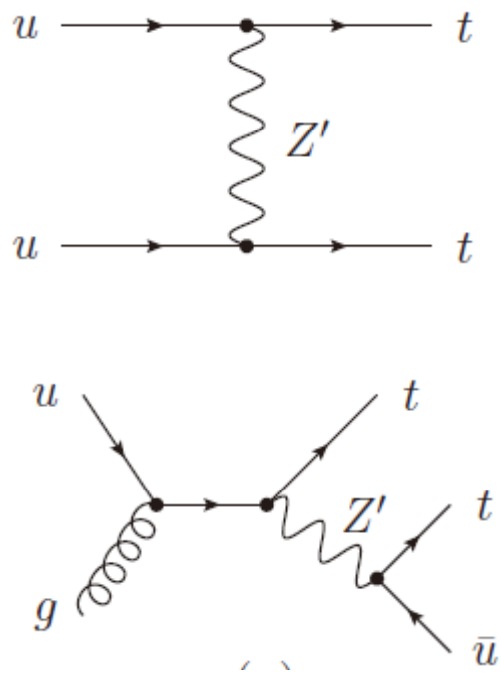
- $m_h = 180 \text{ GeV}$  : conflict with Higgs mass bounds at CMS and ATLAS?
- The bounds are weaker because new decay channels are open.

$$h \rightarrow t\bar{u}, h \rightarrow \Phi + \text{anything}$$

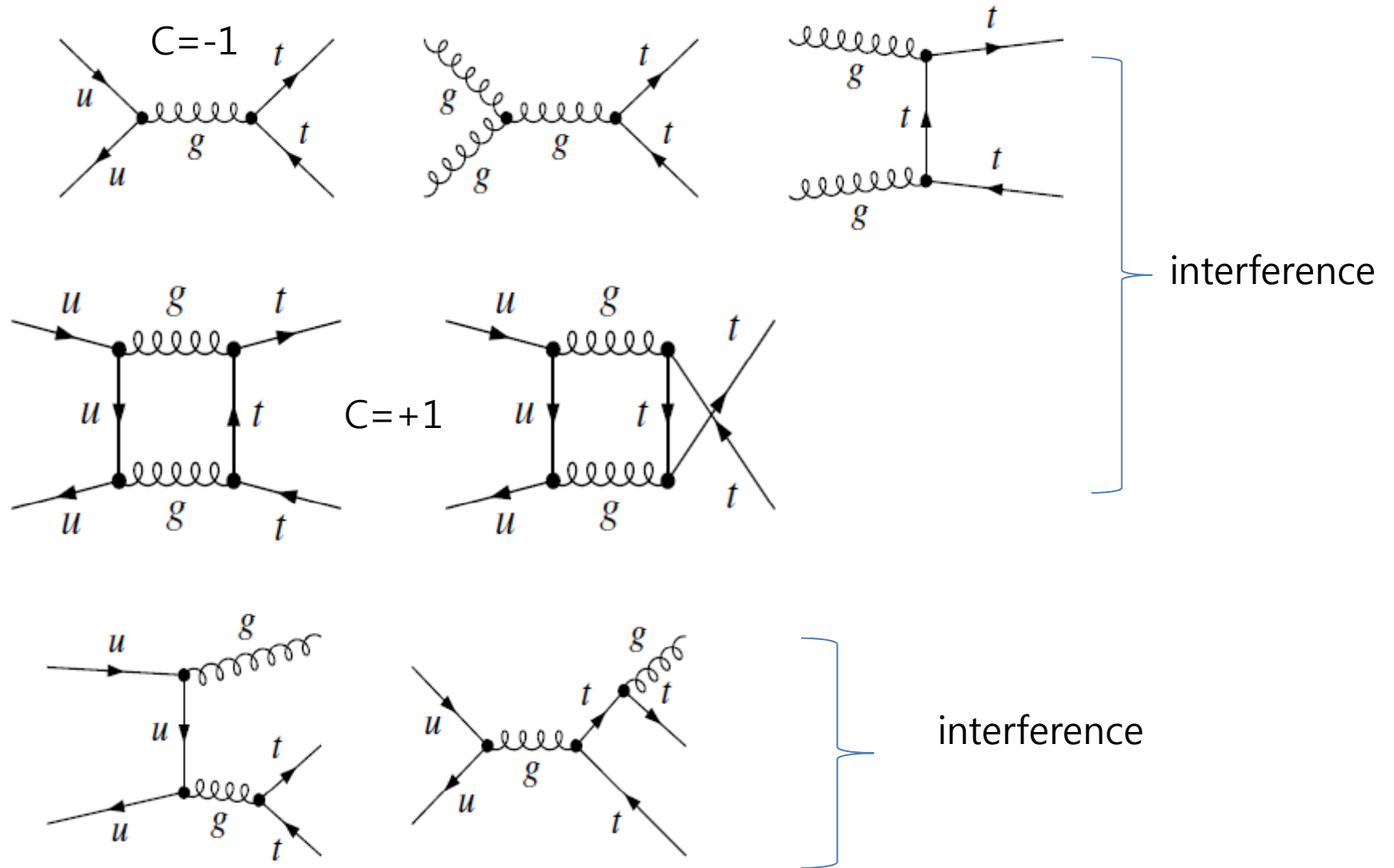
# Same sign top quark pair production at LHC

$$\mathcal{L} = g_W \bar{u} \gamma^\mu (f_L P_L + f_R P_R) t Z'_\mu + \text{h.c.},$$

Berger et al, 1101.5626

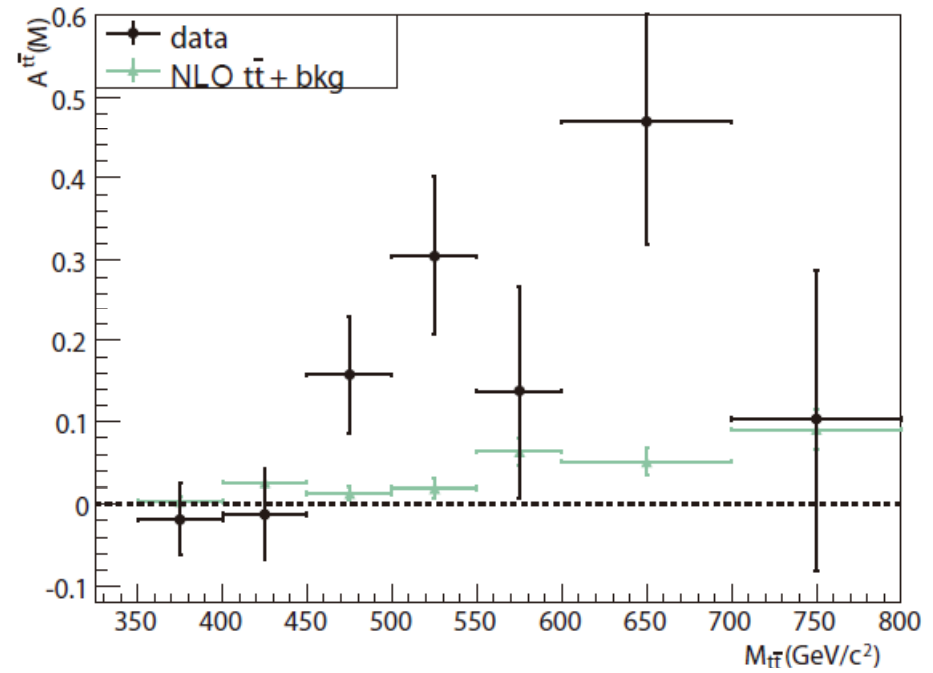
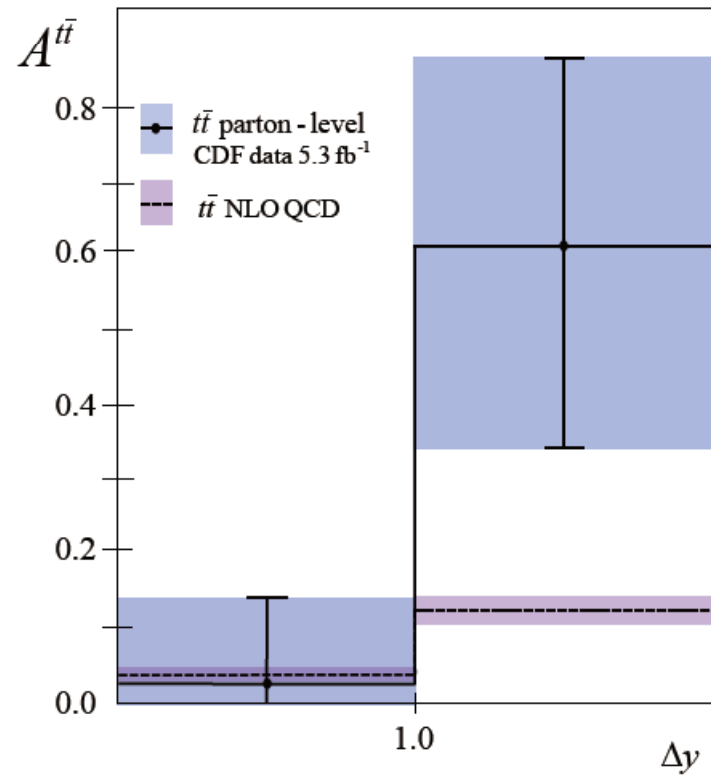


# Top $A_{FB}$ at Tevatron



# Top $A_{FB}^{t\bar{t}}$ at Tevatron

CDF, 1101.0034

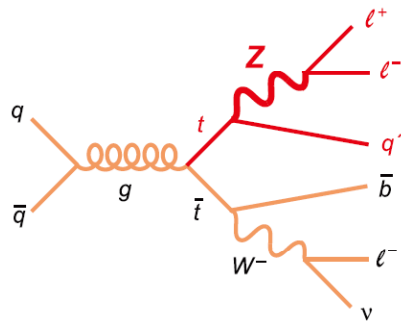


# Flavor-changing neutral currents

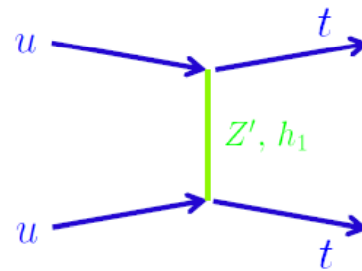
- FCNCs are suppressed in the SM (GIM mechanism).
- probe of new physics.
- FCNCs between down-type quarks and up-charm quarks.

$$K^0 - \bar{K}^0, B^0 - \bar{B}^0, B_s - \bar{B}_s, D^0 - \bar{D}^0 \text{ mixing.}$$

- top quark : no bound state.
- Which processes are proper for the test of the top FCNC?



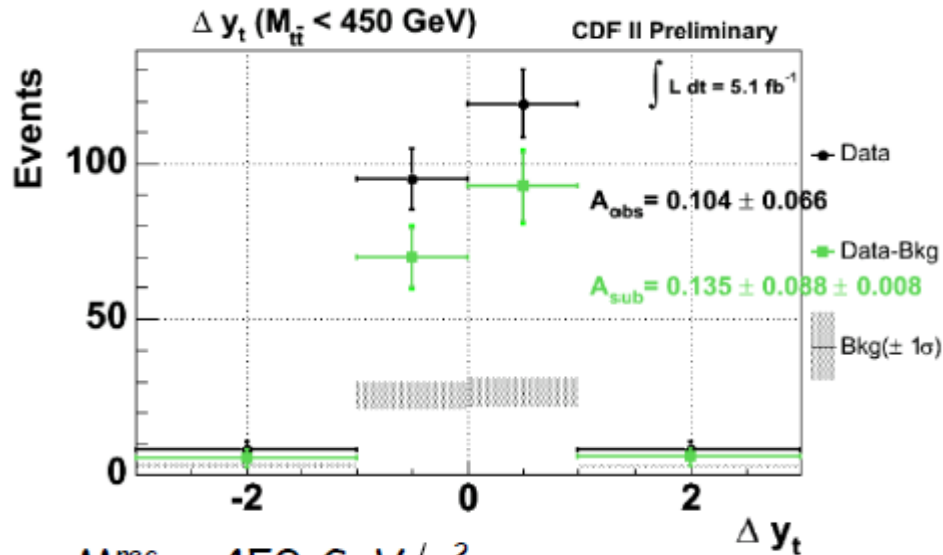
top decay



same sign top pair production

# Top $A_{FB}$ at CDF

From Hyunsoo Kim's talk, Tev2011



$$M_{t\bar{t}}^{\text{rec}} < 450 \text{ GeV}/c^2$$

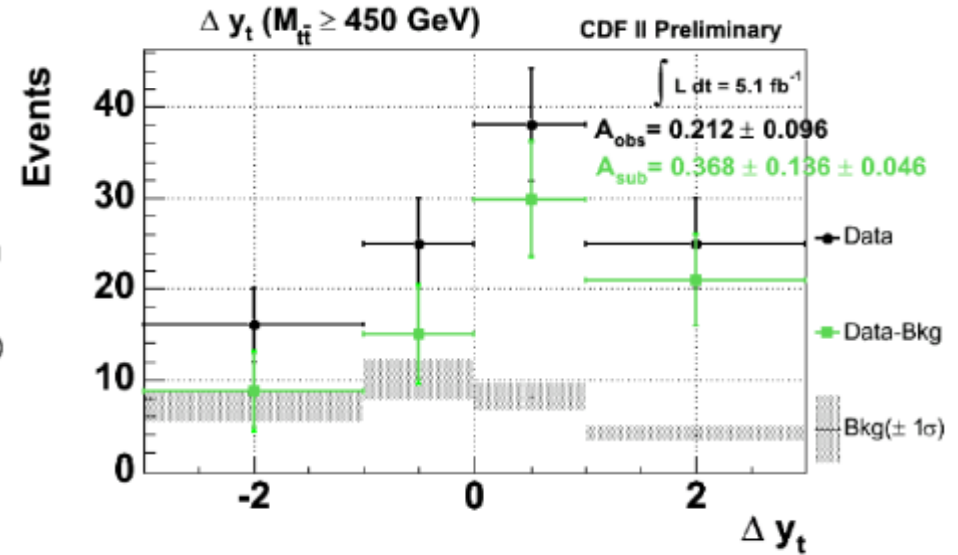
$$A_{\text{FB}}^{\text{sub}} = 0.135 \pm 0.088 (\text{stat.}) \pm 0.008 (\text{bkg shape})$$

$$A_{\text{FB}}^{\text{true}} = 0.417 \pm 0.253 (\text{stat.}) \pm 0.023 (\text{bkg shape}) \pm 0.017 (\text{MC stat.}) \pm ???$$

$$M_{t\bar{t}}^{\text{rec}} > 450 \text{ GeV}/c^2$$

$$A_{\text{FB}}^{\text{sub}} = 0.368 \pm 0.136 (\text{stat.}) \pm 0.046 (\text{bkg shape})$$

$$A_{\text{FB}}^{\text{true}} = 0.411 \pm 0.162 (\text{stat.}) \pm 0.055 (\text{bkg shape}) \pm 0.010 (\text{MC stat.}) \pm ???$$





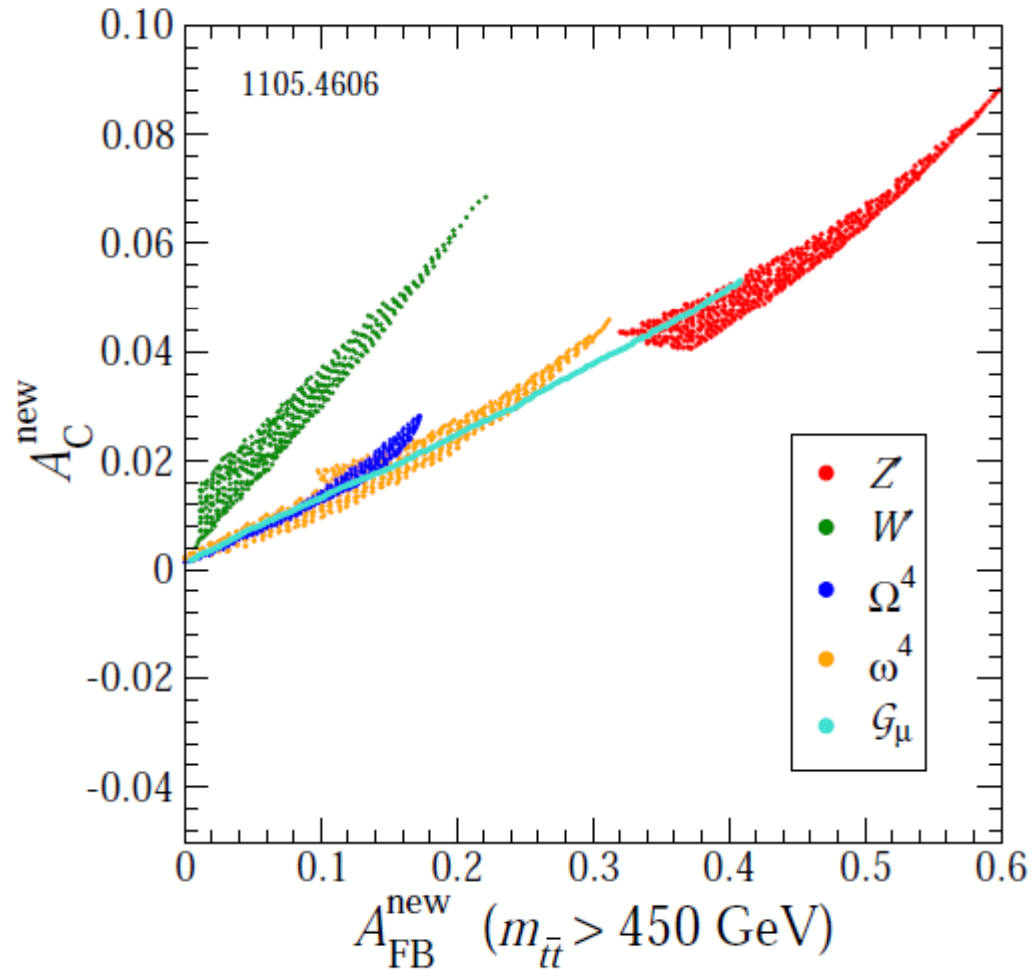
# Wjj excess at CDF

---

CDF, PRL104, 101801 (2011)

- $p\bar{p} \rightarrow W(\rightarrow \ell\nu) + jj$  with data of  $4.3 \text{ fb}^{-1}$ .
- diboson channel.
- background of the Higgs boson search.
- no significant excess for WW or WZ.
  - confirmed by D0.
- increase the jet  $E_T$  threshold from 20 GeV to 30 GeV.
  - interest in a higher mass range.

# Prediction for LHC charge asymmetry

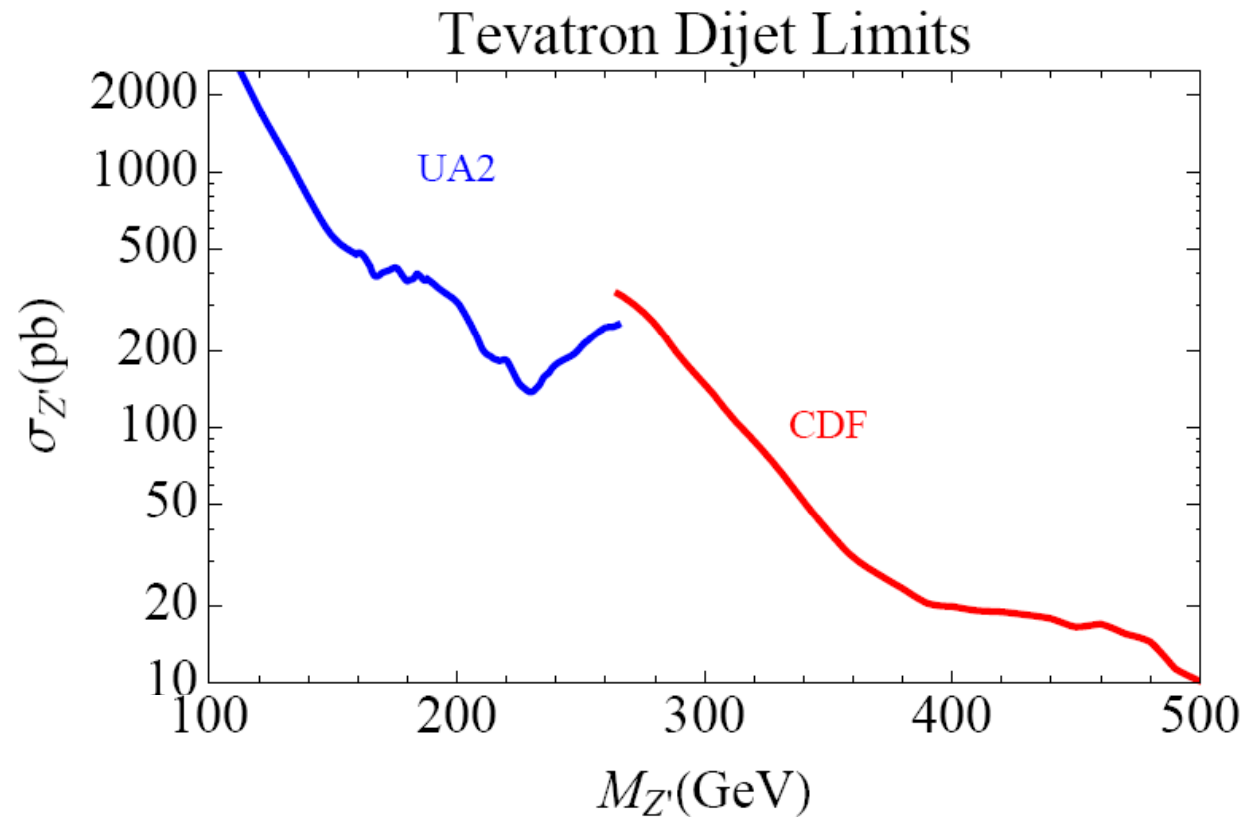


From Saavedra's talk, TOP2011

# Dijet production

Fan, Krohn, Langacker, Yavin, 1106.1682

- an extra resonance couples to the quarks : bound from  $p\bar{p} \rightarrow jj$ .



- assume  $m_{Z'} < 200$  GeV and  $m_{H^+} = 270$  GeV.