

Mass and Rapidity dependent Top Forward-Backward Asymmetry

Dong-Won JUNG

KIAS, Korea

Based on

Phys.Lett.B691 (2010) 238-242,

Phys.Lett.B708 (2012) 157-161,

and more works in progress

with P. Ko (KIAS), J. S. Lee (JNU) and S. Nam (KISTI)

September 14, 2012 @ 2nd PHENO., KIAS

Introduction

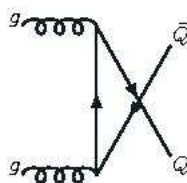
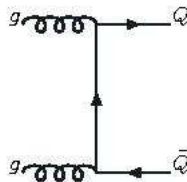
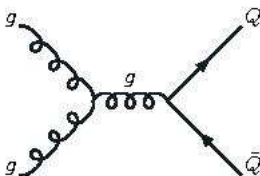
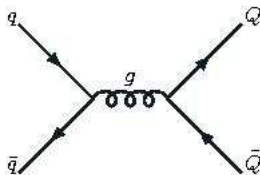
Effective Field Theory (EFT) approach

Mass and Rapidity dependent A_{FB}

Summary and Prospect

Tevatron with $\sqrt{S} = 1.96$ TeV,

- $\sigma_{t\bar{t}} = 7.50 \pm 0.48$ pb
- $q\bar{q} \rightarrow t\bar{t} \sim 85\%$
- $gg \rightarrow t\bar{t} \sim 15\%$



Recent Results @Tevatron ¹

- **CDF, l+jets:**

$$A_{FB}^{t\bar{t}} = 0.158 \pm 0.074 \text{ [} 0.058 \pm 0.009 \text{]}$$

- **CDF, dileptons:**

$$A_{FB}^{t\bar{t}} = 0.42 \pm 0.16$$

- **Combined:**

$$A_{FB}^{t\bar{t}} = 0.201 \pm 0.067 \quad cf. \quad A_{FB}^{t\bar{t}} = 0.196 \pm 0.065 \text{ @ D0}$$

- **CDF, l+jets:**

$$A_{FB}^{t\bar{t}}(|\Delta y| < 1.0) = 0.088 \pm 0.042 \pm 0.022 \text{ [} 0.043 \text{]}$$

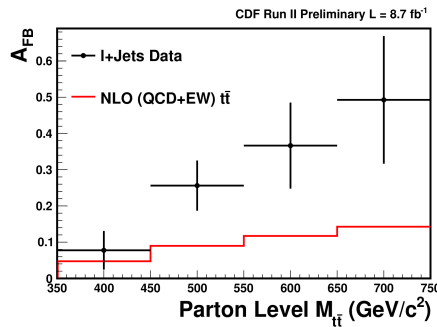
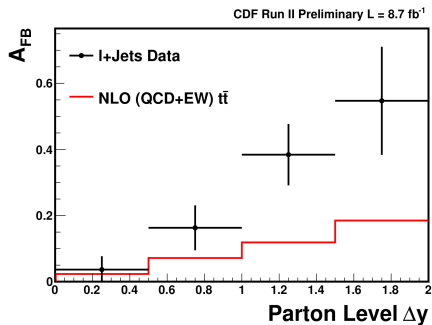
$$A_{FB}^{t\bar{t}}(|\Delta y| \geq 1.0) = 0.433 \pm 0.097 \pm 0.050 \text{ [} 0.139 \text{]}$$

- **CDF, l+jets:**

$$A_{FB}^{t\bar{t}}(M_{t\bar{t}} < 450 \text{ GeV}) = -0.078 \pm 0.048 \pm 0.024 \text{ [} 0.047 \text{]}$$

$$A_{FB}^{t\bar{t}}(M_{t\bar{t}} \geq 450 \text{ GeV}) = 0.296 \pm 0.059 \pm 0.031 \text{ [} 0.100 \text{]}$$

¹The values in the squared bracket are SM predictions. 



Effective Field Theory (EFT) approach

- If the NP particles are heavy enough, we can adopt the viewpoint of effective Lagrangian.
- Assuming $SU(2)_L \times U(1)_Y$ and the custodial symmetry $SU(2)_R$ for the light quark sector,

$$\mathcal{L}_6 = \frac{g_s^2}{\Lambda^2} \sum_{A,B} \left[C_{1q}^{AB} (\bar{q}_A \gamma_\mu q_A) (\bar{t}_B \gamma^\mu t_B) \right. \\ \left. + C_{8q}^{AB} (\bar{q}_A T^a \gamma_\mu q_A) (\bar{t}_B T^a \gamma^\mu t_B) \right]$$

where $T^a = \lambda^a/2$, $\{A, B\} = \{L, R\}$, and $L, R \equiv (1 \mp \gamma_5)/2$ with $q = (u, d)^T, (s, c)^T$.

- Cross section up to $O(1/\Lambda^2)$ is calculated, keeping only the **INTERFERENCE** between the SM and NP.

Amplitude

- **The amplitude for $q(p_1) + \bar{q}(p_2) \rightarrow t(p_3) + \bar{t}(p_4)$**

$$\overline{|\mathcal{M}|^2} \simeq \frac{4 g_s^4}{9 \hat{s}^2} \left\{ 2 m_t^2 \hat{s} \left[1 + \frac{\hat{s}}{2 \Lambda^2} (C_1 + C_2) \right] s_{\hat{\theta}}^2 + \frac{\hat{s}^2}{2} \left[\left(1 + \frac{\hat{s}}{2 \Lambda^2} (C_1 + C_2) \right) (1 + c_{\hat{\theta}}^2) + \hat{\beta}_t \left(\frac{\hat{s}}{\Lambda^2} (C_1 - C_2) \right) c_{\hat{\theta}} \right] \right\}$$

where $C_1 \equiv C_{8q}^{LL} + C_{8q}^{RR}$ and $C_2 \equiv C_{8q}^{LR} + C_{8q}^{RL}$

- $\hat{s} = (p_1 + p_2)^2$, $\hat{\beta}_t^2 = 1 - 4m_t^2/\hat{s}$, and $s_{\hat{\theta}} \equiv \sin \hat{\theta}$ and $c_{\hat{\theta}} \equiv \cos \hat{\theta}$, with $\hat{\theta}$ being the polar angle between the incoming quark and the outgoing top quark in the $t\bar{t}$ rest frame.

Result for Integrated A_{FB}

- **Convolutd with CTEQ6L and K factor 1.3.**
- A_{FB} must be expanded with $\sigma_{NLO}^{int}, \sigma_{NP}^{int}$,

$$A_{FB} \simeq A_{FB}^{SM} + \Delta\sigma_{NP}^{int}/\sigma_0^{SM}.$$

- **For validity check, we impose three criteria,**

$$\sigma = \sigma_{SM} + \sigma_{NP}^{int} + \sigma_{NP},$$

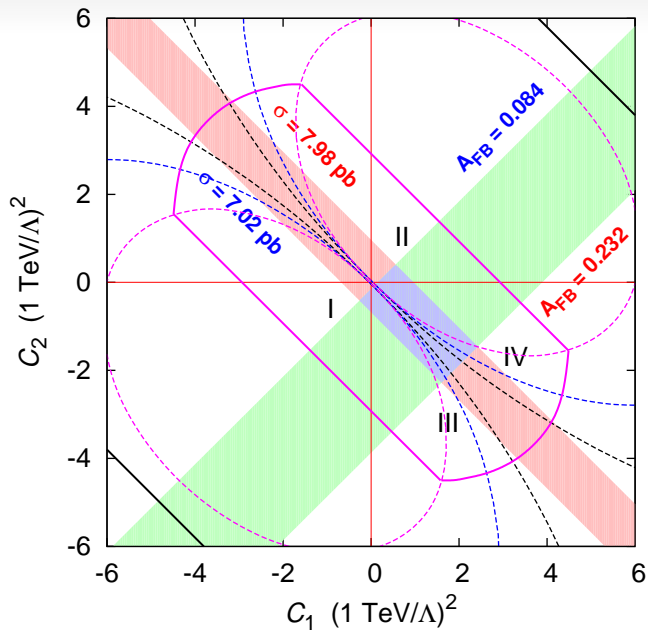
$$|\sigma_{NP}^{int}| < r \times \sigma_{SM} \quad (\text{straight})$$

$$\sigma_{NP} < r \times |\sigma_{NP}^{int}| \quad (\text{two ellipses adjacent at the origin})$$

$$\sigma_{NP} < r^2 \times \sigma_{SM} \quad (\text{ellipses centered at the origin})$$

with $r = 0.3, 0.5, 1.0$.

- $\Delta\sigma_{t\bar{t}} \equiv \sigma_{t\bar{t}} - \sigma_{t\bar{t}}^{SM} \propto (C_1 + C_2),$
 $\Delta A_{FB} \equiv A_{FB} - A_{FB}^{SM} \propto (C_1 - C_2).$



- General Lagrangian

$$\begin{aligned}
 \mathcal{L}_{\text{int}} = & g_s V_8^{a\mu} \sum_A \left[g_{8q}^A (\bar{q}_A \gamma_\mu T^a q_A) + g_{8t}^A (\bar{t}_A \gamma_\mu T^a t_A) \right] \\
 & + g_s \left[\tilde{V}_1^\mu \sum_A \tilde{g}_{1q}^A (\bar{t}_A \gamma_\mu q_A) + \tilde{V}_8^{a\mu} \sum_A \tilde{g}_{8q}^A (\bar{t}_A \gamma_\mu T^a q_A) + \text{h.c.} \right] \\
 & + g_s \left[\tilde{S}_1 \sum_A \tilde{\eta}_{1q}^A (\bar{t}_A q) + \tilde{S}_8^a \sum_A \tilde{\eta}_{8q}^A (\bar{t}_A T^a q) + \text{h.c.} \right],
 \end{aligned}$$

- Integrating out heavy fields, the Wilson coefficients are,

$$\begin{aligned}
 \frac{C_{8q}^{RR}}{\Lambda^2} &= -\frac{g_{8q}^R g_{8t}^R}{m_{V_{8R}}^2} - \frac{2|\tilde{g}_{1q}^R|^2}{m_{\tilde{V}_{1R}}^2} + \frac{1}{N_c} \frac{|\tilde{g}_{8q}^R|^2}{m_{\tilde{V}_{8R}}^2}, \\
 \frac{C_{8q}^{LL}}{\Lambda^2} &= -\frac{g_{8q}^L g_{8t}^L}{m_{V_{8L}}^2} - \frac{2|\tilde{g}_{1q}^L|^2}{m_{\tilde{V}_{1L}}^2} + \frac{1}{N_c} \frac{|\tilde{g}_{8q}^L|^2}{m_{\tilde{V}_{8L}}^2}, \\
 \frac{C_{8q}^{LR}}{\Lambda^2} &= -\frac{g_{8q}^L g_{8t}^R}{m_{V_8}^2} - \frac{|\tilde{\eta}_{1q}^L|^2}{m_{\tilde{S}_{1L}}^2} + \frac{1}{2N_c} \frac{|\tilde{\eta}_{8q}^L|^2}{m_{\tilde{S}_{8L}}^2}, \\
 \frac{C_{8q}^{RL}}{\Lambda^2} &= -\frac{g_{8q}^R g_{8t}^L}{m_{V_8}^2} - \frac{|\tilde{\eta}_{1q}^R|^2}{m_{\tilde{S}_{1R}}^2} + \frac{1}{2N_c} \frac{|\tilde{\eta}_{8q}^R|^2}{m_{\tilde{S}_{8R}}^2},
 \end{aligned}$$

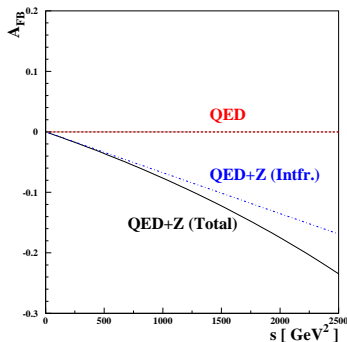
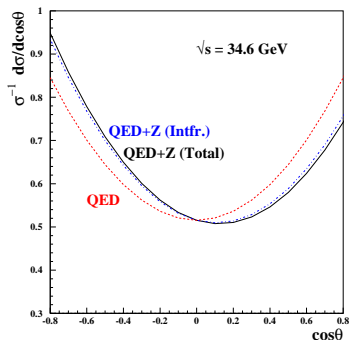
Summary Table

New particles	couplings	C_1	C_2	1 σ favor
V_8 (spin-1 FC octet)	$g_{8q,8t}^{L,R}$	indef.	indef.	\checkmark
\tilde{V}_1 (spin-1 FV singlet)	$\tilde{g}_{1q}^{L,R}$	—	0	\times
\tilde{V}_8 (spin-1 FV octet)	$\tilde{g}_{8q}^{L,R}$	+	0	\checkmark
\tilde{S}_1 (spin-0 FV singlet)	$\tilde{\eta}_{1q}^{L,R}$	0	—	\checkmark
\tilde{S}_8 (spin-0 FV octet)	$\tilde{\eta}_{8q}^{L,R}$	0	+	\times
S_2^α (spin-0 FV triplet)	η_3	—	0	\times
$S_{13}^{\alpha\beta}$ (spin-0 FV sextet)	η_6	+	0	\checkmark

Mass and Rapidity dependent A_{FB}

- **Old Wisdom from Weak Interaction, RETRA :**
Far below the Z pole mass,

$$A_{FB}(s) \simeq -\frac{3G_F}{\sqrt{2}} \frac{s}{4\pi\alpha} (g_L - g_R)^2 \equiv kG_F s \simeq -7.18 G_F s.$$



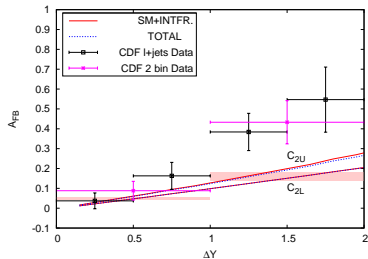
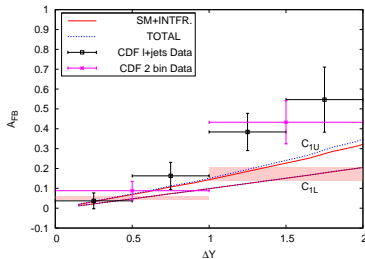
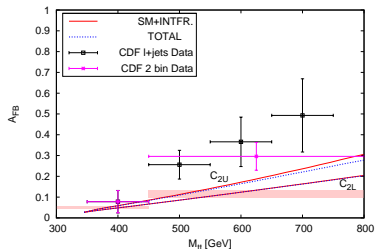
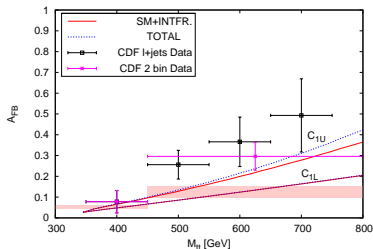
The case for $q\bar{q} \rightarrow t\bar{t}$

$$\begin{aligned}\hat{A}_{\text{FB}}(M_{t\bar{t}}) &= \frac{\hat{\beta}_t \frac{\hat{s}}{\Lambda^2} (C_1 - C_2)}{\frac{8}{3} \left[1 + \frac{\hat{s}}{2\Lambda^2} (C_1 + C_2) \right] + \frac{16\hat{s}}{3m_t^2} \left[1 + \frac{\hat{s}}{2\Lambda^2} (C_1 + C_2) \right]} \\ &\simeq \frac{3\hat{\beta}_t \frac{\hat{s}}{\Lambda^2} (C_1 - C_2)}{8 + 16\frac{m_t^2}{\hat{s}}}.\end{aligned}$$

- **FB asymmetry near the threshold is approximately linear in \hat{s} .**
- **Some nontrivial structure like wiggles or it changes the shape, one can say more about the underlying physics.**

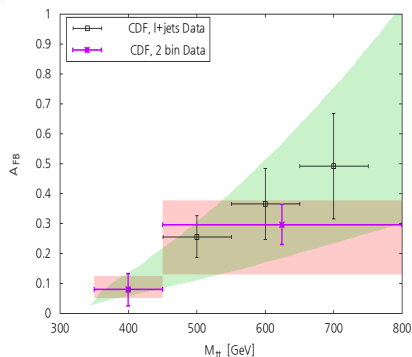
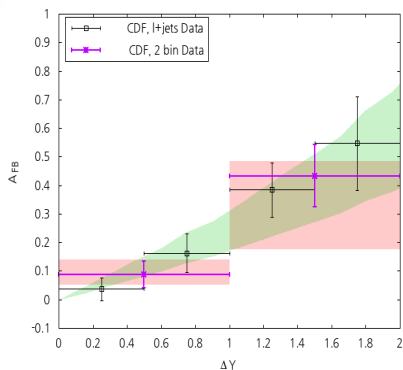
- Most models predict that only one of C_1 or C_2 nonzero.

$$(C_1, C_2) = (0.15 \sim 0.97, 0) \quad \text{or} \quad (C_1, C_2) = (0, -0.67 \sim -0.15).$$



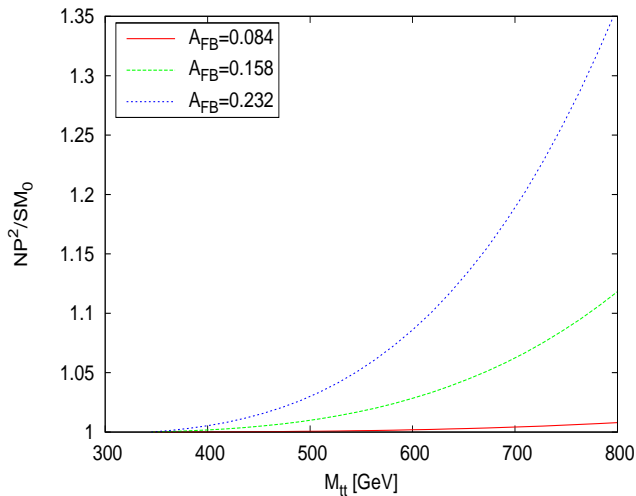
More general analysis

- Through the region $A_{FB}^{t\bar{t}} = 0.158 \pm 0.074 [0.058 \pm 0.009]$.
- Assuming that $C_1 + C_2 = 0$, which means no contribution to the cross section.
→ Reasonable assumption since variation of σ is smaller than A_{FB} .
- To minimize the NP^2 contribution, we take $C_{8q}^{RR} = C_{8q}^{LL} = \frac{1}{2}C_1$ and $C_{8q}^{RL} = C_{8q}^{LR} = -\frac{1}{2}C_1$. In this case there is no contribution to \hat{A}_{FB} from NP^2 terms.
- NP^2 increases as larger $M_{t\bar{t}}$ and ΔY . → Validity of EFT??



- Well fit!
- Central value of A_{FB} is favored.
- Validity issue : $NP^2/SM0$ goes upto 0.08%, 12% and 35% in the high M_{tt} and ΔY region.

e.g.,



Summary and Prospect

- M_{tt} and ΔY dependent \hat{A}_{FB} can be well fit with EFT approach.
- Preferred region for the Wilson coefficients can be selected : Model discrimination?
- For large M_{tt} and ΔY region there is validity issue of EFT approach.
- LHC study and longitudinal observables etc. ..?