
Probing the fourth generation Majorana neutrino dark matter

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arXiv:1110.2930

outline

□ Introduction

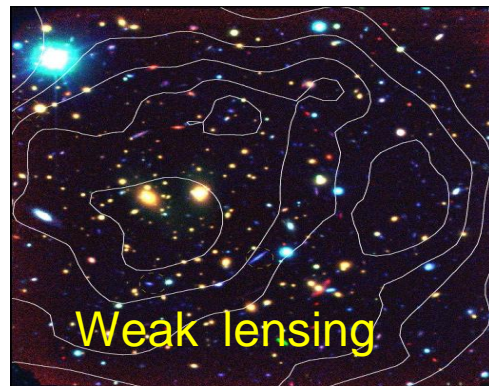
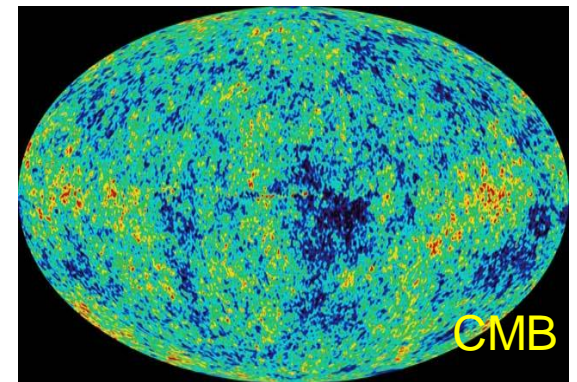
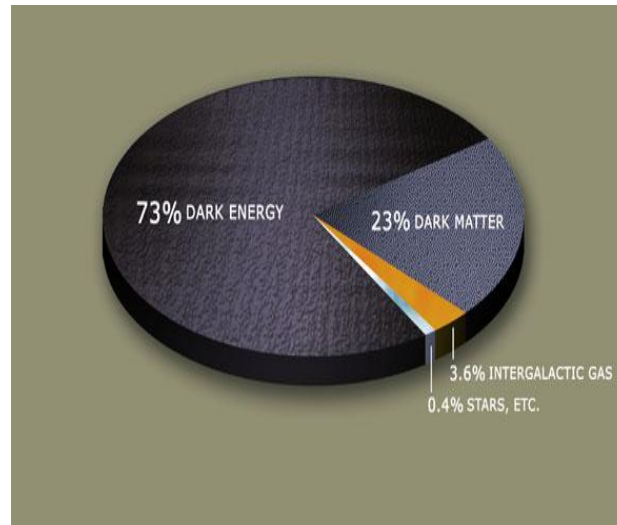
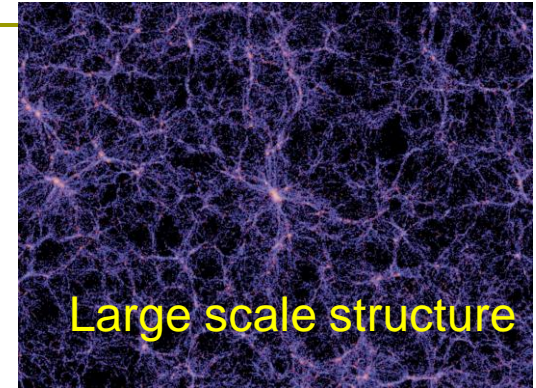
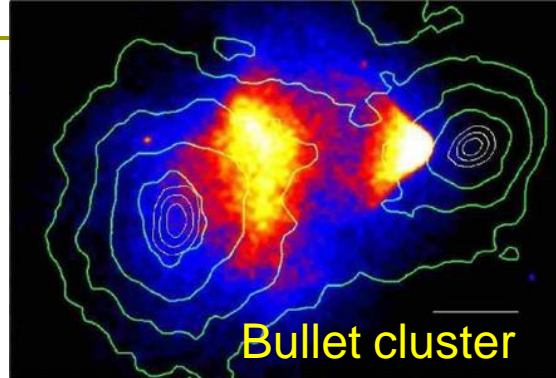
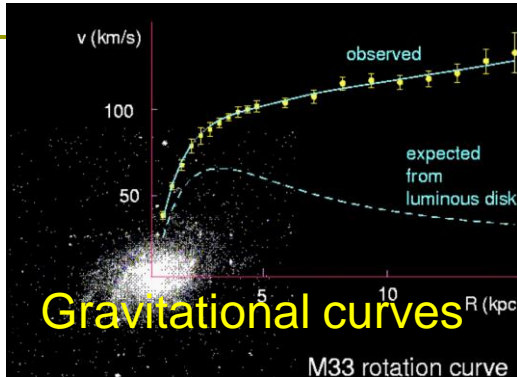
- Brief overview on current DM searches
- Neutrinos as DM candidates
- The 4th generation models

□ A 4th generation model with Majoran neutrino DM

- The 4th generation with $U(1)'_F$.
- Correlation between relic density and event rate in direct detection searches
- Numerical results

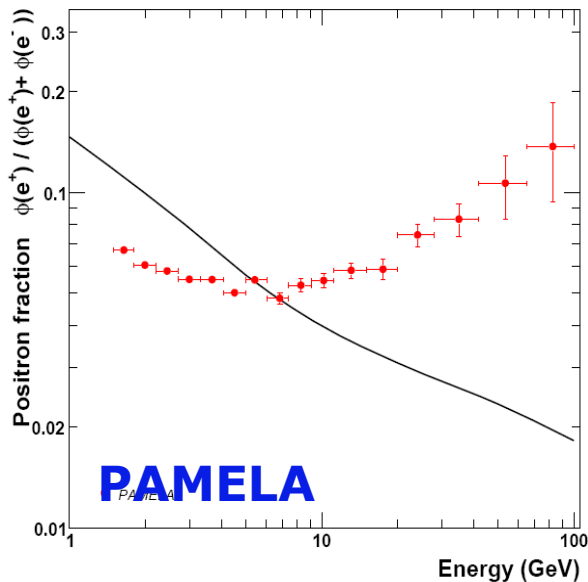
□ Conclusions

DM from gravitational effects

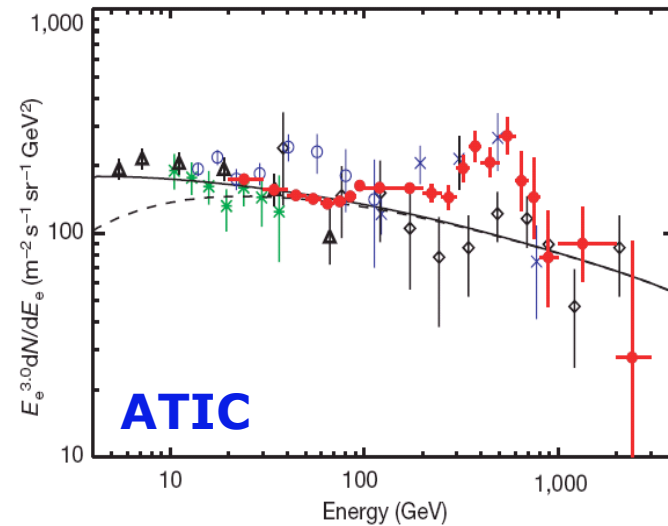
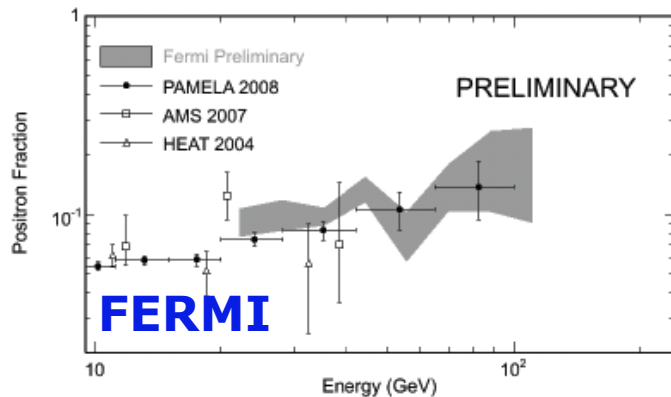


Hints of DM ?

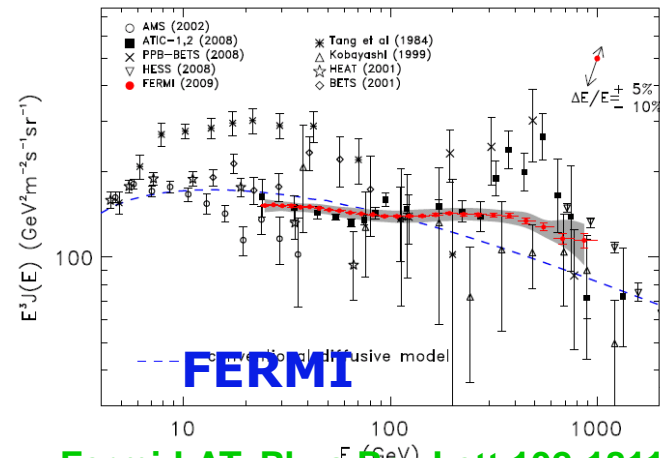
the cosmic-ray lepton anomalies



PAMELA, Nature 458, 607 (2009)



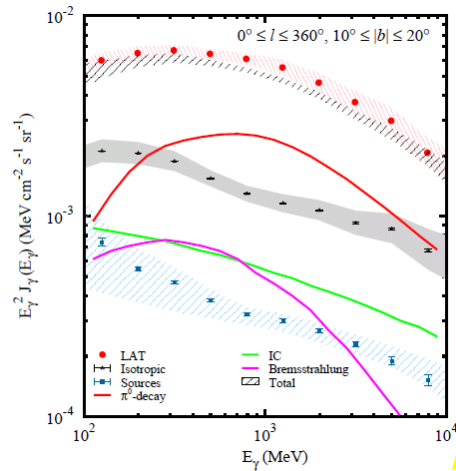
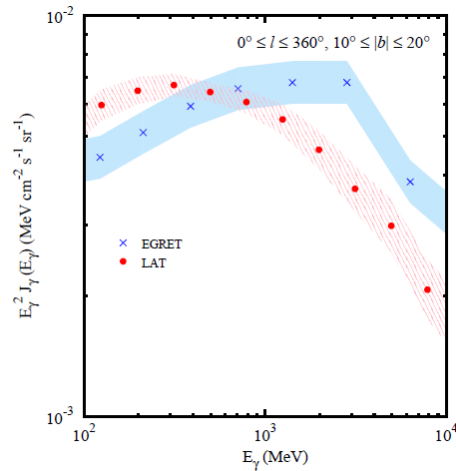
ATIC, Nature, 456, 2008,362-365



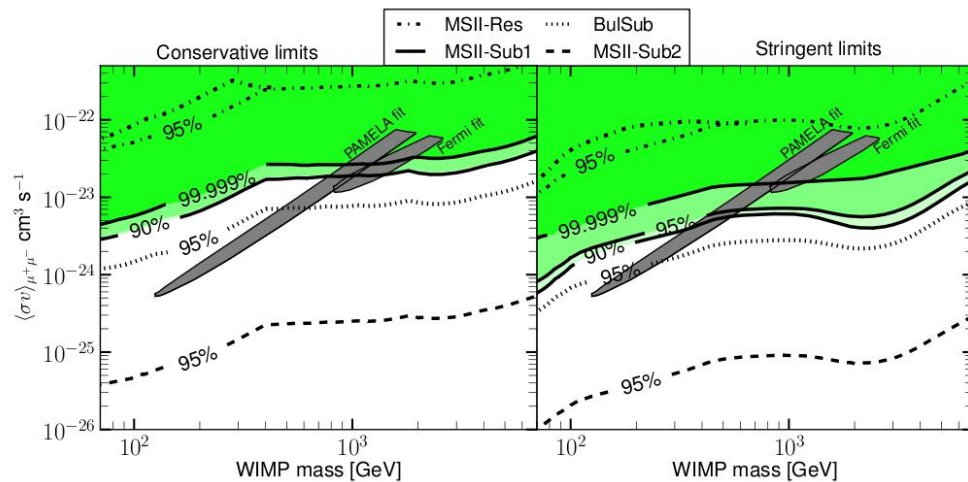
Fermi-LAT, Phys.Rev.Lett.102:181101,2009

Constraints from cosmic gamma-rays

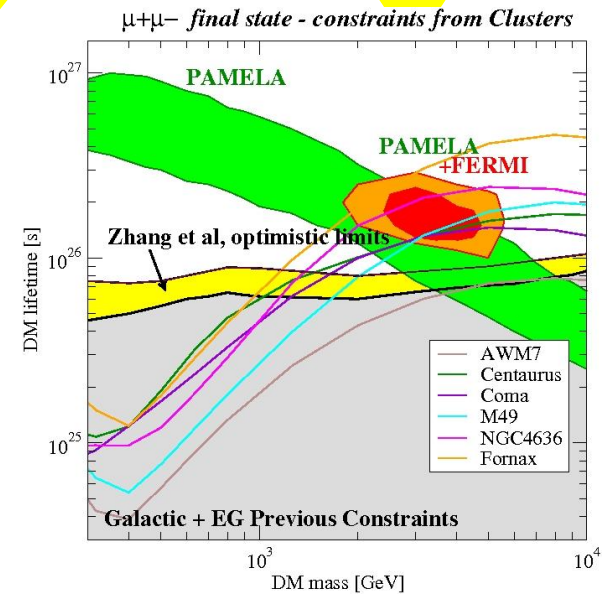
Fermi-LAT, Phys.Rev.Lett.103:251101, 2009



DM annihilation/decay constrained by the null results on cosmic gamma-rays



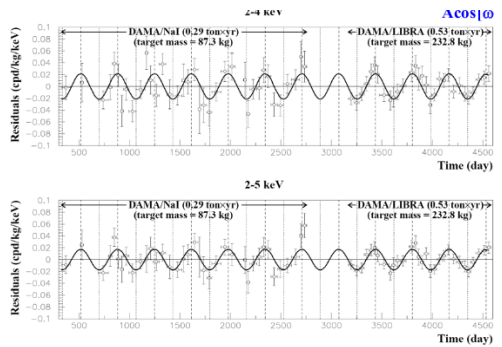
Abdo, et.al, arXiv:1002.4415



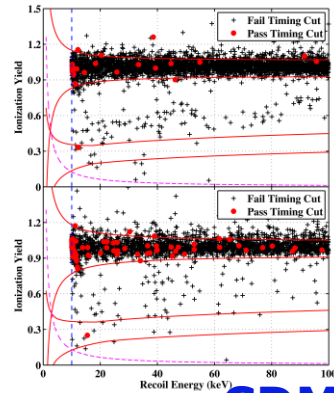
Dugger et.al, arXiv:1009.5988

Dark matter direct searches

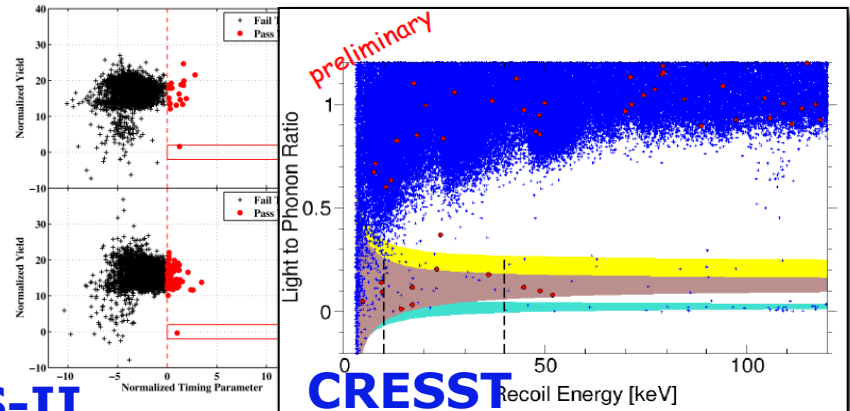
H.Kim talk



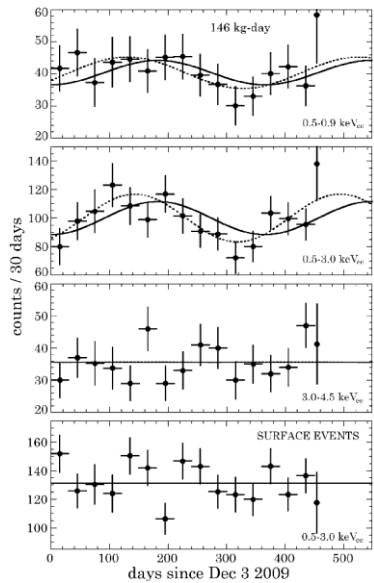
DAMA



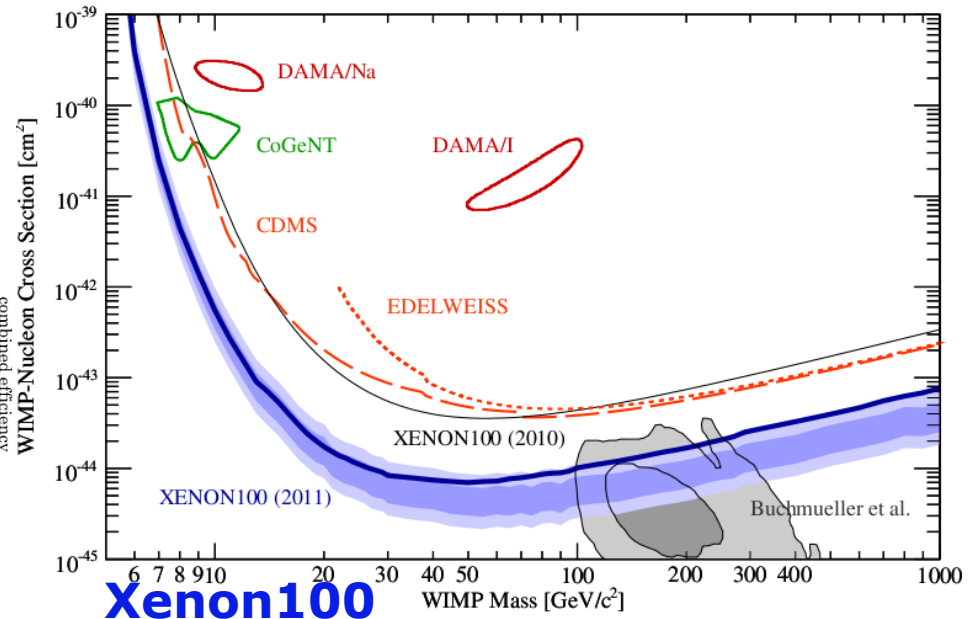
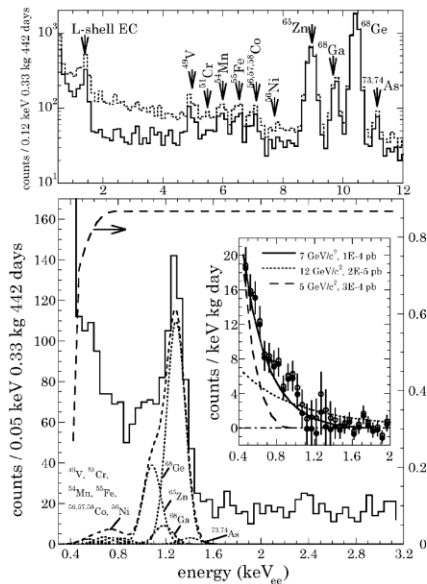
CDMS-II



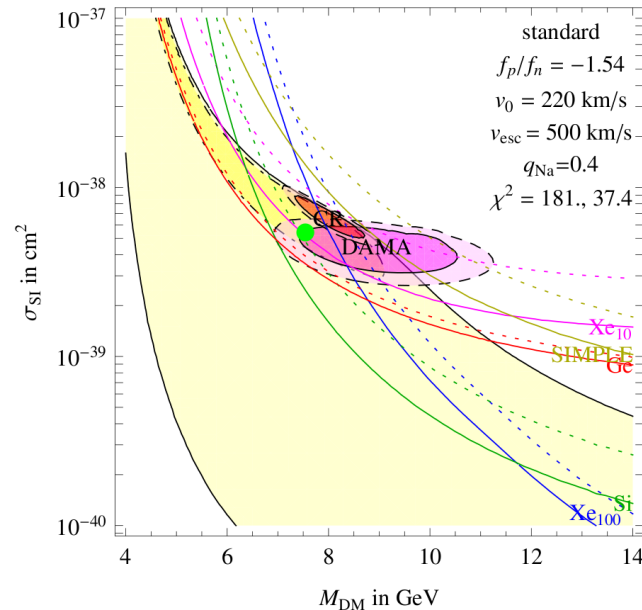
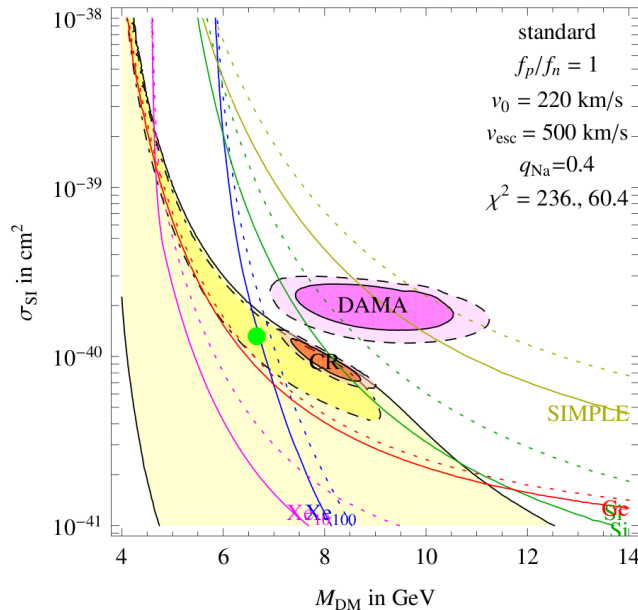
CRESST



CoGeNT



Xenon100



Explanations

- Inelastic scattering ?
- Isospin-violating DM ?
- Velocity suppressed interaction
- Momentum dependent scattering
- Resonant scattering

Uncertainties

- halo DM velocity distribution
- Quenching factors
- Channeling effects
-

Stability of DM

□ Using extra symmetries

motivated to evade phenomenological constraints

- R-Parity: MSSM, LSP
- KK-Parity: UED, LKP
- T-Parity: Little higgs,

motivated by DM

- Z_2 : SM+scalar DM
- $U(1)'$: SM+fermion DM (**E.J.Chun, Y. Omura, C.Yu, H.S. Lee,.. talk**)

□ Using the known symmetries

- $SU(2)$: Minimal DM model: SM + scalar/fermion with high representation, (**X.G. He talk**), $SU(2)$ “pion” (**R.Hill talk**)
- P, CP : LR symmetric model + scalar

Scalar DM protected by P/CP in LR models

Adding gauge-singlet in to the LR model

W.L.Guo, Y.L.Wu, YFZ,
Phys.Rev.D79:055015(2009)
Phys.Rev.D82:095004(2010)
Phys.Rev.D81:075014(2010)

$$\phi = \begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}, \chi = \begin{pmatrix} \chi_1^0 & \chi_1^+ \\ \chi_2^- & \chi_2^0 \end{pmatrix},$$

$$\Delta_{L,R} = \begin{pmatrix} \delta_{L,R}^+/\sqrt{2} & \delta_{L,R}^{++} \\ \delta_{L,R}^0 & -\delta_{L,R}^+/\sqrt{2} \end{pmatrix},$$

$$S = \frac{1}{\sqrt{2}}(S_\sigma + iS_D)$$

spontaneous CP violation assumed

P- and CP-transformations

	P	CP
$\phi \rightarrow$	$\phi^\dagger \rightarrow$	ϕ^*
$\chi \rightarrow$	$\chi^\dagger \rightarrow$	χ^*
$\Delta_{L(R)} \rightarrow$	$\Delta_{R(L)} \rightarrow$	$\Delta_{L(R)}^*$
$S \rightarrow$	$S \rightarrow$	S^*

Terms forbidden

$$(S - S^*)^{1,3}$$

$$(S - S^*)\text{Tr}(\phi^\dagger \phi)$$

$$(S - S^*)\text{Tr}(\Delta_L^\dagger \Delta_L + \Delta_R^\dagger \Delta_R)$$

Residual Z_2

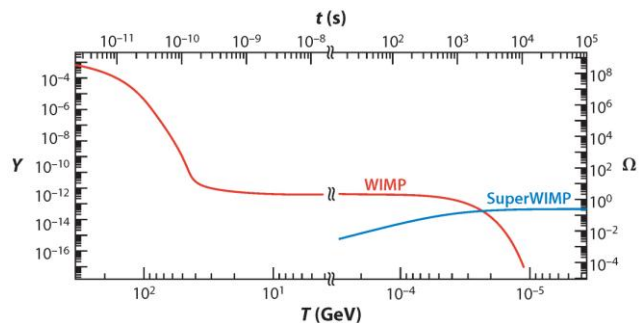
S_D a DM candidate

The relic density of DM

- Thermal: decouple from thermal equilibrium
- Nonthermal: never reached thermal equilibrium (super weak)

Nonthermal generation of DM

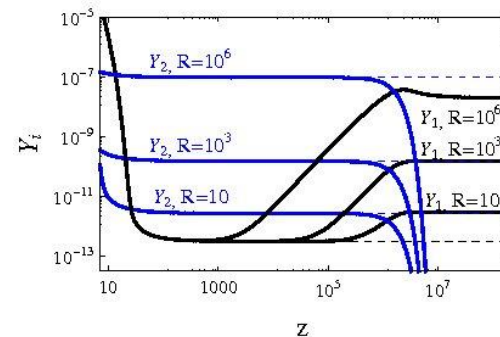
- gravitational interaction
- decay of unstable particles



J.L.Feng 2004

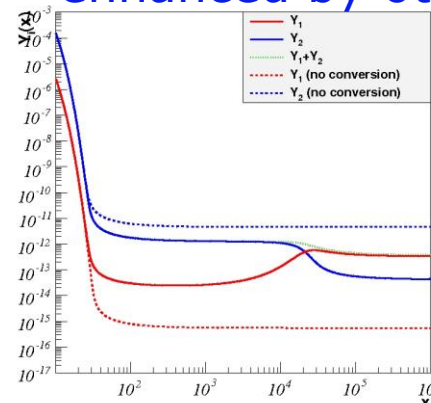
- Late decay may affect BBN, CMB
 - DM may get warm
- by transitions from other particles

- Thermal DM density enhanced by late decay of unstable states



Zupan, etal, 2009

- Thermal DM density enhanced by other DMs



Neutrinos as DM candidates

□ Very light SM neutrinos (hot DM)

CMB bound: $\Omega_\nu h^2 < 0.0067 (90\% CL)$

Disfavored by large scale structure formation

□ KeV sterile neutrinos (warm DM)

Less substructure, favored by observations

Life-time longer than the Univers

Non-thermal generation

Constraints: X-ray, BBN, Ly-alpha, ...

hard to detect directly ?

□ Heavy neutrinos (cold DM)

Heavy active Dirac/Majorana neutrino
cannot make up the whole DM

(not necessary in multi-DM models !)

Heavy sterile neutrino ?

GUT scale

DM ?

EW scale

cold DM?

KeV

warm DM

eV

hot DM



A 4th generation model with neutrino DM

- SM4: Simplest extension to the SM.

$$SM + \left(\begin{matrix} u_{4L} \\ d_{4L} \end{matrix} \right), \left(\begin{matrix} \nu_{4L} \\ e_{4L} \end{matrix} \right), u_{4R}, d_{4R}, (\nu_{4R}), e_{4R}$$

$$V_{CKM4} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} & V_{ub'} \\ V_{cd} & V_{cs} & V_{cb} & V_{cb'} \\ V_{td} & V_{ts} & V_{tb} & V_{tb'} \\ V_{t'd} & V_{t's} & V_{t'b} & V_{t'b'} \end{pmatrix}$$

- New sources of CP violation

6 real parameters, 3 phases in CKM4

Large CP violation possible

- Effective Hamiltonian

W.S.Hou, 0803.1234

$$H_{eff} = \frac{G_F}{\sqrt{2}} \left[\lambda_u (C_1^u O_1 + C_2^u O_2 + \sum C_i^u O_i) + \lambda_c \sum C_i^c O_i - \lambda_{t'} \sum \Delta C_i^{t'} O_i \right]$$

Can help in explaining the πK CP puzzle:

A.Soni, 0807.1871

$$\Delta A_{CP} = A_{CP}(\pi^0 K^-) - A_{CP}(\pi^+ K^-) = (14.4 \pm 2.9)\%$$

Constraints on the SM4

□ 4th generation quarks

LHC:

$m_{u_4} \geq 450 \text{ GeV}$ from searching for $u_4 \bar{u}_4 \rightarrow WbW\bar{b} \rightarrow b3j\ell^\pm E_T$,

$m_{d_4} \geq 490 \text{ GeV}$ from searching for $d_4 \bar{d}_4 \rightarrow WtW\bar{t} \rightarrow \ell^\pm \ell^\pm b3jE_T$

**R. Contino,
A. Koryot**

H.Yokoda talk

Tevatron:

Assuming 4-3 transition and 100% Br

$m_{u_4} \geq 335 \text{ GeV}$ assuming mass splitting less than $m_W : u_4 \rightarrow Wq$

$m_{d_4} \geq 385 \text{ GeV}$, assuming $m_{u_4} > m_{d_4}$ from $d_4 \rightarrow Wt$

G. Kribs etal, 0706.3718

J.Erler,P.Langacker,1003.3211

□ SM higgs with 4th family

LHC ruled out:

$$120\text{GeV} \leq m_H \leq 600\text{GeV}$$

EW precision test:

Mass difference between u_4
and d_4 have to be small

Constraints on the SM4 leptons

□ 4th generation leptons

$$m_{e_4} \geq 100.8 \text{ GeV} \text{ from } e_4 \rightarrow \nu_4 W^\pm$$

□ 4th generation neutrinos

- Unstable neutrinos $\nu_4 \rightarrow (e, \mu, \tau) W^+$

$$m_{\nu_4} \geq (101.3, 101.5, 90.3) \text{ GeV} \text{ (unstable Dirac)}$$

$$m_{\nu_4} \geq (89.5, 90.7, 80.5) \text{ GeV} \text{ (unstable Majorana)}$$

- Stable neutrinos -> weaken EWPT bounds, dark matter ?

$$m_{\nu_4} \geq 45.0(39.5) \text{ GeV} \text{ (stable Dirac(Majorana))}$$

The 4th generation leptons must have quite different nature

A 4th generation model with stable neutrinos

- A 4th generation U(1)' model

$$q_{iL} = \begin{pmatrix} u_{iL} \\ d_{iL} \end{pmatrix}, \ell_{iL} = \begin{pmatrix} \nu_{iL} \\ e_{iL} \end{pmatrix}, u_{iR}, d_{iR}, \nu_{iR}, e_{iR} \quad (i = 1, \dots, 4), \phi_a, \phi_b$$

- Gauge interaction

$$SU(2) \otimes U(1)_Y \otimes U(1)_F$$

- Anomaly-free assignments

$U(1)_F$	charges	
u_i, d_i	Q_q	$u_4, d_4 \quad -3Q_q$
ν_i, e_i	Q_L	$\nu_4, e_4 \quad -3Q_L$
ϕ_a	$-2Q_L$	$\phi_b \quad +6Q_L$

1) vector-like interaction

$$[U(1)_F]^3, [SU(3)_C]^2, [gravity]^2 U(1)_F \rightarrow 0$$

2) due to $\sum (-Y_{qL} + Y_{qR}) = 0$ and $\sum (-Y_{\ell L} + Y_{\ell R}) = 0$

$$U(1)_Y [U(1)_F]^2 \rightarrow 0$$

3) from $\sum_{i=1}^4 Q_{qi} = 0$ and $\sum_{i=1}^4 Q_{Li} = 0$

$$[SU(2)_L]^2 U(1)_F \rightarrow 0$$

Anomalies canceled between generations

New interactions

□ Interactions

YFZ, arXiv:110.2930

$$\begin{aligned}
 \mathcal{L} = & \bar{f}_i i\gamma^\mu D_\mu f_i + (D_\mu \phi_a)^\dagger (D_\mu \phi_a) + (D_\mu \phi_b)^\dagger (D_\mu \phi_b) \\
 & - Y_{ij}^d \bar{q}_{iL} H d_{iR} - Y_{ij}^u \bar{q}_{iL} \tilde{H} u_{iR} - Y_{ij}^e \bar{\ell}_{iL} H e_{iR} - Y_{ij}^v \bar{\ell}_{iL} \tilde{H} e_{iR} \\
 & - \frac{1}{2} Y_{ij}^m \overline{v_{iR}^c} \phi_a v_{jR} \quad (i, j = 1, 2, 3) - \frac{1}{2} Y_4^m \overline{v_{4R}^c} \phi_b v_{4R} - V(\phi_a, \phi_b, H) + \text{H.c.}
 \end{aligned}$$

The 4x4 CKM matrix is block diagonal

-> no mixing at tree level

□ Mass of Z' from $U(1)'$ breaking

$$\langle \phi_a \rangle = v_{a,b} / \sqrt{2}$$

$$m_{Z'}^2 = g_F^2 (Q_a^2 v_a^2 + Q_b^2 v_b^2),$$

small g_F suppresses $Z' \rightarrow \ell \bar{\ell}$,

can evade severe LHC bounds

muon $g-2$

$$\Delta a_\mu \approx \frac{g_F^2}{12\pi^2} \frac{m_\mu^2}{m_{Z'}^2} \leq 3.9 \times 10^{-9}$$

lead to a bound

$$\sqrt{Q_a^2 v_a^2 + Q_b^2 v_b^2} \geq 1.4 \times 10^2 \text{ GeV}$$

Dark matter in the model

□ Possible dark matter

- Lightest active neutrino ($\sim eV$)

$$V_{1L}, (V_{2L}, V_{3L})$$

- Light sterile neutrinos ($\sim KeV$)

$$V_{1R}, V_{2R}, V_{3R}$$

- Lightest 4th neutrino ($\sim 100 GeV$)

$$V_{4L}, V_{4R}$$

- Neutral 4th bound-states ?

□ The 4th heavy Majorana neutrino as DM component

- allow

$$r_{\Omega} \equiv \frac{\Omega_{\chi_1}}{\Omega_{DM}} < 1,$$

- assumption on halo DM

$$r_{\rho} \equiv \frac{\rho_1}{\rho_0} \approx \frac{\Omega_{\chi_1}}{\Omega_{DM}},$$

A multi-component DM model?

$$\sum \Omega_i = \Omega_{DM} = 0.11$$

- nontrivial correlation between relic density and the event-rate of direct detection.
- very predictive, constrained by the current exp., clear prediction for SI scattering

Neutrino masses and interaction

□ 4th Neutrino mass matrix

$$m_\nu = \begin{pmatrix} 0 & m_D \\ m_D & m_M \end{pmatrix}, \quad m_1 = \left(\frac{s}{c}\right)m_D, \quad \text{and} \quad m_2 = \left(\frac{c}{s}\right)m_D. \quad \tan 2\theta = \frac{2m_D}{m_M}.$$

□ Mass eigenstates (the lightest one is stable: U(1)' protected)

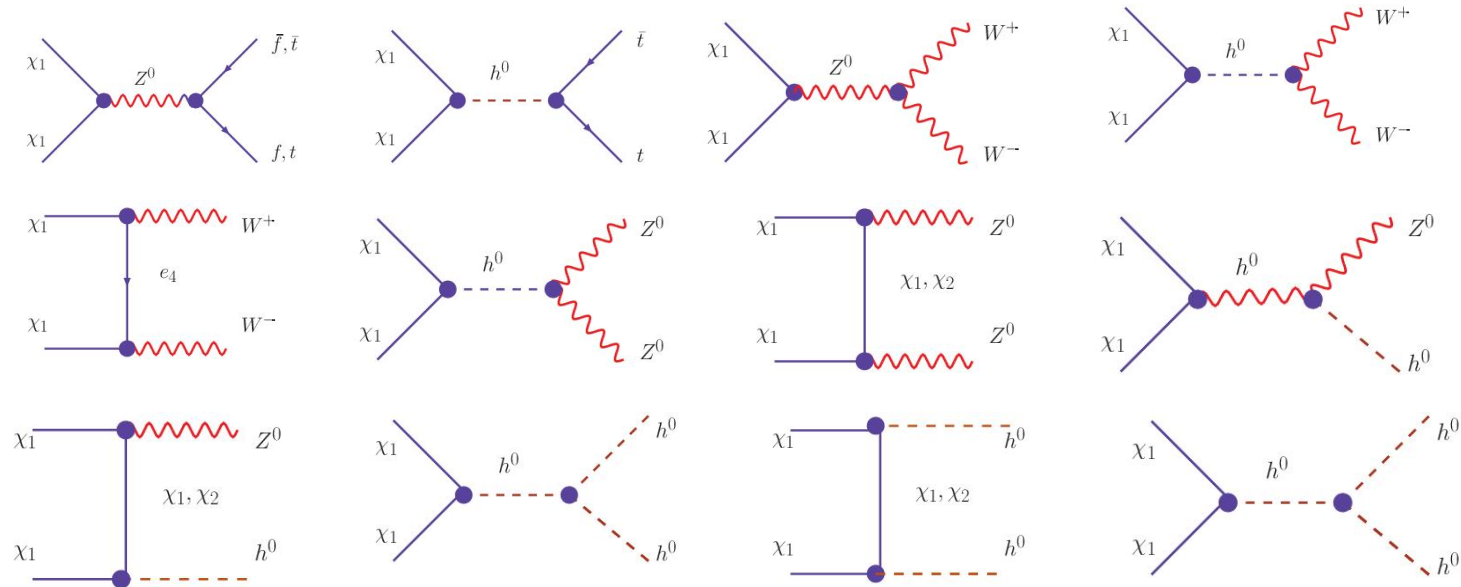
$$V_{1L}^{(m)} = -i(cV_L - sV_R^c), \quad V_{2L}^{(m)} = sV_L + cV_R^c,$$

□ Interactions with Z and SM higgs

$$\mathcal{L}_{NC} = \frac{g_1}{4 \cos \theta_W} \left[-c^2 \bar{\chi}_1 \gamma^\mu \gamma^5 \chi_1 - s^2 \bar{\chi}_2 \gamma^\mu \gamma^5 \chi_2 + 2ics \bar{\chi}_1 \gamma^\mu \chi_2 \right] Z_\mu,$$

$$\mathcal{L}_Y = -\frac{m_1}{v_H} \left(\frac{c}{s}\right) \left[cs \bar{\chi}_1 \chi_1 + cs \bar{\chi}_2 \chi_2 - i(c^2 - s^2) \bar{\chi}_1 \gamma^5 \chi_2 \right] h^0.$$

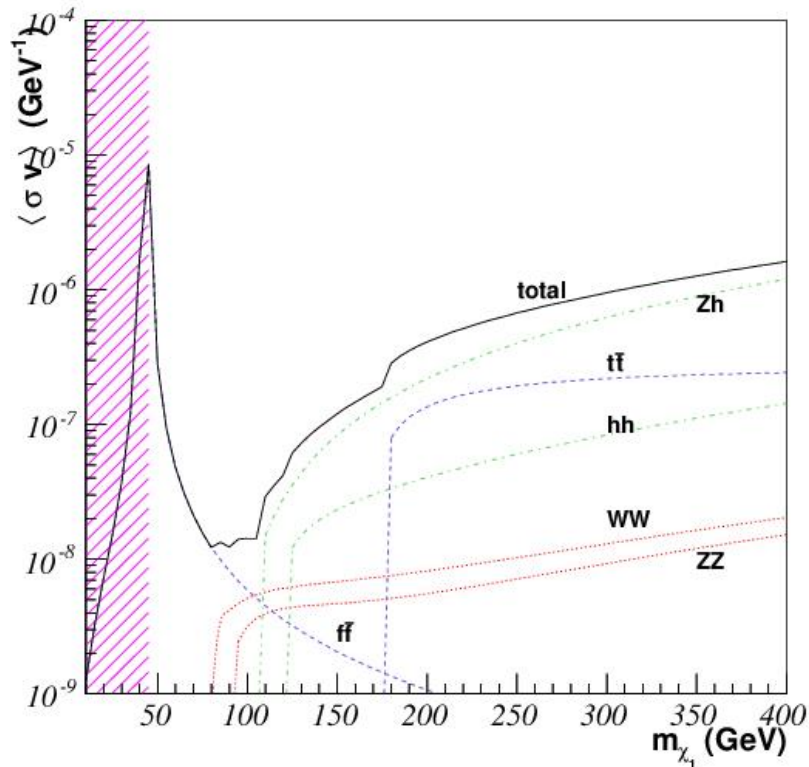
Annihilation channels



Thermally averaged cross section

$$\langle \sigma v \rangle = \frac{1}{8m_1^2 T K_2^2(m_1/T)} \int_{4m_1^2}^{\infty} ds \sigma(s - 4m_1^2) \sqrt{s} K_1\left(\frac{\sqrt{s}}{T}\right),$$

Annihilation cross-sections



Annihilation channels

1) For $m_\chi < m_Z / 2$, $f\bar{f}$ is dominant

$$\langle \sigma v \rangle \approx G_F^2 m_\chi^2 \langle v^2 \rangle (C_V^2 + C_A^2) c^4, \text{ suppressed}$$

2) For $m_\chi > m_W$, W^+W^- channel opens

$$\langle \sigma v \rangle \approx G_F^2 m_\chi^2 \langle v^2 \rangle \beta_W s^2 |D_Z|^2 c^4, \text{ suppressed}$$

3) For $m_\chi > m_Z$, Z^0Z^0 channel opens

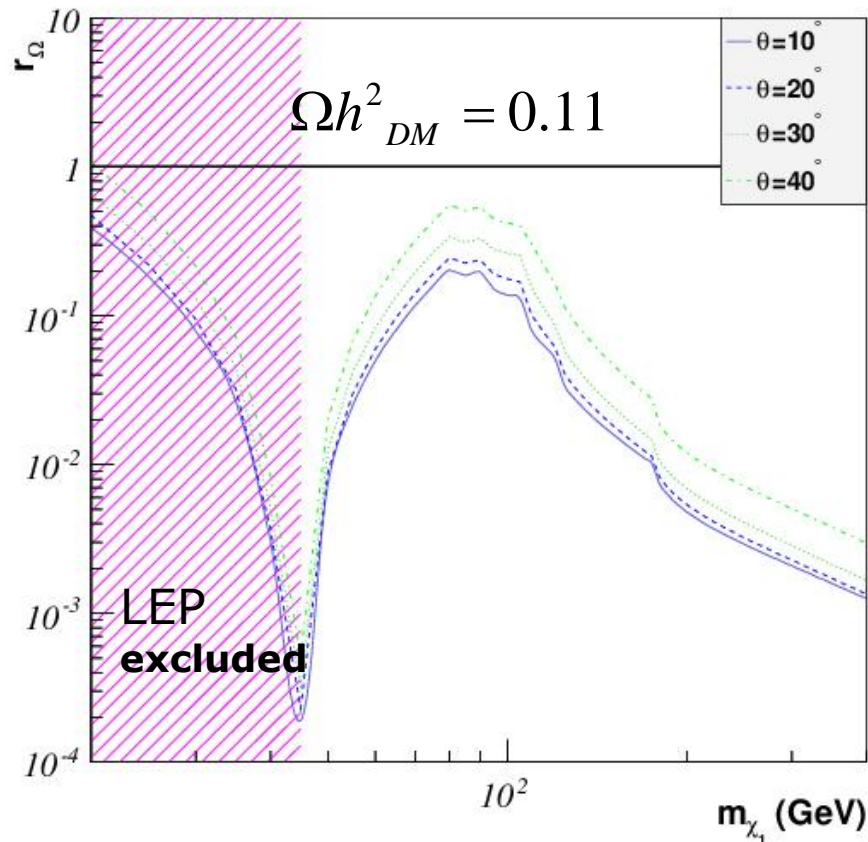
4) For $m_\chi > m_H$, Z^0h^0 , h^0h^0 channels open

5) For $m_\chi > m_t$, $t\bar{t}$ channel opens

Different from Dirac neutrino DM

- $f\bar{f}$ channels are helicity and velocity suppressed
- WW channel is velocity suppressed
- Zh more important

Relic density of heavy neutrino DM



strong mass-dependence

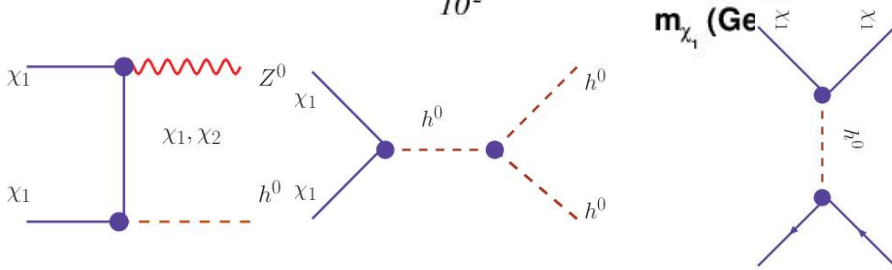
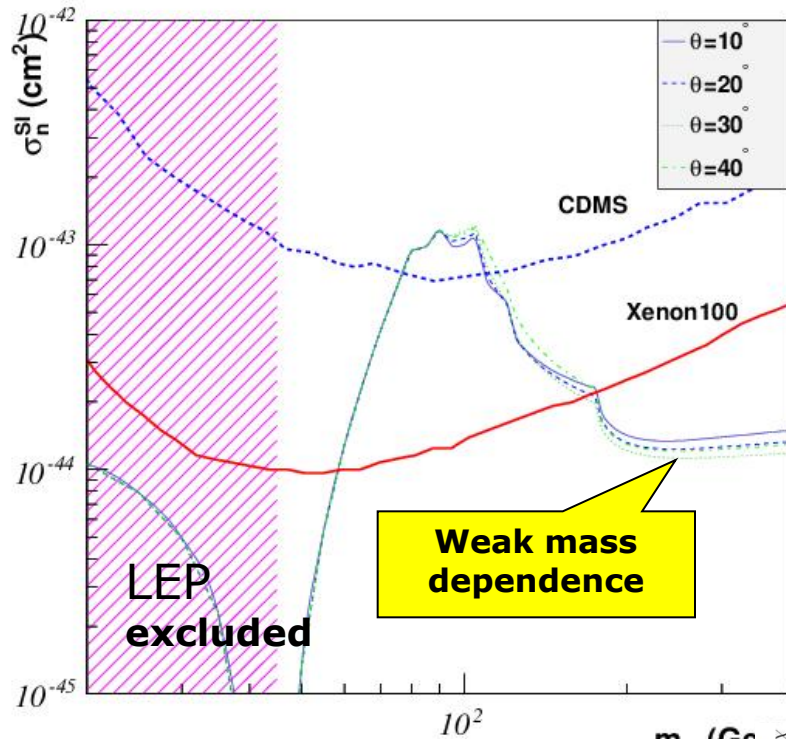
Relic density

$$\Omega h^2 \approx \frac{1.07 \times 10^9 \text{ GeV}^{-1}}{\sqrt{g_*} M_{pl} \int_{x_F}^{\infty} \frac{\langle \sigma v \rangle}{x^2} dx},$$

- No upper bound on neutrino mass
- Majorana neutrino can make up O(10%) of DM at 100 GeV, O(1%) at 200 GeV.

SI cross section

YFZ, arXiv:110.2930



Relic density

SI scattering

- Event rate rescaled

$$r_\rho \equiv \frac{\rho_1}{\rho_0} \approx \frac{\Omega_{\chi_1}}{\Omega_{DM}}, \quad \sigma^{SI} \rightarrow \tilde{\sigma}^{SI}$$

(mass dependent)

- Xenon ruled out mass range 55-175 GeV

- Correlation

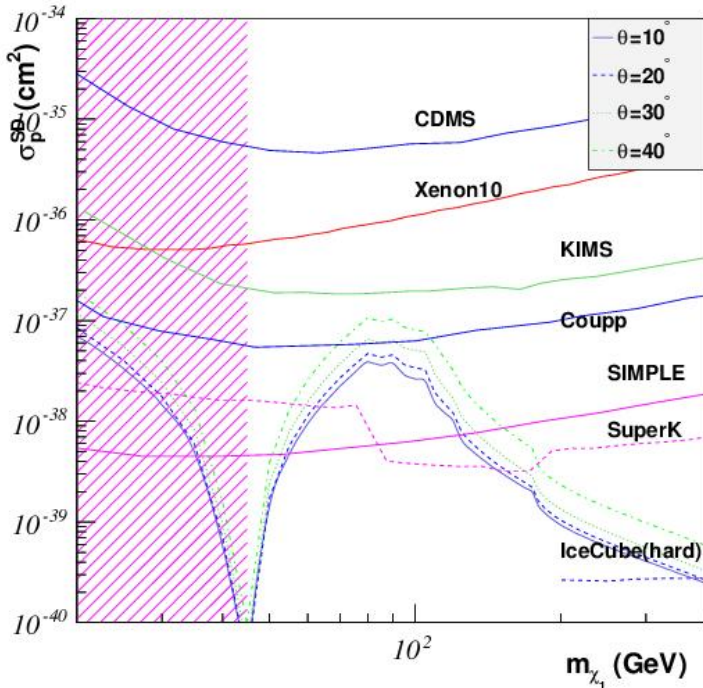
- Low halo density and large SI cross section all related to large Yukawa coupling
- > cancellation
- weak mixing angle/mass dependences

Prediction:

$$\tilde{\sigma}^{SI} \approx 1.5 \times 10^{-44} \text{ cm}^2 \text{ for } m_\chi > 175 \text{ GeV}$$

can be tested soon

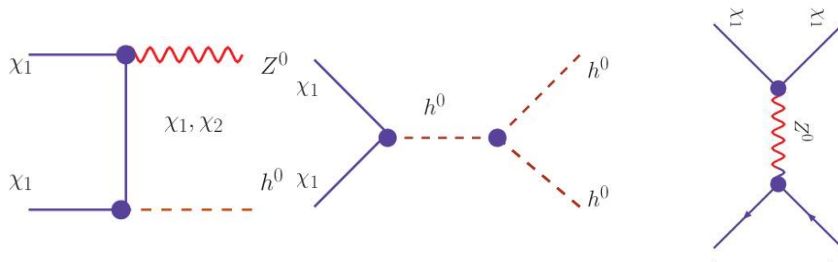
SD cross section for DM-proton



- Event rate follow the relic density
- SIMPLE ruled out mass range 50-150 GeV, compatible with Xenon100 on SI cross section
- Prediction

$$\tilde{\sigma}_p < 3 \times 10^{-39} \text{ cm}^2 \text{ for } m_\chi > 175 \text{ GeV}$$

- Can be reached by indirect search exp. SuperK and IceCube. The current SuperK/IceCube bounds are model dependent.



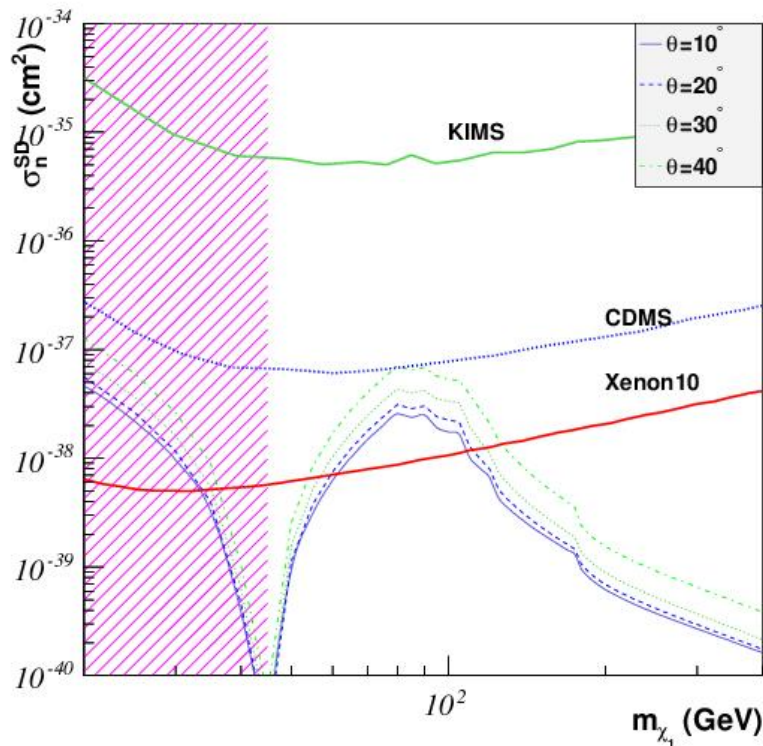
Relic density

SI scattering

$$\sigma_N^{SD} = \frac{32}{\pi} G_F^2 \mu_n^2 \frac{J+1}{J} \left(a_p \langle S_p \rangle + a_n \langle S_n \rangle \right)^2,$$

$$d_u = -d_d = -d_s = \frac{G_F}{\sqrt{2}}.$$

SD cross section for DM-neutron



- Scattering off proton and neutron are similar
- much weaker bounds compared with DM-proton case.
- The Xenon10 ruled out 60-120 GeV

Conclusions

- Heavy stable neutrino as a cold DM component is highly testable in direct detection experiments even it contribute to a small fraction of the total DM relic density
- We consider a 4th generation model with an anomaly-free U(1) gauge symmetry which keep the heavy 4th Majorana neutrino stable. The models is less constrained by the current LHC data.
- The Xenon100 data constrain the mass of the 4th neutrino to be heavier than 175 GeV. For heavier neutrino DM, the prediction for SI cross section is $1.5 \times 10^{-44} \text{cm}^2$ which is insensitive to the neutrino mass due to the correlation between relic density and the SI cross section.
- The prediction for SD cross section is within the reach of SuperK and IceCube.