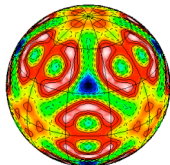


# Iridates

## The New Frontier

Mike Norman

Materials Science Division  
Argonne National Laboratory



KIAS, Dec. 19, 2009

Spin Liquids are insulators that have large Curie-Weiss temperatures, but do not order or show spin glass behavior. Particularly relevant are those with  $S=1/2$ .

## PHYSICS

# An End to the Drought of Quantum Spin Liquids

Patrick A. Lee

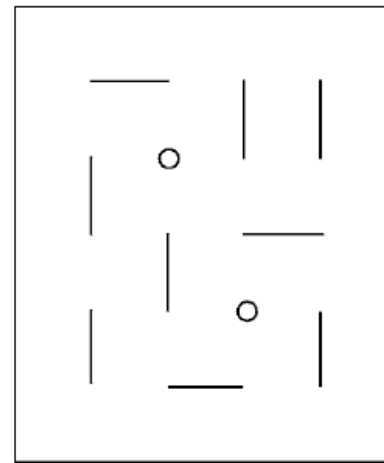
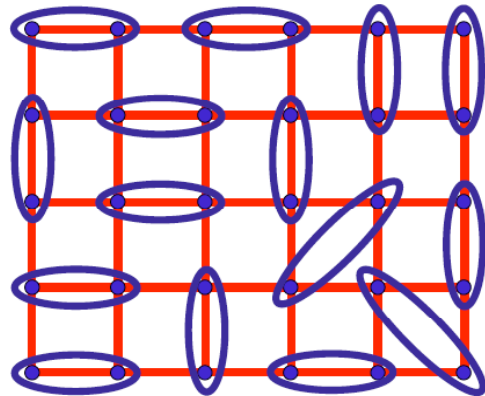
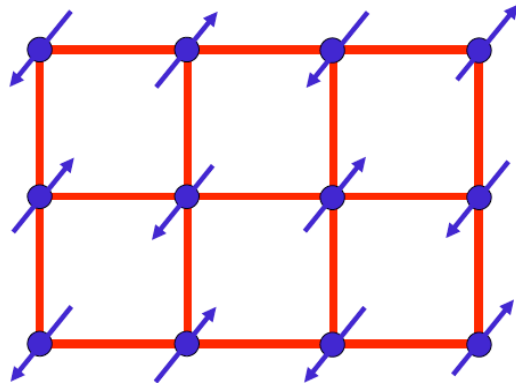
After decades of searching, several promising examples of a new quantum state of matter have now emerged.

Science 321, 1306 (2008)

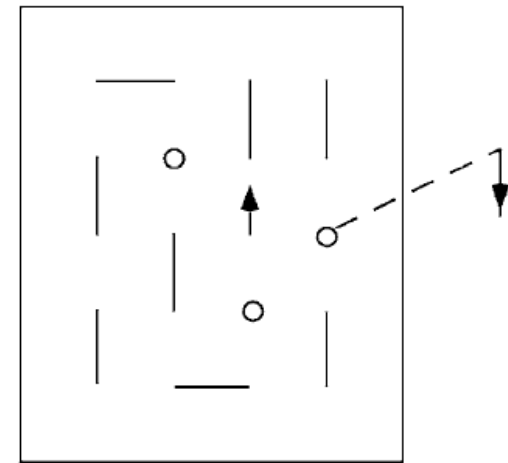
- $\kappa$ -(BEDT-TTF)<sub>2</sub>Cu<sub>2</sub>(CN)<sub>3</sub> (triangular)
- ZnCu<sub>3</sub>(OH)<sub>6</sub>Cl<sub>2</sub> (kagome)
- Na<sub>4</sub>Ir<sub>3</sub>O<sub>8</sub> (hyper-kagome)

(See also article in Physics Today by Barbara Goss Levi, page 16, February 2007)

## RVB - a liquid of spin singlets



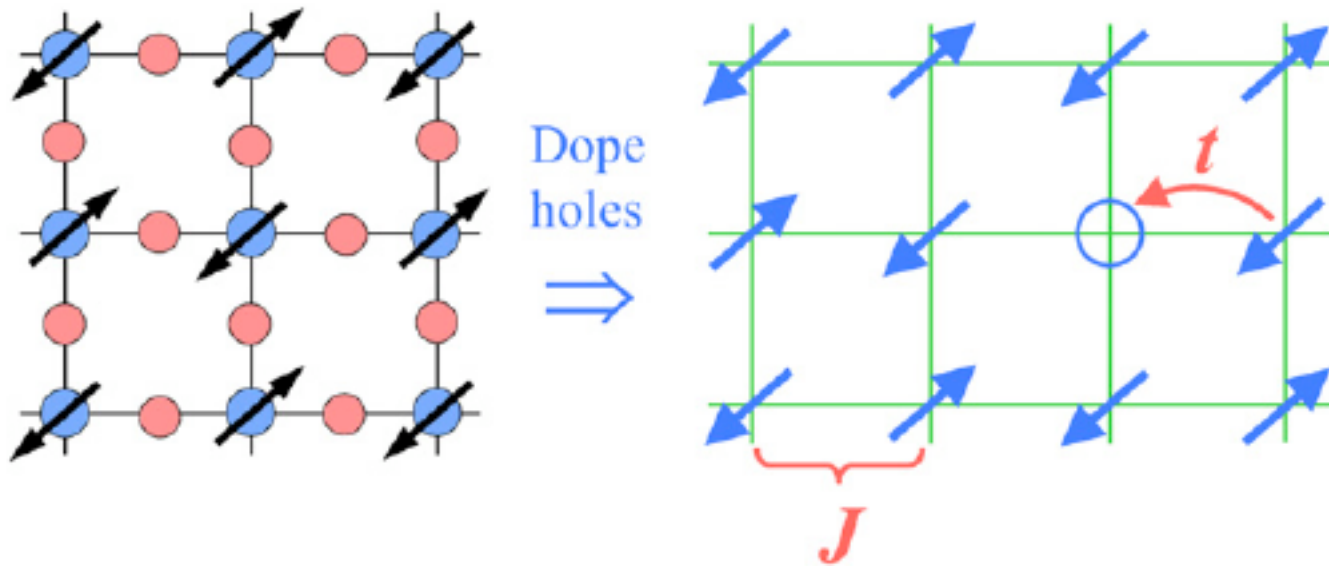
(a)



(b)

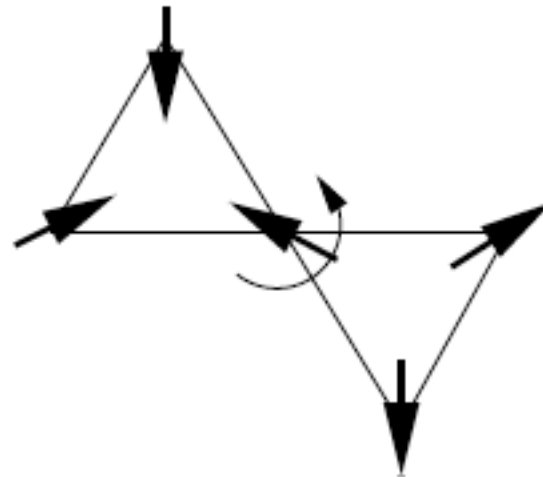
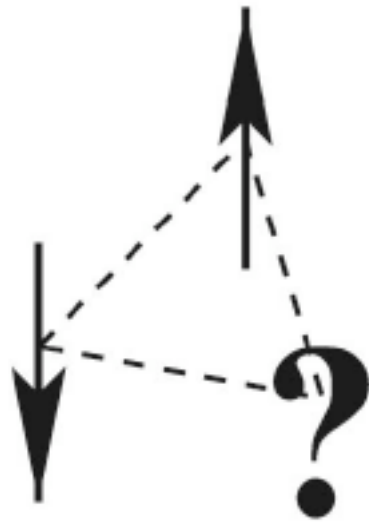
In the uniform RVB state, the  $S=1/2$  excitations (spinons) possess a zero energy Fermi surface (Anderson, 1987)

## Square Lattice (Anderson, 1987)



$J \mathbf{S} \cdot \mathbf{S}$  where  $J = 2 t^2 / U$   
 for  $S = 1/2$   
 Neel -  $J/4$  per bond -  $S_z S_z$   
 Singlet -  $3J/4$  per bond -  $S(S+1)$

## Frustration on the Triangular Lattice (Anderson, 1973)



$$\chi^{-1} = T - \Theta_{CW}$$

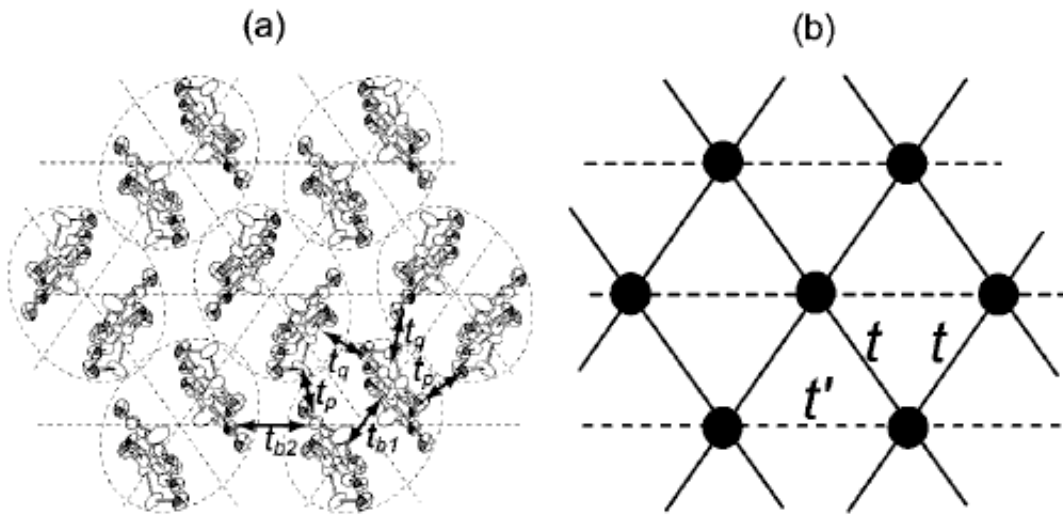
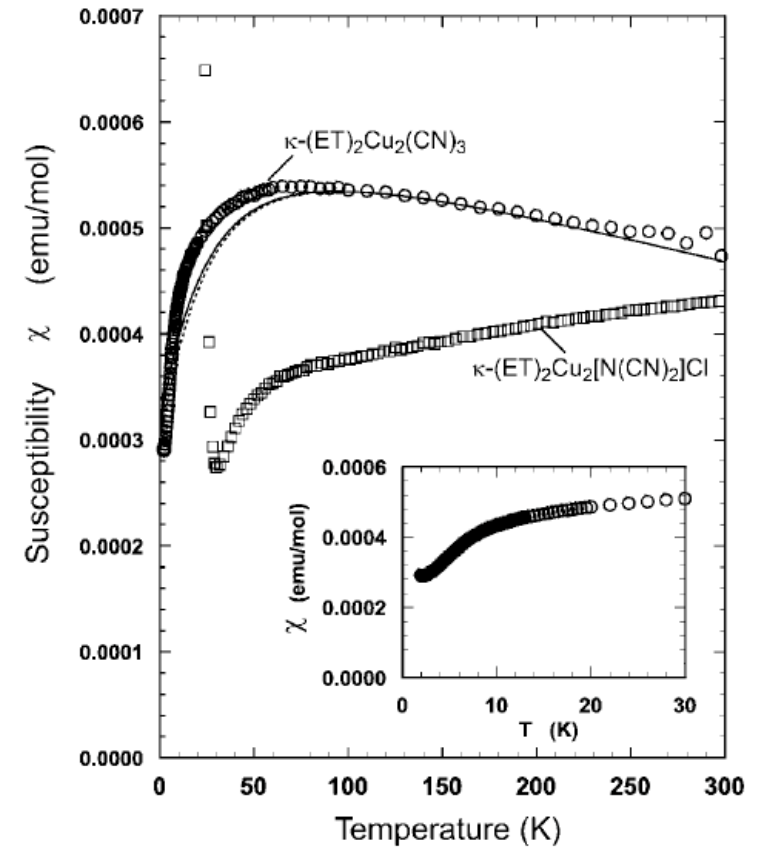
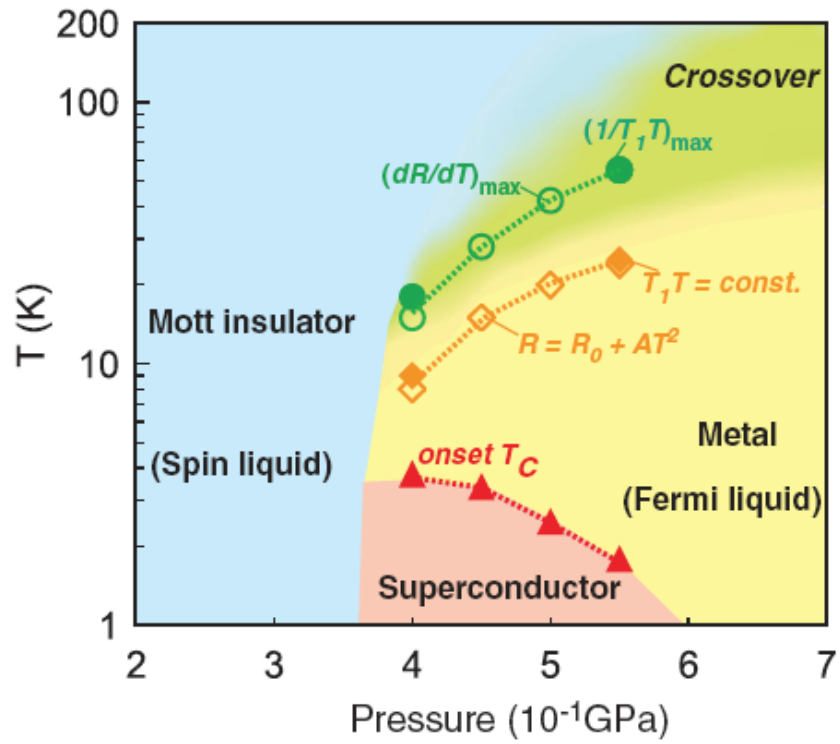


FIG. 1. (a) Crystal structure of an ET layer of  $\kappa\text{-(ET)}_2\text{Cu}_2(\text{CN})_3$  viewed along the long axes of ET molecules

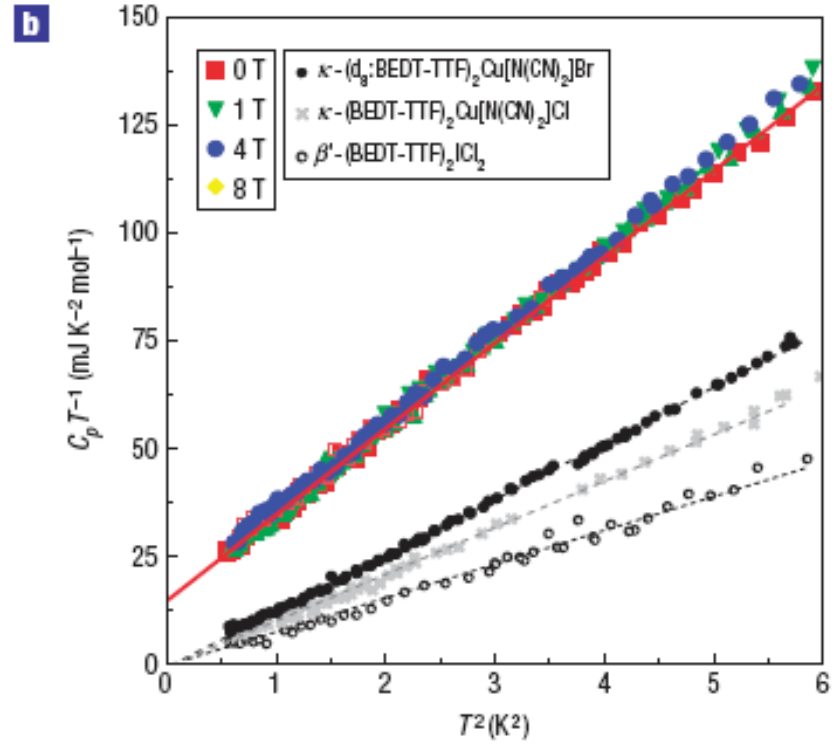
Geiser *et al.*, Inorg. Chem. 30, 2586 (1991)



Shimizu *et al.*, PRL 91, 107001 (2003)

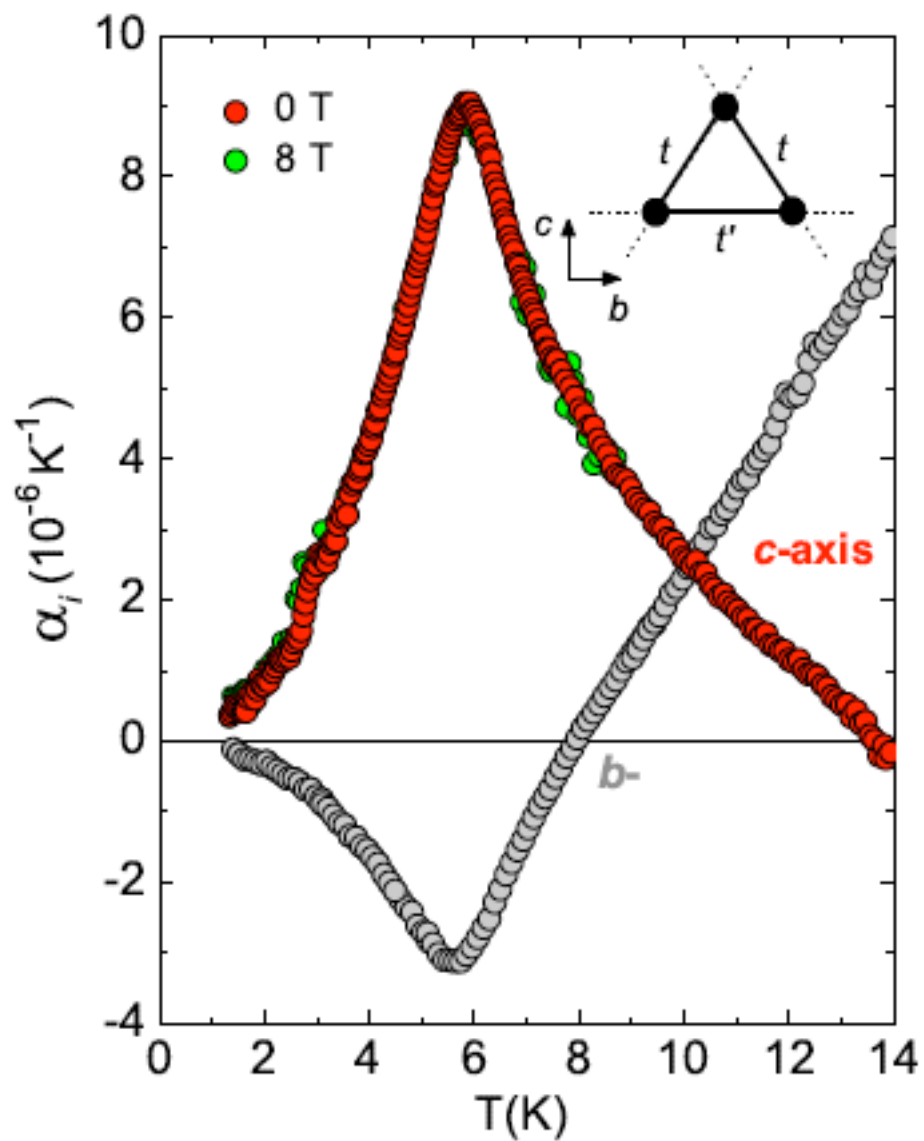


Kurosaki *et al.*, PRL 95, 177001 (2005)



Yamashita *et al.*, Nat. Phys. 4, 459 (2008)

## Ordering at 6 K?





# The Universe is a String-Net Liquid



Herbertsmithite

Zeeya Merali, *New Scientist* (15 March 2007)

# Herbertsmithite $\text{ZnCu}_3(\text{OH})_6\text{Cl}_2$

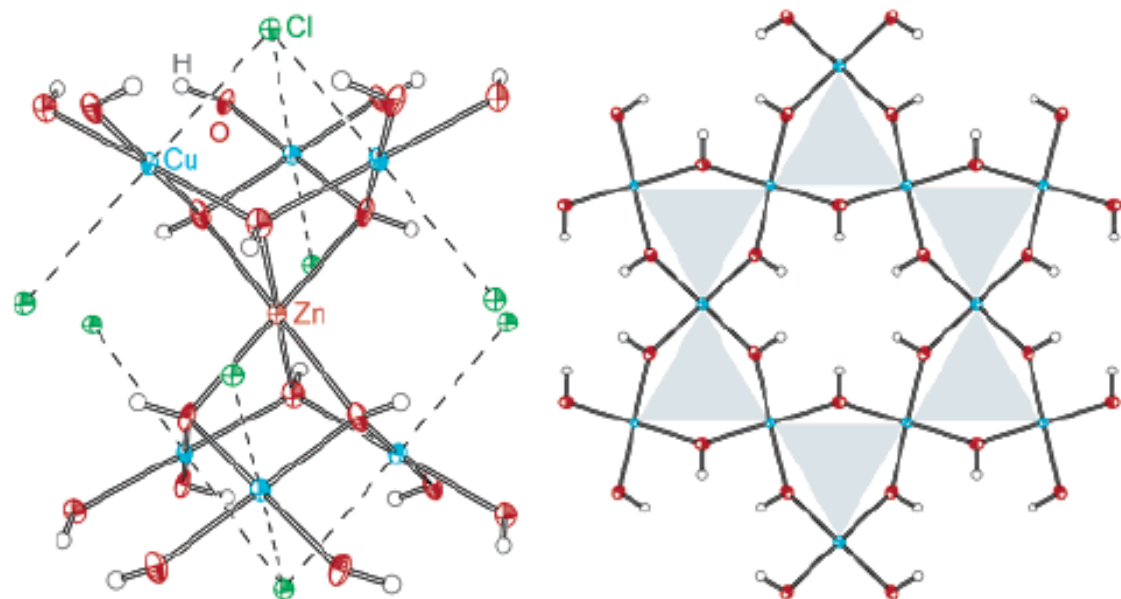
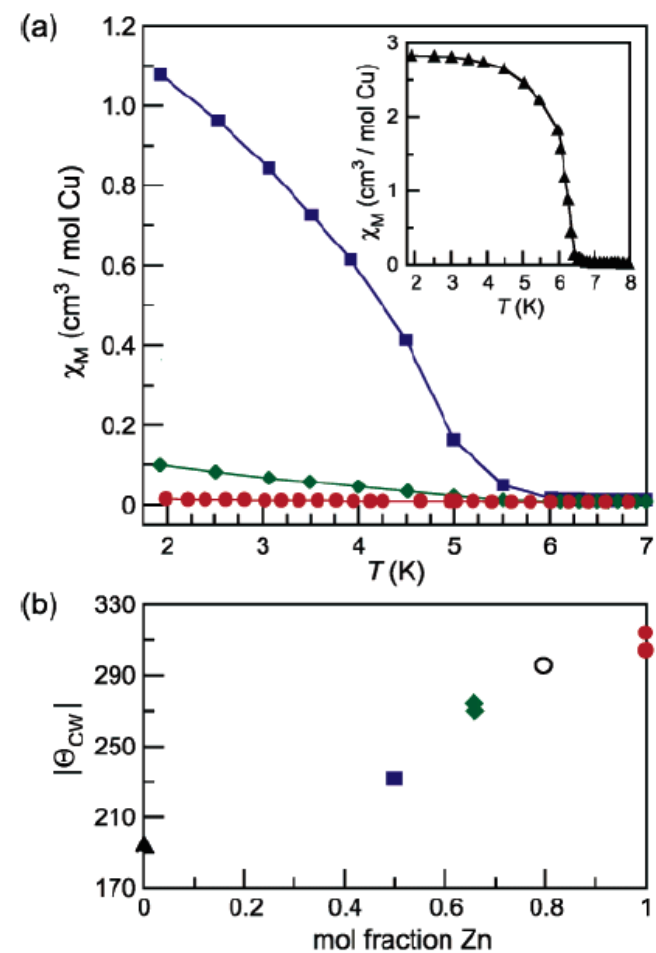
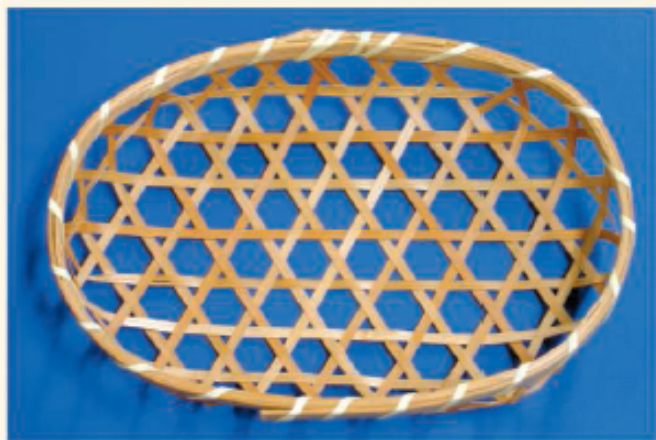
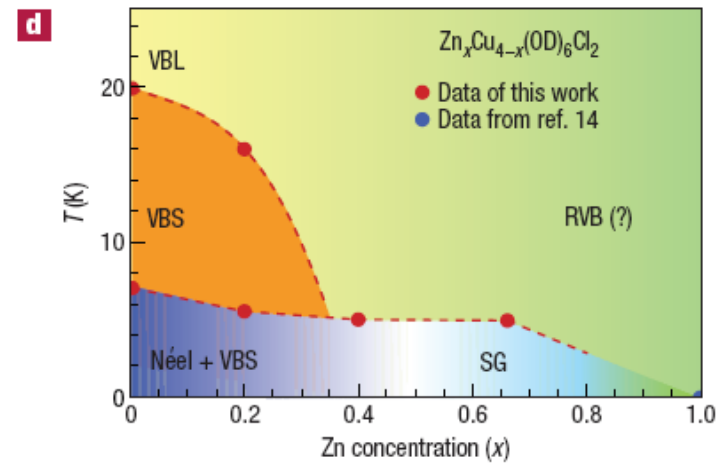
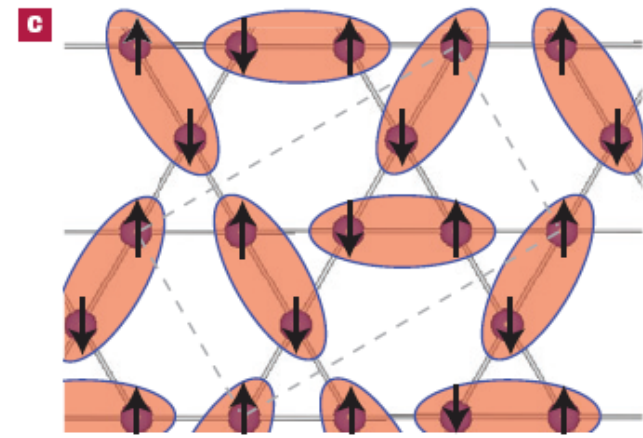
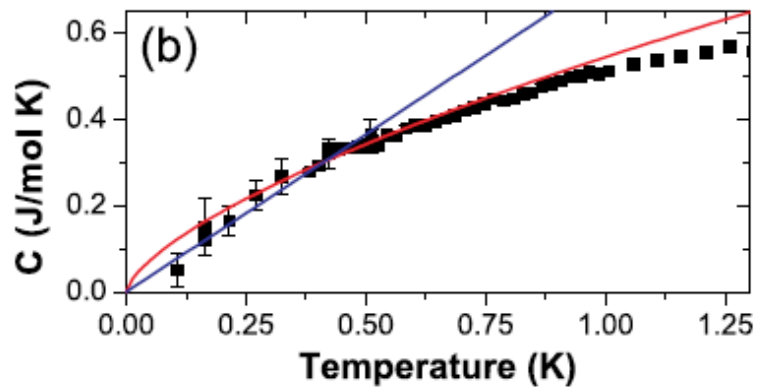
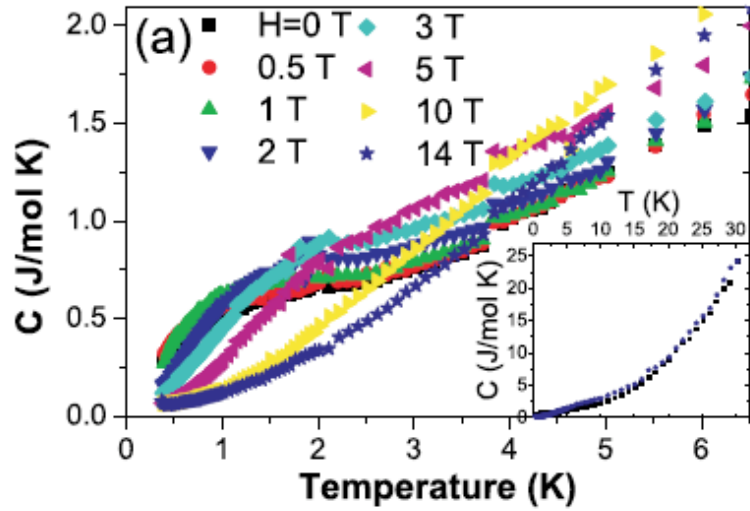


Figure 1. Crystal structure of Zn-paratacamite (I),  $\text{Zn}_{0.33}\text{Cu}_{3.67}(\text{OH})_6\text{Cl}_2$ .



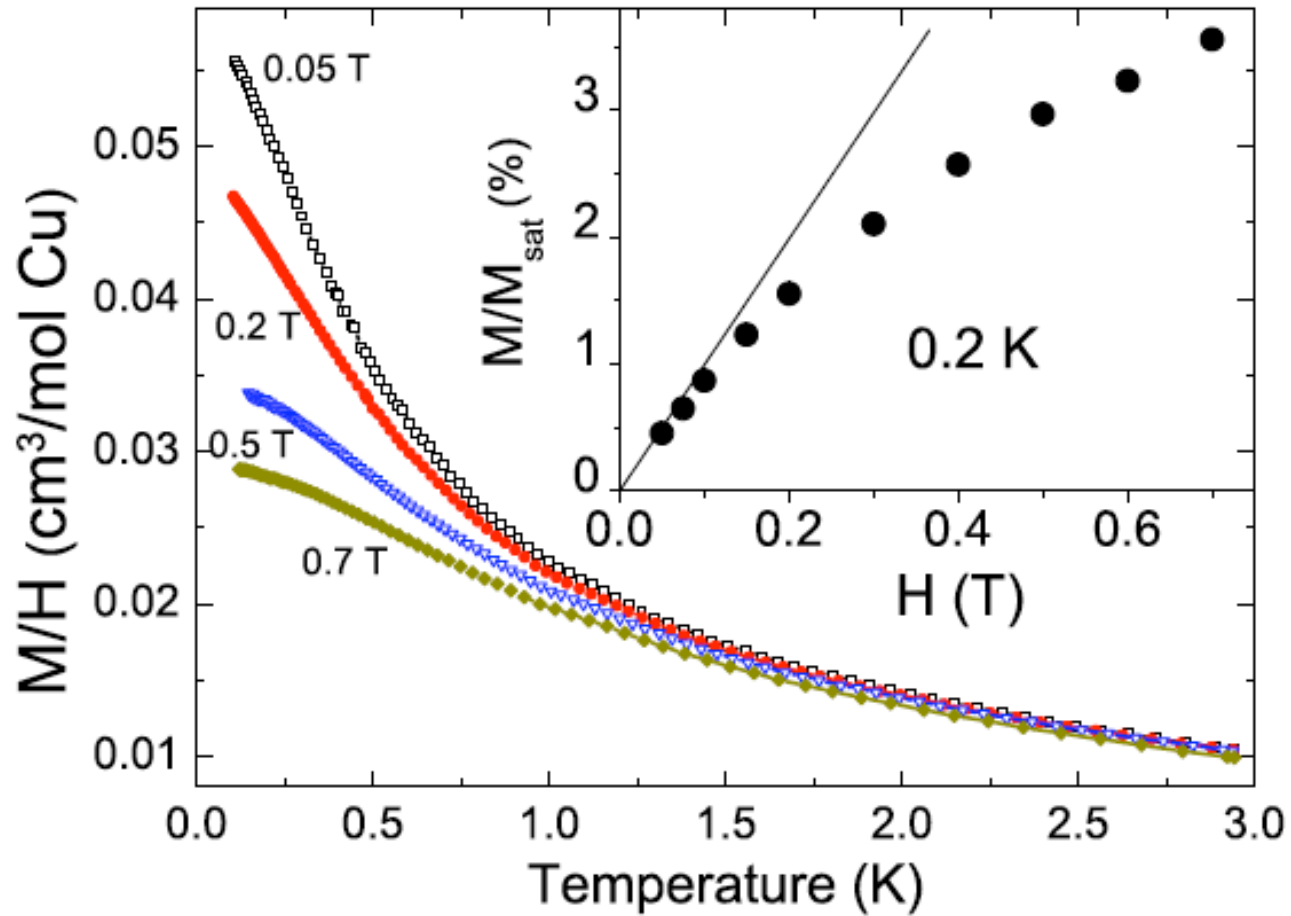
Shores *et al.*, JACS 127, 13462 (2005)



Helton *et al.*, PRL 98, 107204 (2007)

Lee *et al.*, Nat. Matls. 6, 853 (2007)

But, there is a large defect concentration (Cu - Zn)



Bert *et al.*, PRB 76, 132411 (2007)

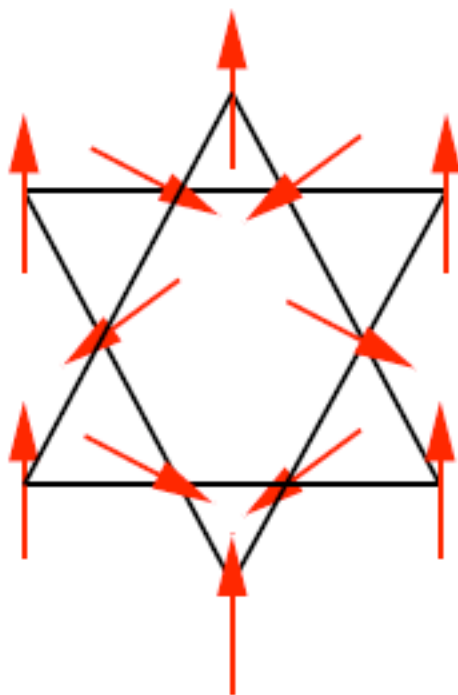
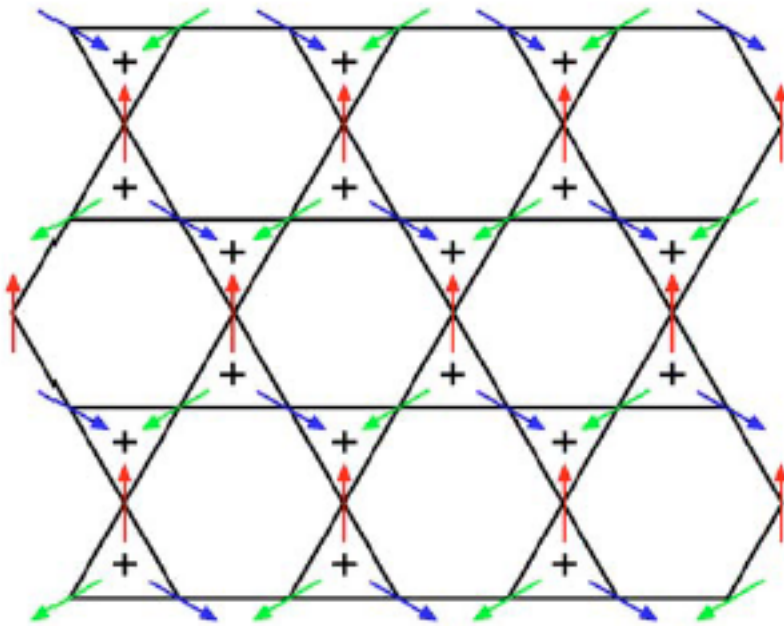


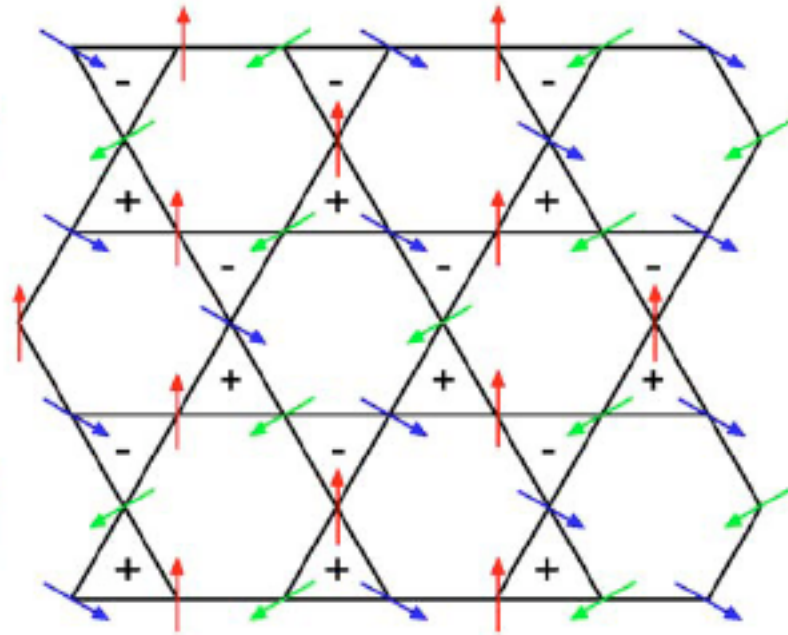
Fig. 1.5. Illustration of how ground state degrees of freedom arise for the Heisenberg model on the kagome lattice: spins on the central hexagon may be rotated together through any angle about the axis defined by the outer spins, without leaving the ground state.

# Classical Kagome Ground States

$q=0$

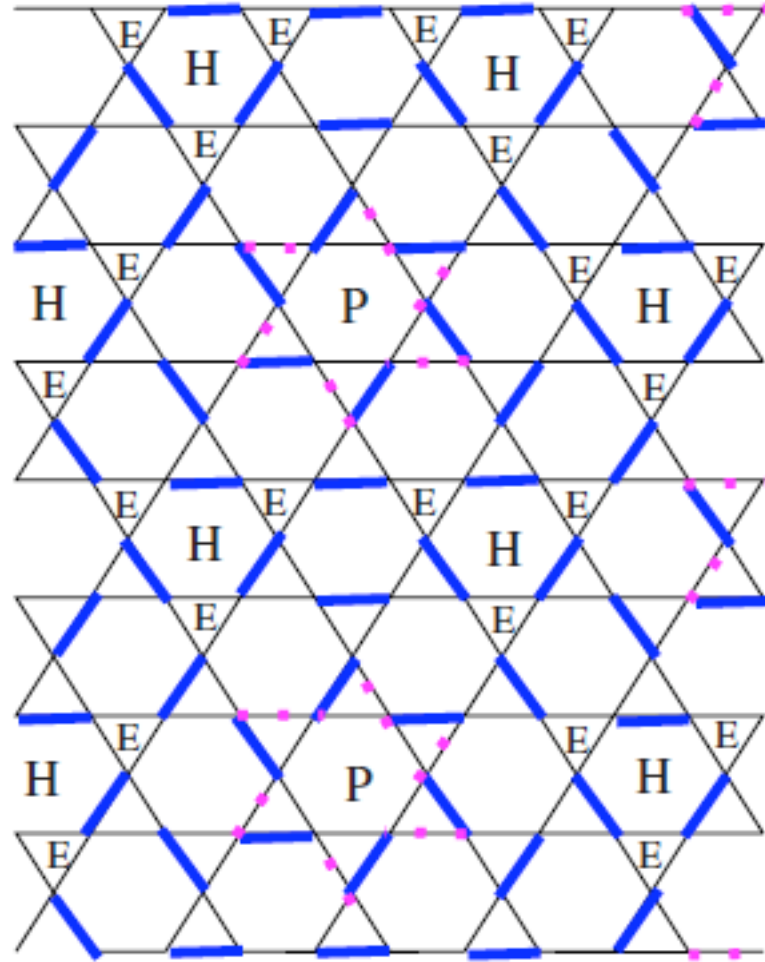


$\sqrt{3} \times \sqrt{3}$



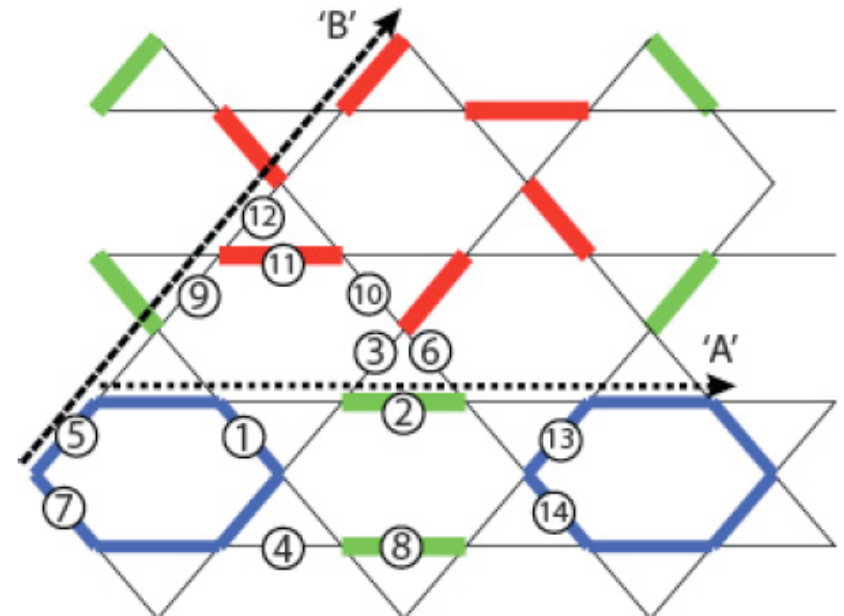
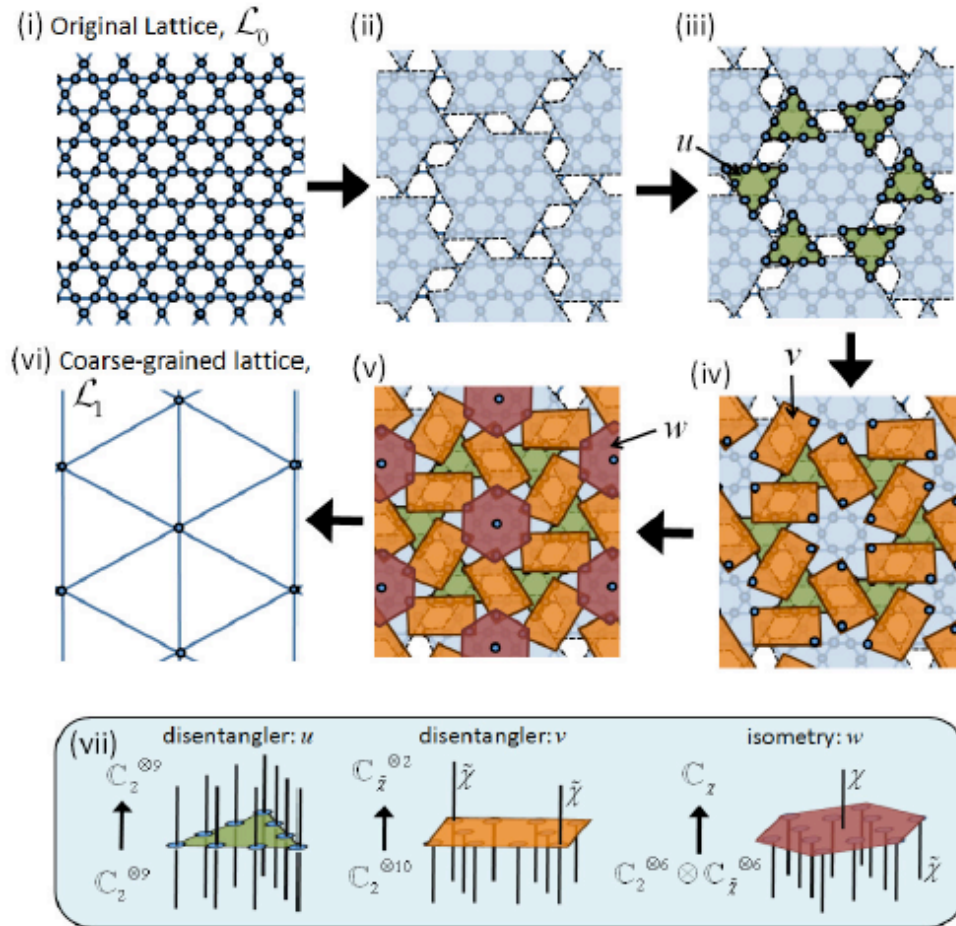
Ryu *et al.*, PRB 75, 184406 (2007)

# Valence Bond Ground State for Kagome?



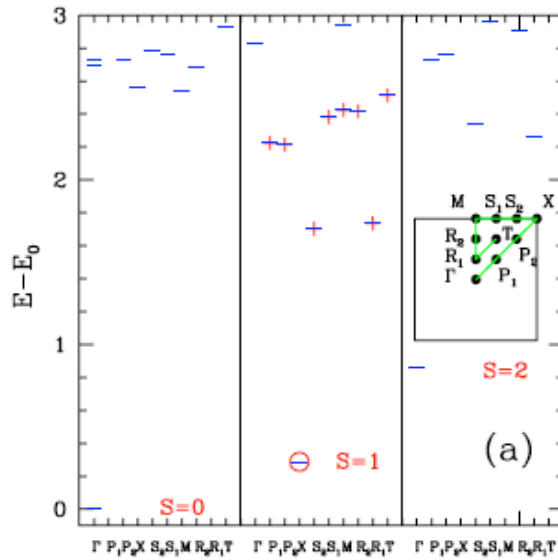
Singh & Huse, PRB 76, 180407 (2007)

# Multi-scale entanglement renormalization ansatz (MERA)

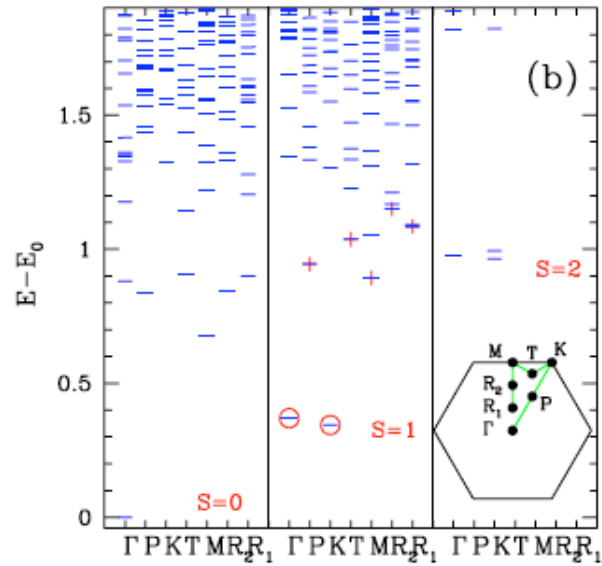




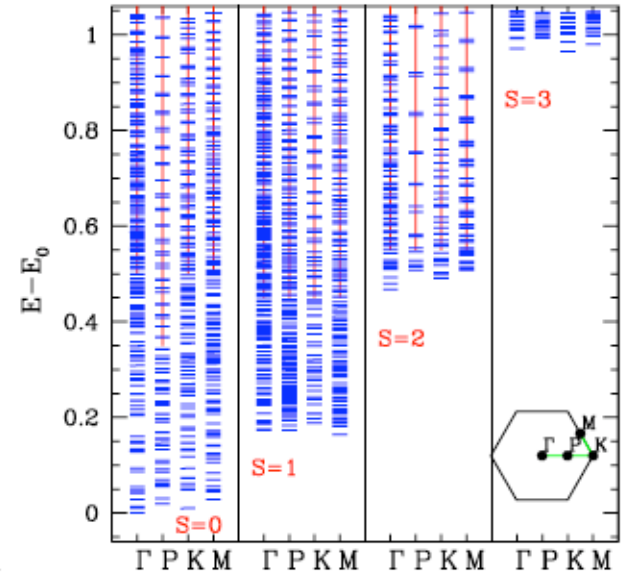
# Exact Diagonalization (36 sites)



square



triangular

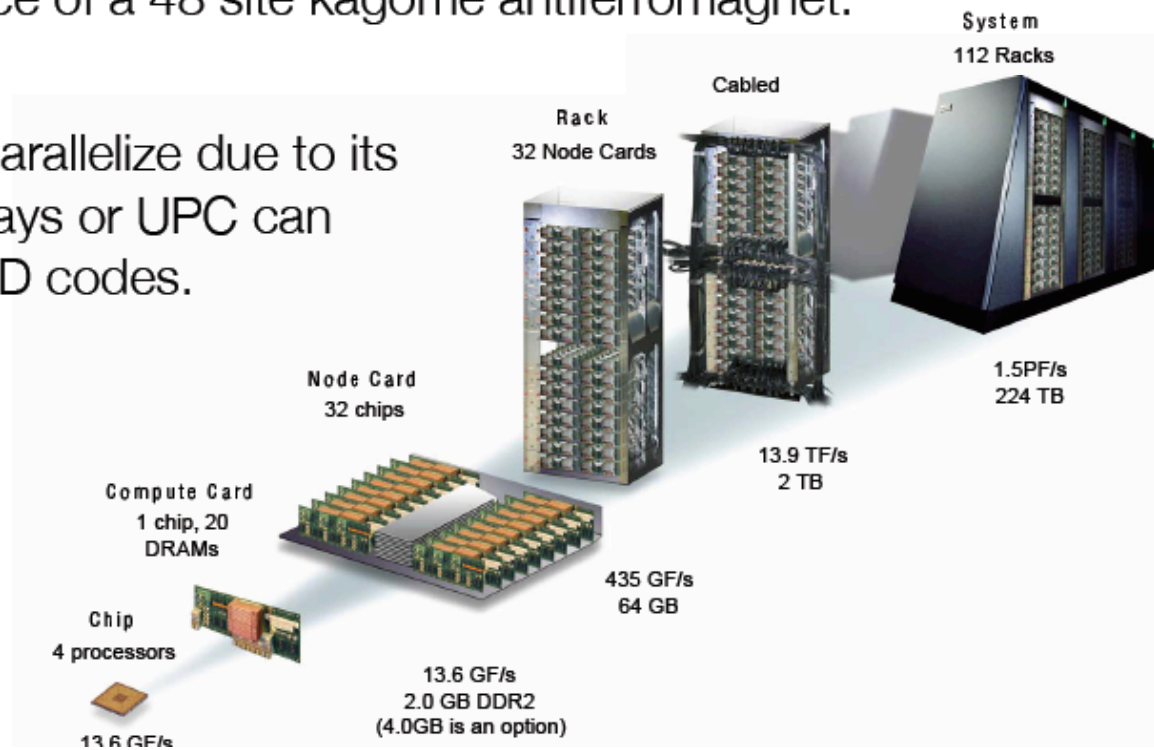


Kagome



# Parallelization: How to harness the petaflop computers ?

- Cutting edge petaflop systems have a huge number of core, but only a moderate amount of node-local memory.
- Next generation ED codes need to be developed in order to attack e.g. the 80 billion Hilbert space of a 48 site kagome antiferromagnet.
- Problem remains difficult to parallelize due to its all-to-all structure. Global Arrays or UPC can help developing distributed ED codes.



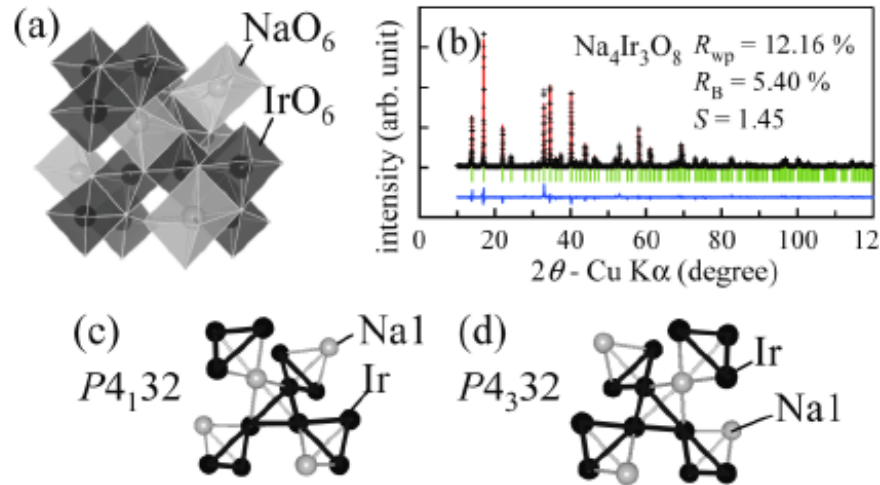
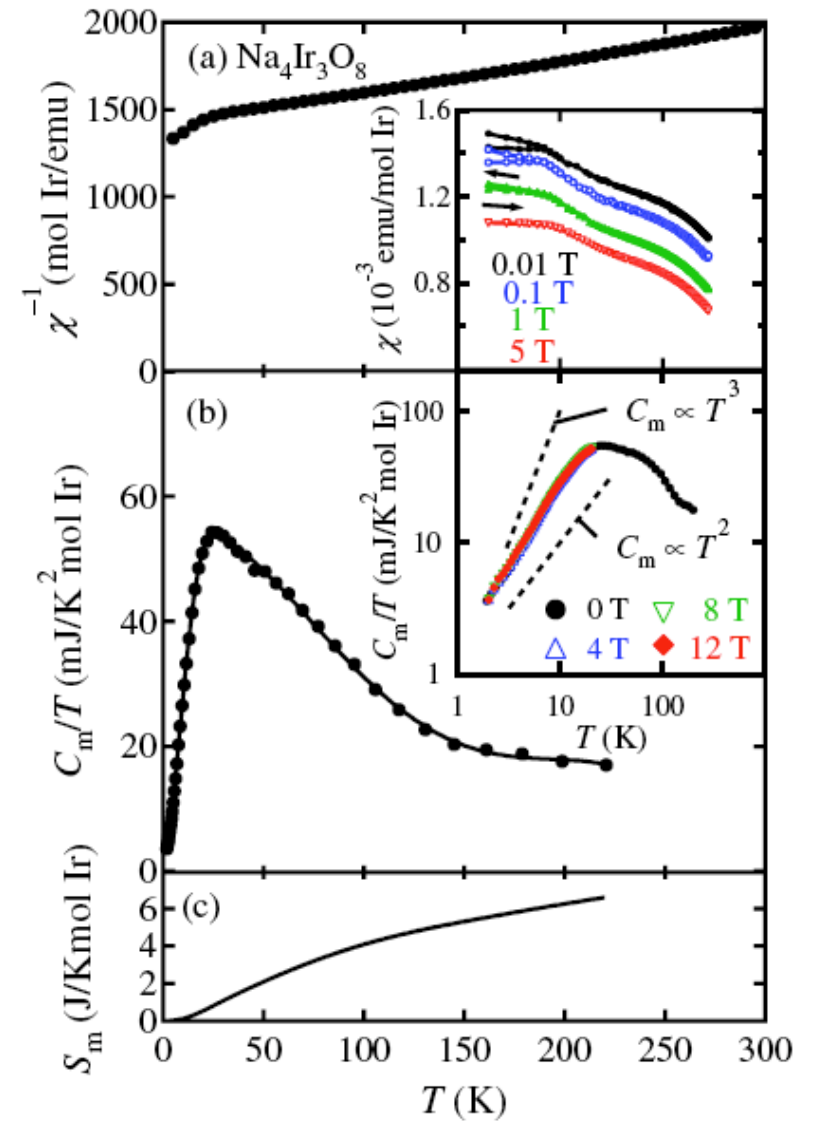


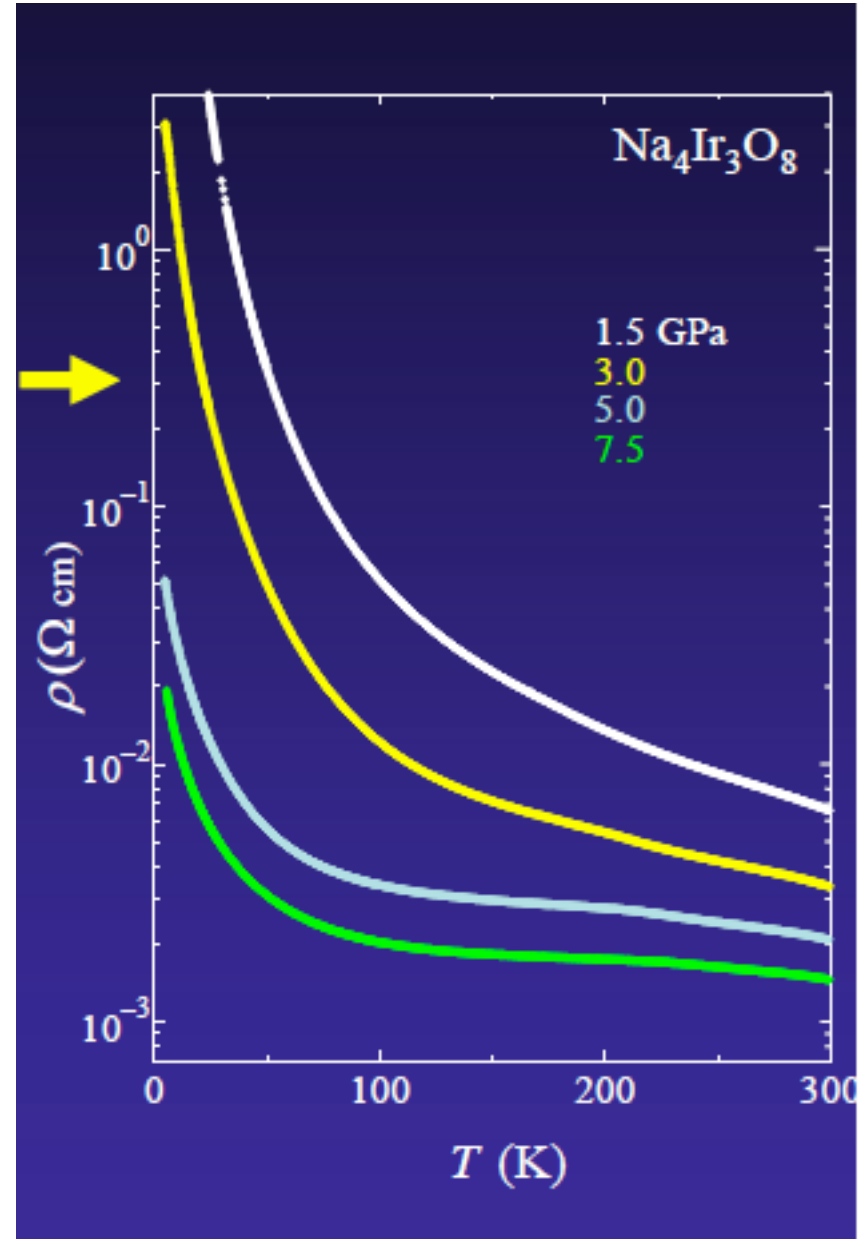
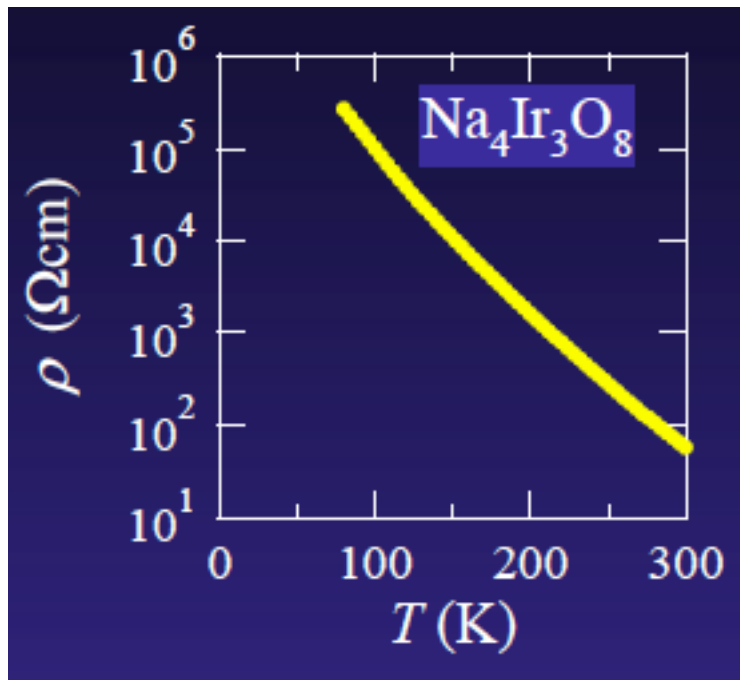
FIG. 1 (color online). (a) Crystal structure of  $\text{Na}_4\text{Ir}_3\text{O}_8$

$$\Theta_{\text{CW}} \sim -650 \text{ K}, \mu_{\text{eff}} \sim 2 \mu_{\text{B}}$$



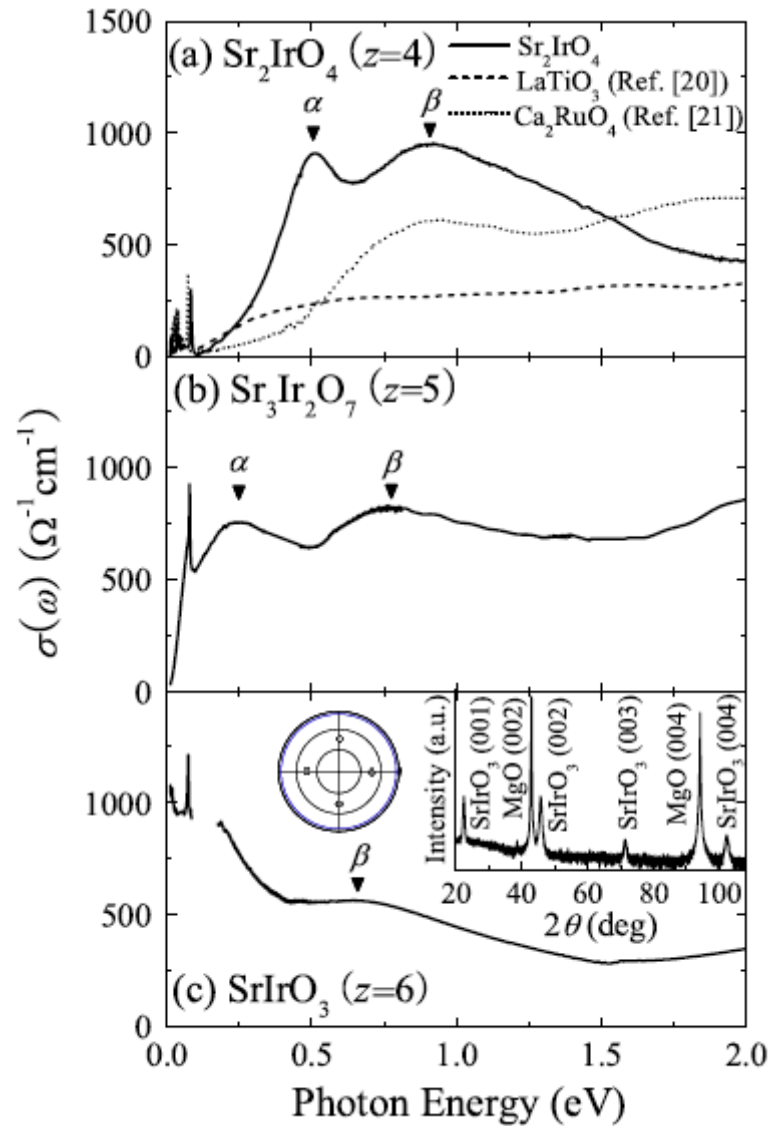
Okamoto *et al.*, PRL 99, 137207 (2007)

## Mott Insulator?



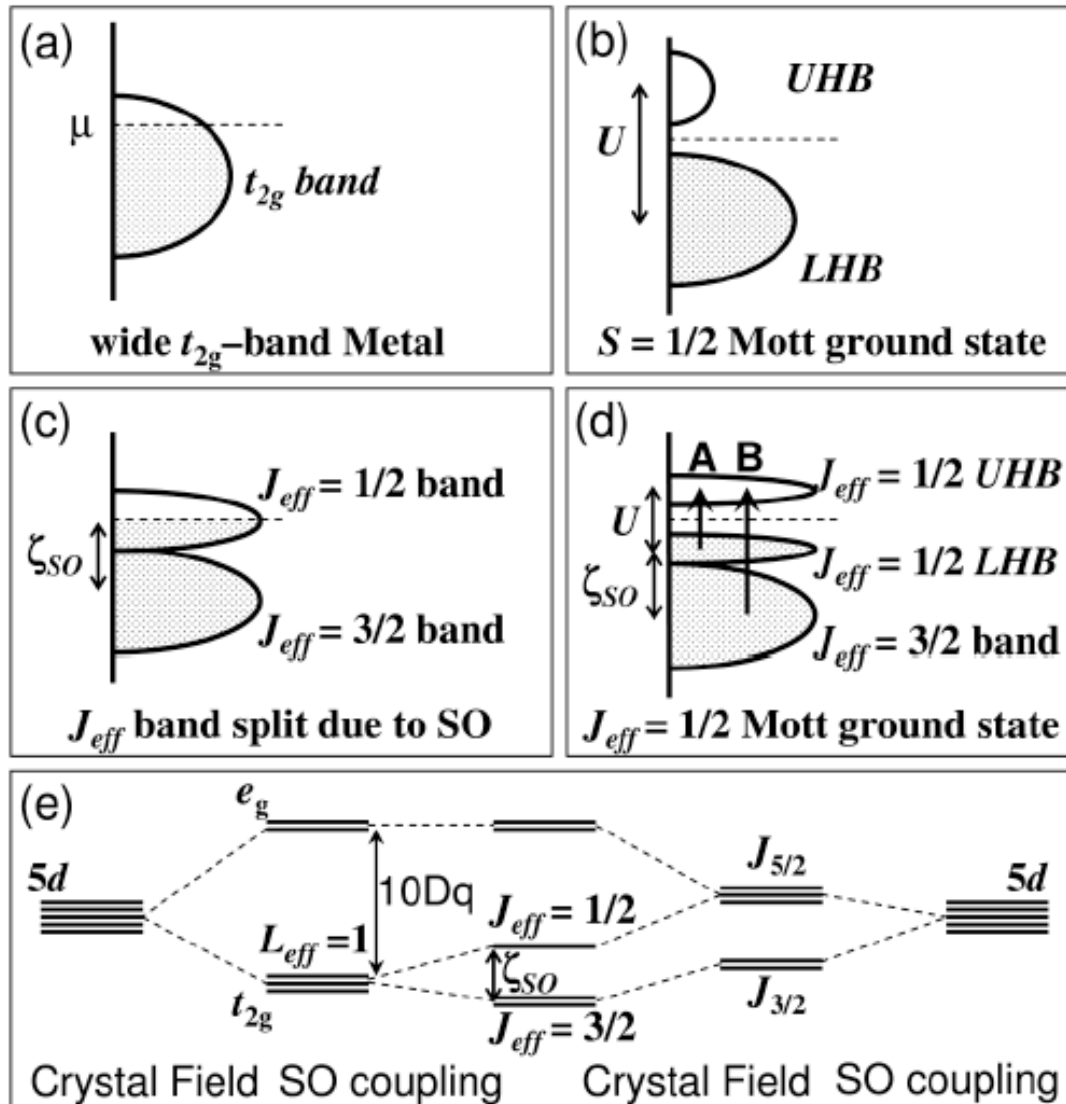
H. Takagi, unpublished

$\text{Sr}_2\text{IrO}_4$  ( $\text{La}_2\text{CuO}_4$  analogue) appears to be a Mott insulator

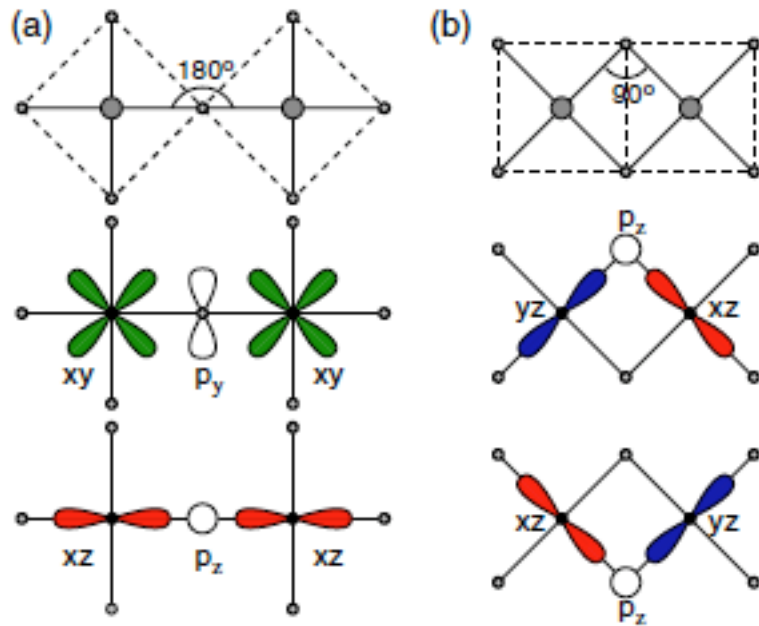


Moon *et al.*, PRL 101, 206402 (2008)

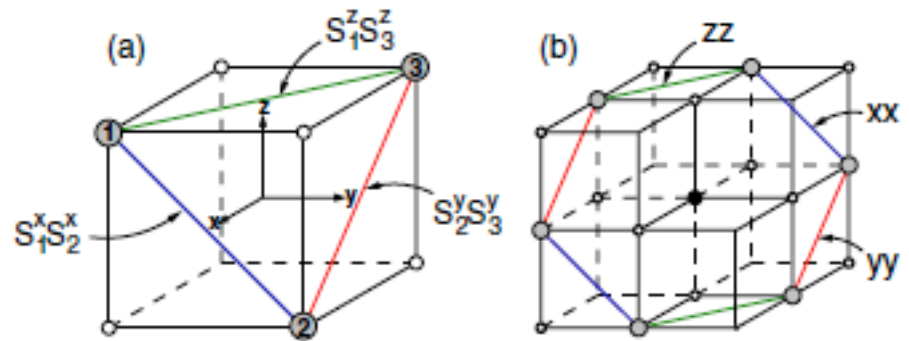
# Spin-orbit plus $d^5$ configuration leads to a half filled doublet



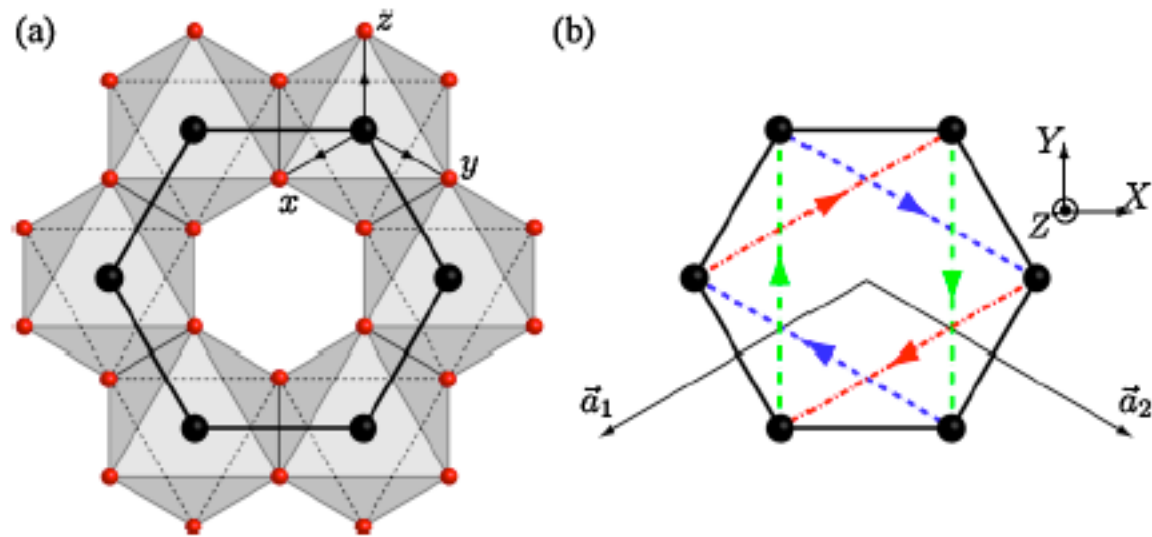
# Realization of the Kitaev Model?



$$\mathcal{H}_{ij}^{(\gamma)} = -JS_i^\gamma S_j^\gamma$$



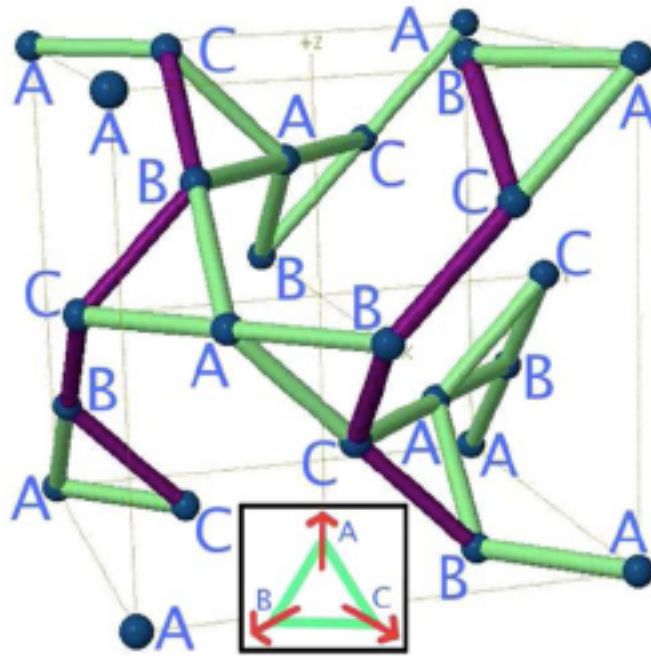
# Quantum Spin Hall Effect in $\text{Na}_2\text{IrO}_3$ ?



Shitade *et al.*, PRL 102, 256403 (2009)



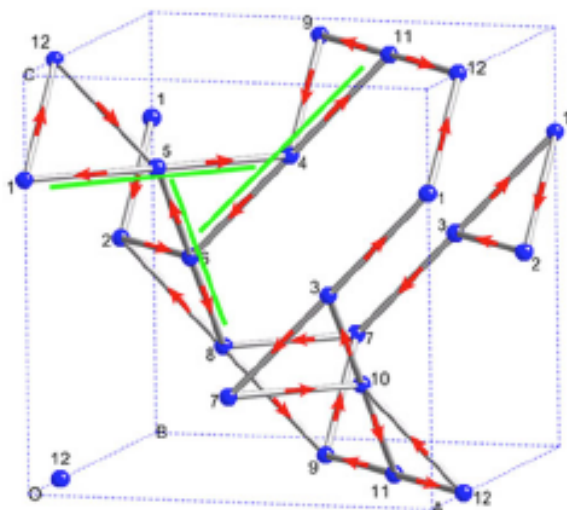
$\text{Na}_4\text{Ir}_3\text{O}_8$  - Ir ions form a network of corner sharing triangles



Semiclassical - J S·S

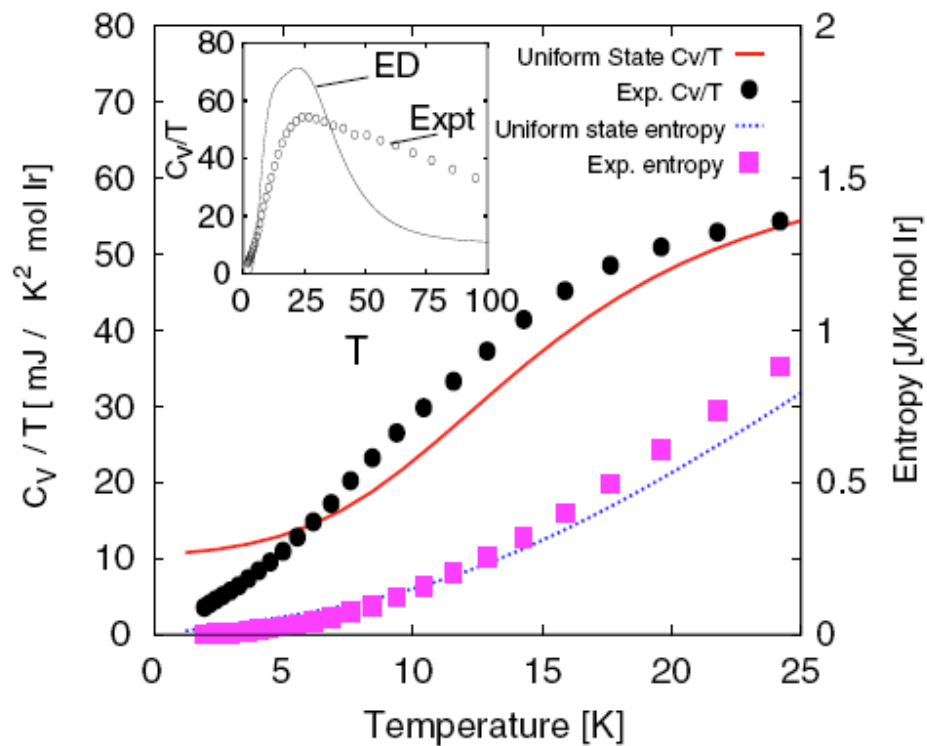
Lawler *et al.*, PRL 100, 227201 (2008)

## 10 site loops



(a)

Zhou *et al.*, PRL 101, 197201 (2008)



Lawler *et al.*, PRL 101, 197202 (2008)

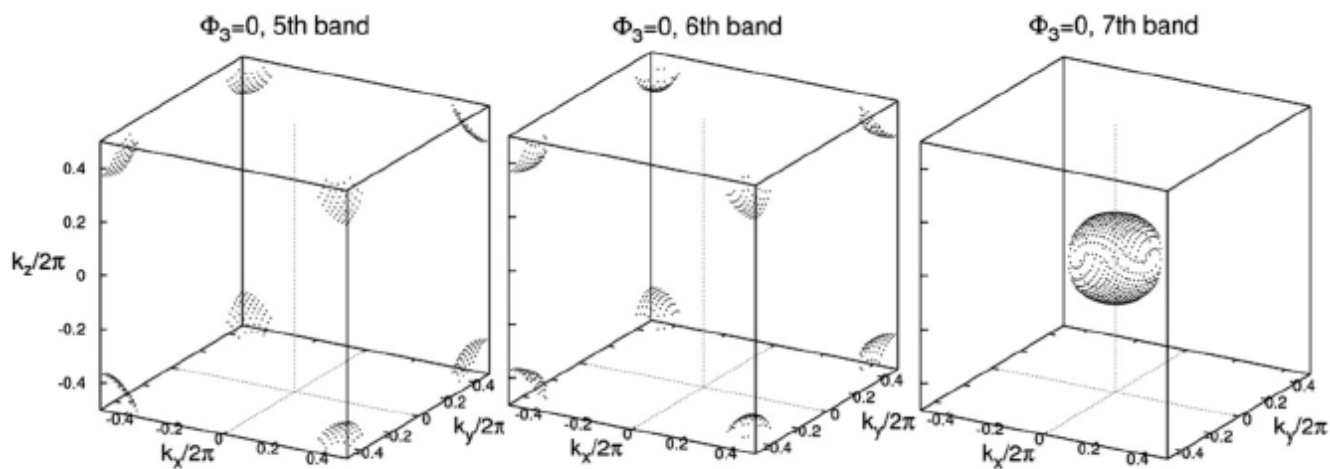


FIG. 2. Fermi surfaces of the zero flux state at half-filling.

# Pairing Instability of Spinon Fermi surface?

## Line Nodes ( $C \sim T^2$ )

1.  $\Delta_{\text{singlet}}$

2.  $\Delta = \Delta_{\text{singlet}} + \Delta_{\text{triplet}}$   
line nodes if  $\Delta_{\text{triplet}}$  is large enough

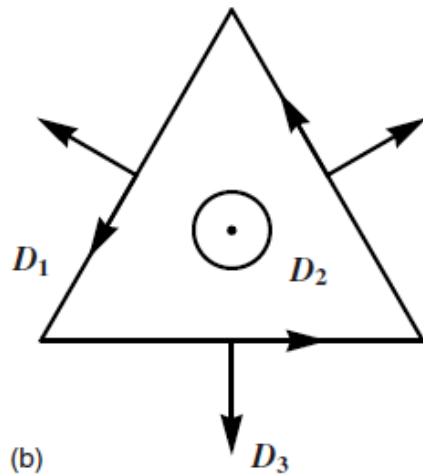
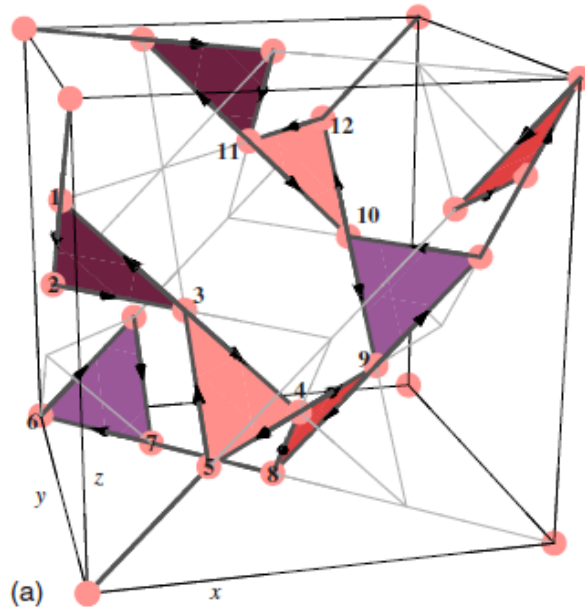
*Zhou et al.*, PRL 101, 197201 (2008)

3.  $\Delta_{\text{triplet}}$

line nodes on zone face

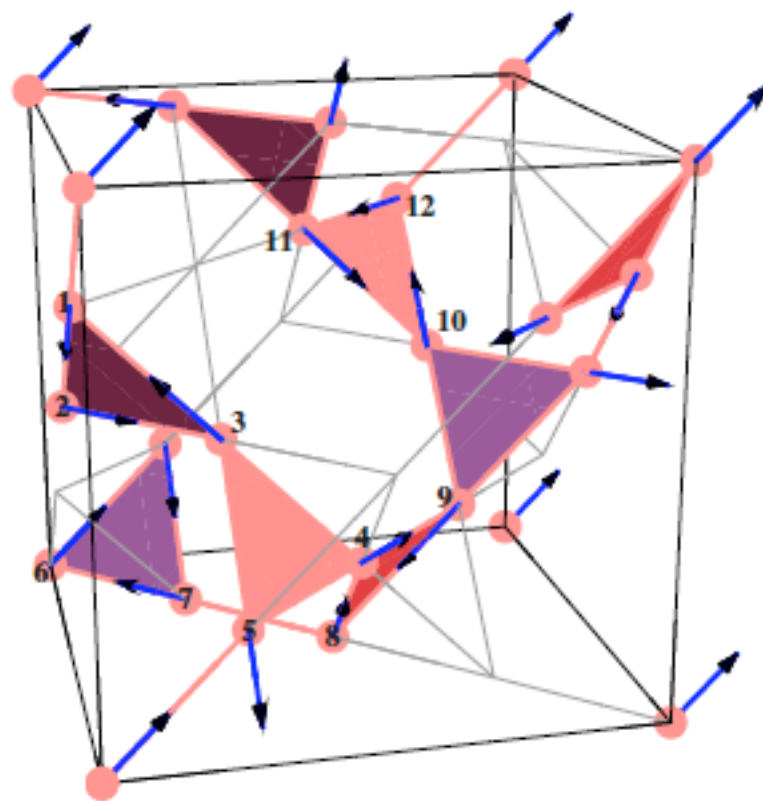
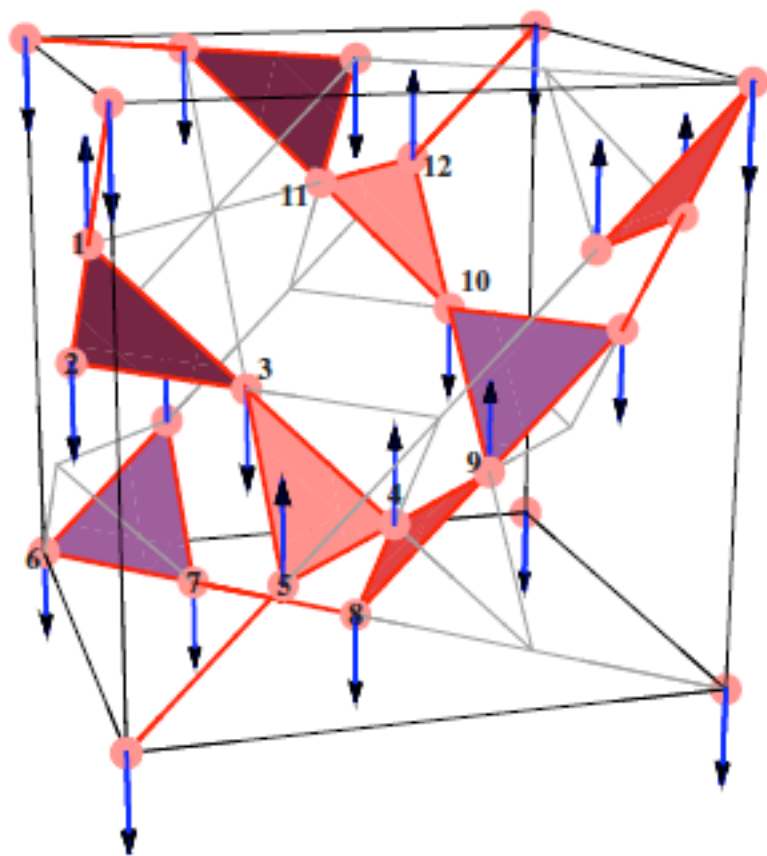
*Micklitz & Norman*, PRB 80, 100506 (2009)

# $\text{Na}_4\text{Ir}_3\text{O}_8$ - anisotropy important due to spin-orbit coupling



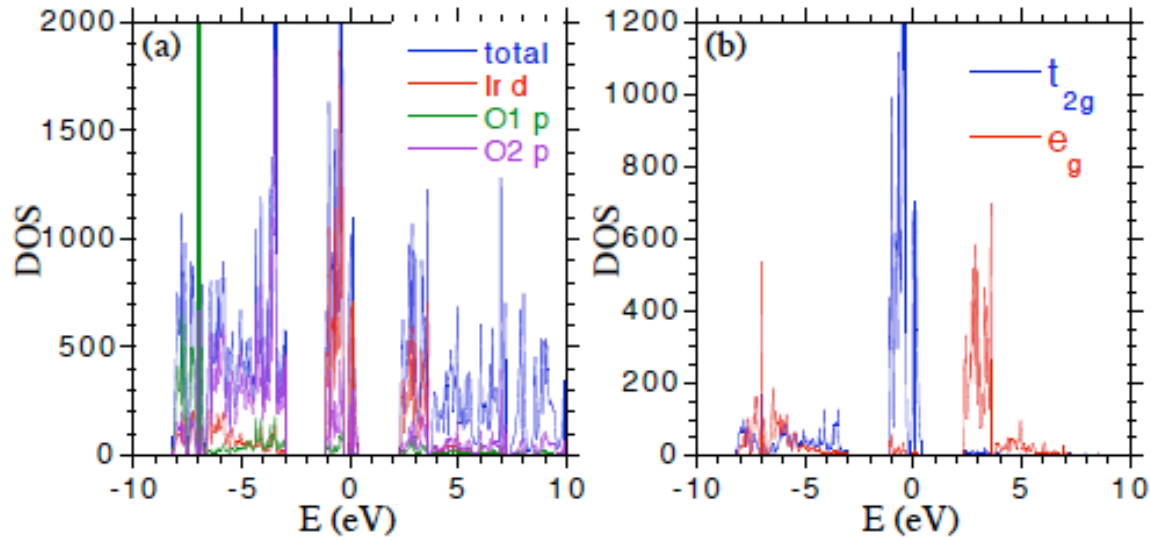
Chen & Balents, PRB 78, 094403 (2008)

# Possible ground states in the presence of anisotropy

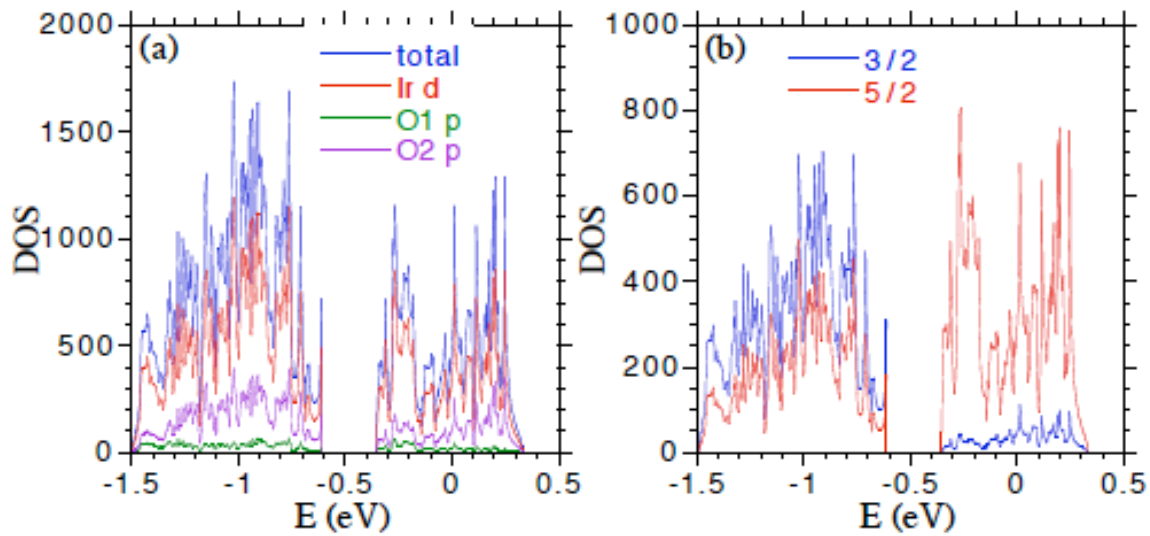


Chen & Balents, PRB 78, 094403 (2008)

# LDA

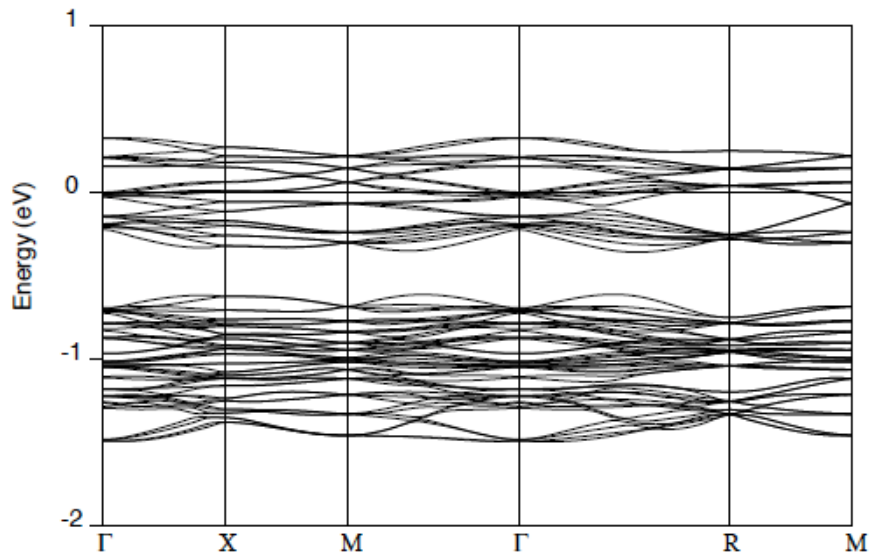


No spin-orbit



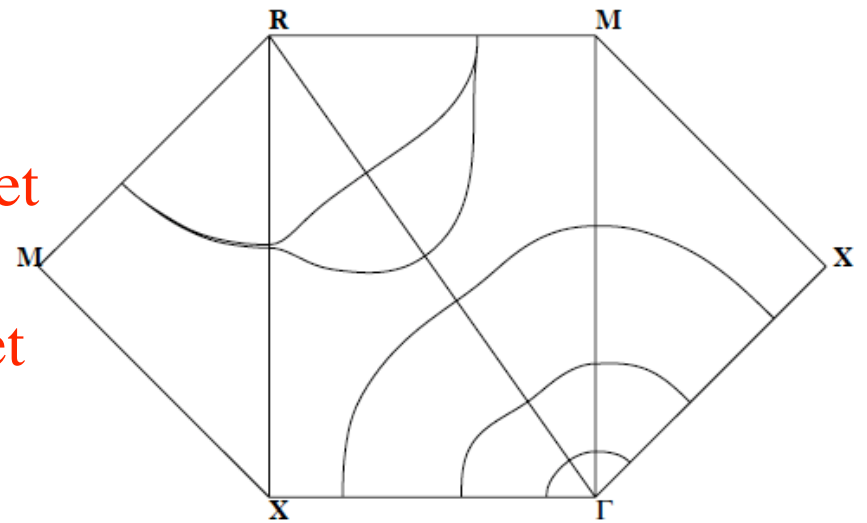
Spin-orbit

## LDA (spin-orbit)

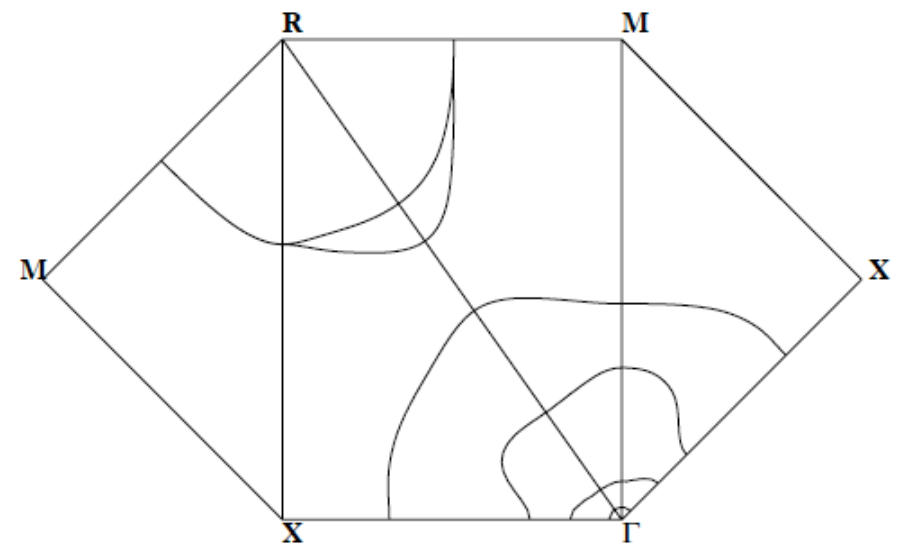
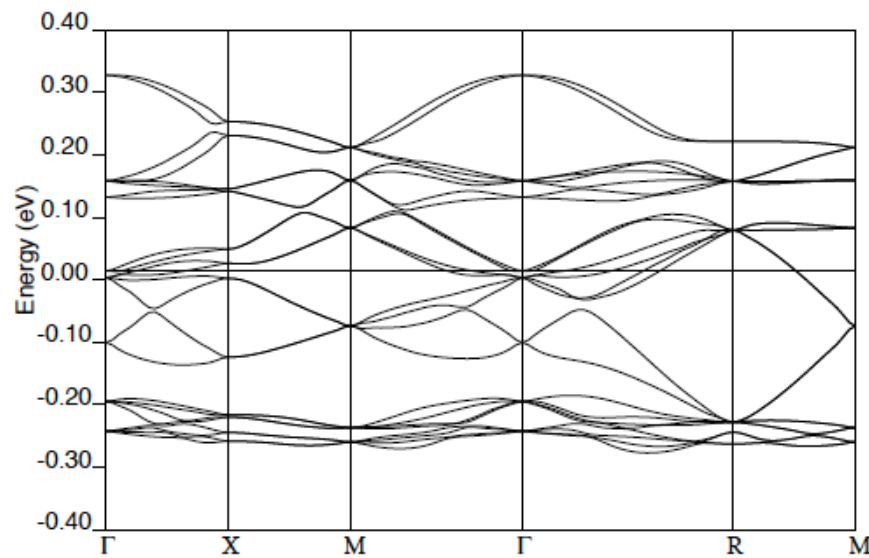


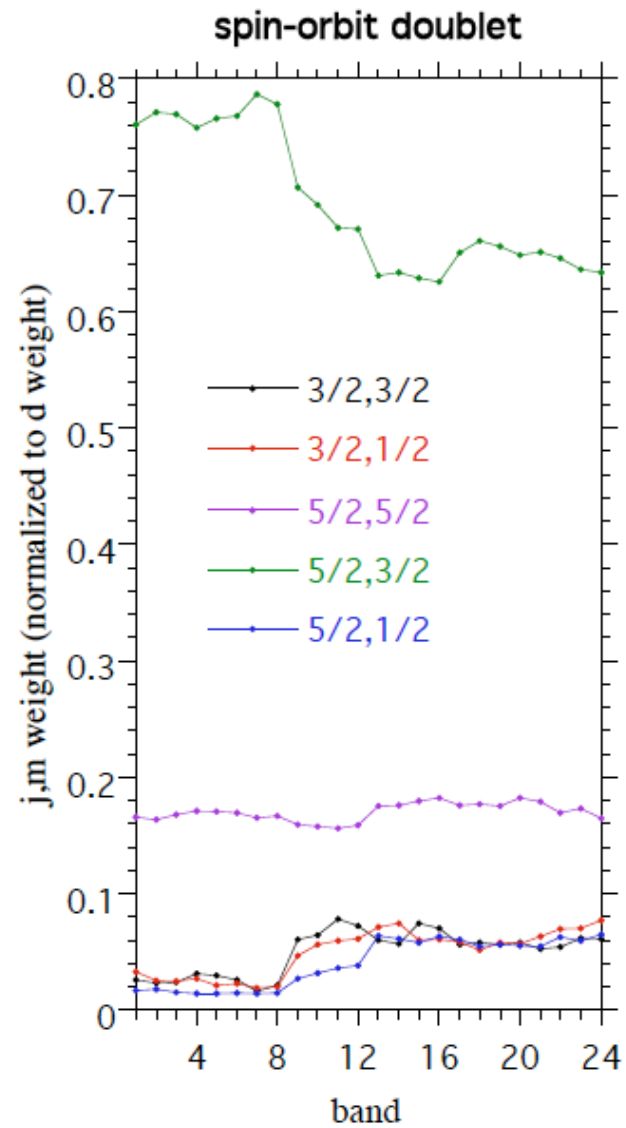
doublet

quartet



## Tight Binding Fit





Effective cubic model - 83%  $|5/2, \pm 3/2\rangle$  + 17%  $|5/2, \pm 1/2\rangle$



## Direct Ir-Ir Exchange is Isotropic!

$$J = 2 t^2/U$$

where

$$t = \frac{1}{4} t_{dd}^{\sigma} + \frac{1}{3} t_{dd}^{\pi} + \frac{5}{12} t_{dd}^{\delta}$$

But superexchange is present which should give rise to anisotropic exchange plus Dzyaloshinskii-Moriya

$$a_{i3\uparrow} = \frac{1}{\sqrt{3}}((-i)d_{i,xz\downarrow} + d_{i,yz\downarrow} + d_{i,xy\uparrow})$$

	$2p_x$	$2p_y$	$2p_z$	$5p_x$	$5p_y$	$5p_z$
$A, xz$	$t$	0	0	0	0	0
$A, yz$	0	$t$	0	0	0	$-t$
$A, xy$	0	0	0	$-t$	0	0
$B, xz$	0	0	0	$-t$	0	0
$B, yz$	0	0	$t$	0	$-t$	0
$B, xy$	$t$	0	0	0	0	0

$$\mathcal{H}_{AB} = -JS_A^x S_B^x + JS_A^y S_B^y + JS_A^z S_B^z,$$

## Tight Binding Fit (18 parameters)

Tight binding hopping parameters in eV

The on-site energies are:

$$\epsilon_{O1} = -5.1913, \epsilon_{O2} = -4.0954, \epsilon_{t_{2g}} = -1.7983, \epsilon_{e_g} = 0.7987$$

The spin-orbit coupling  $\lambda = 0.6386$

	$\sigma$	$\pi$	$\delta$
Ir-O1	-2.5579	0.3186	
Ir-O2	-2.1070	1.1817	
Ir-Ir	-0.6372	0.0719	0.1545
O1-O1	0.6156	0.0354	
O2-O2	0.5463	-0.2121	
O1-O2	-1.1327	0.0663	

## Anisotropic Exchange

Exchange integrals in meV ( $U=0.5$  eV,  $U_p=0$ )

$$H_{ex} = (J_d + J_s)\mathbf{S}_n \cdot \mathbf{S}_m + \mathbf{D}^{nm} \cdot (\mathbf{S}_n \times \mathbf{S}_m) + \mathbf{S}_n \cdot \mathbf{\Gamma}^{nm} \cdot \mathbf{S}_m$$

Ir site  $m$  is along (0,1,-1) relative to Ir site  $n$

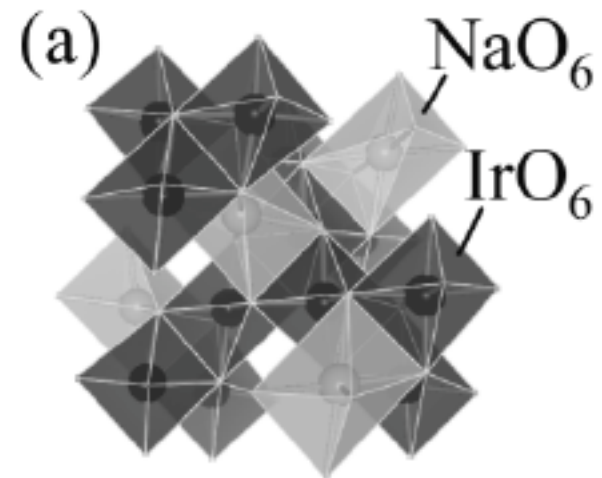
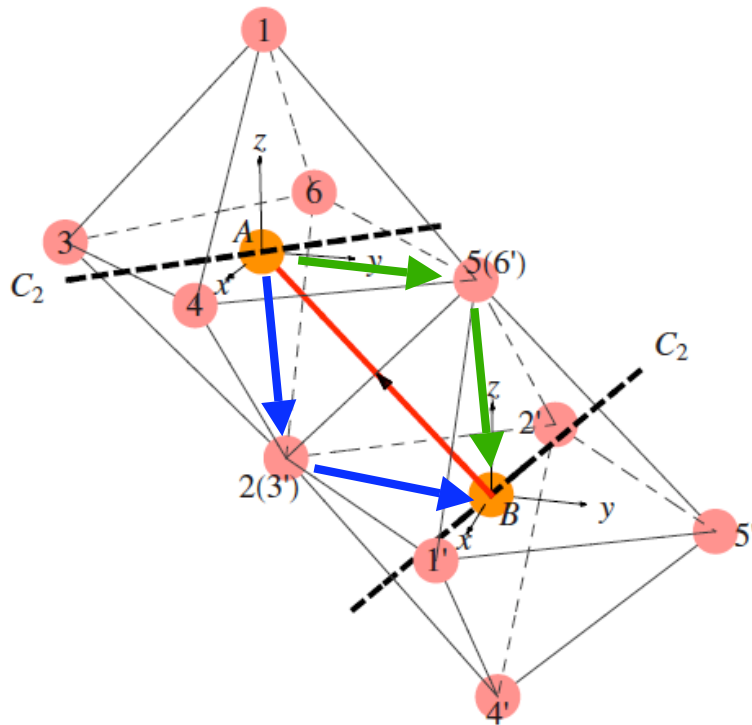
$$J_d = 20.1 \text{ meV}, J_s = 12.9 \text{ meV}$$

$$J_{ii} = J_d + J_s + \Gamma_{ii}$$

$i$ ( $jk$ )	$x$ ( $xy$ )	$y$ ( $xz$ )	$z$ ( $yz$ )
$D_i$	47.3	1.3	-4.7
$\Gamma_{ii}$	36.1	-36.8	-36.2
$\Gamma_{jk}$	2.1	-8.3	-0.2
$J_{ii}$	69.2	-3.7	-3.1

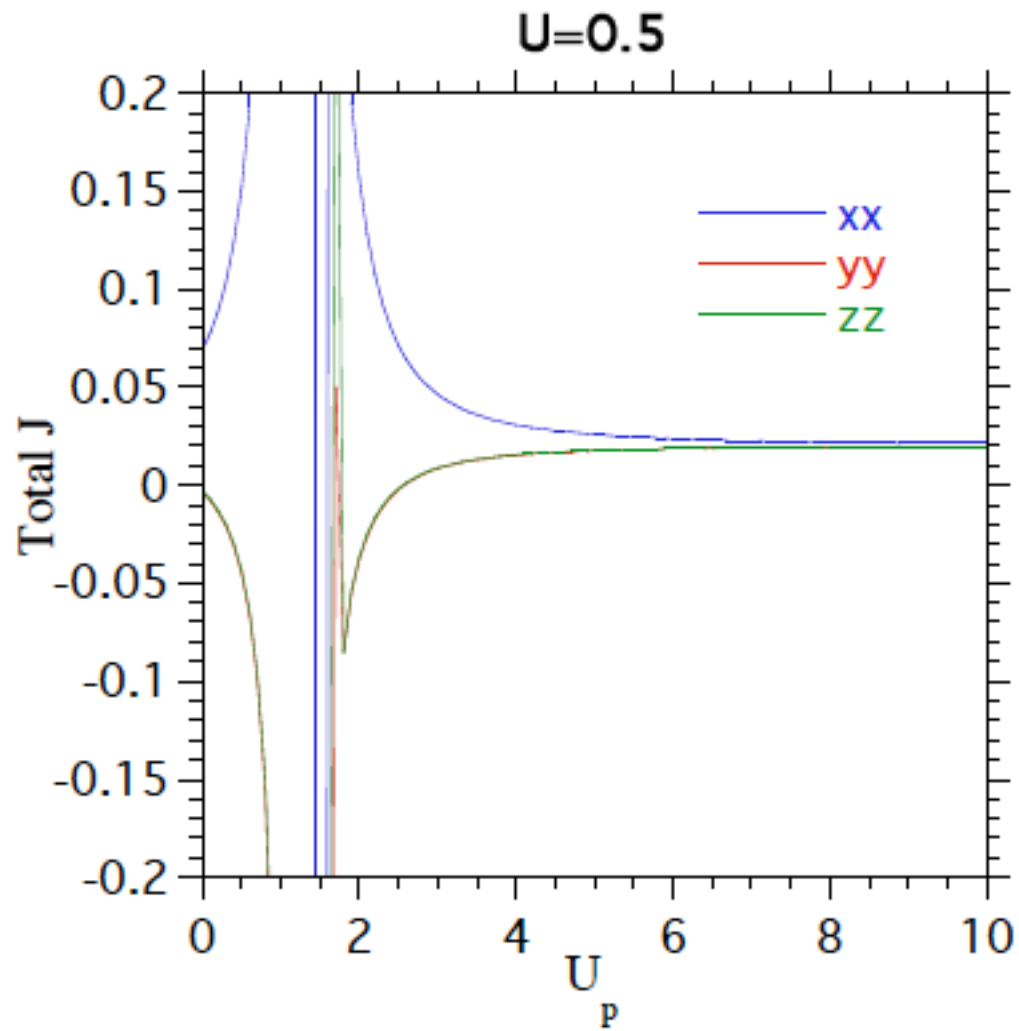
# How to Get Rid of the Superexchange

Dephasing due to partial Na site occupation?



# How to Get Rid of the Superexchange

Large  $U_p$ ?



# How to Detect a Spinon Fermi Surface

Norman & Micklitz, PRL 102, 067204 (2009)

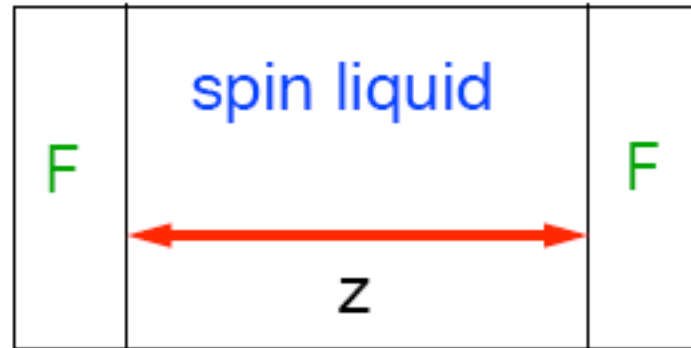


FIG. 1: (Color online) The proposed experiment involves two ferromagnetic layers (F) with a spin liquid spacer of variable thickness,  $z$ . Depending on the sign of the oscillatory coupling, the two ferromagnets will be aligned or anti-aligned.

Parkin *et al.*, PRL 64, 2304 (1990)

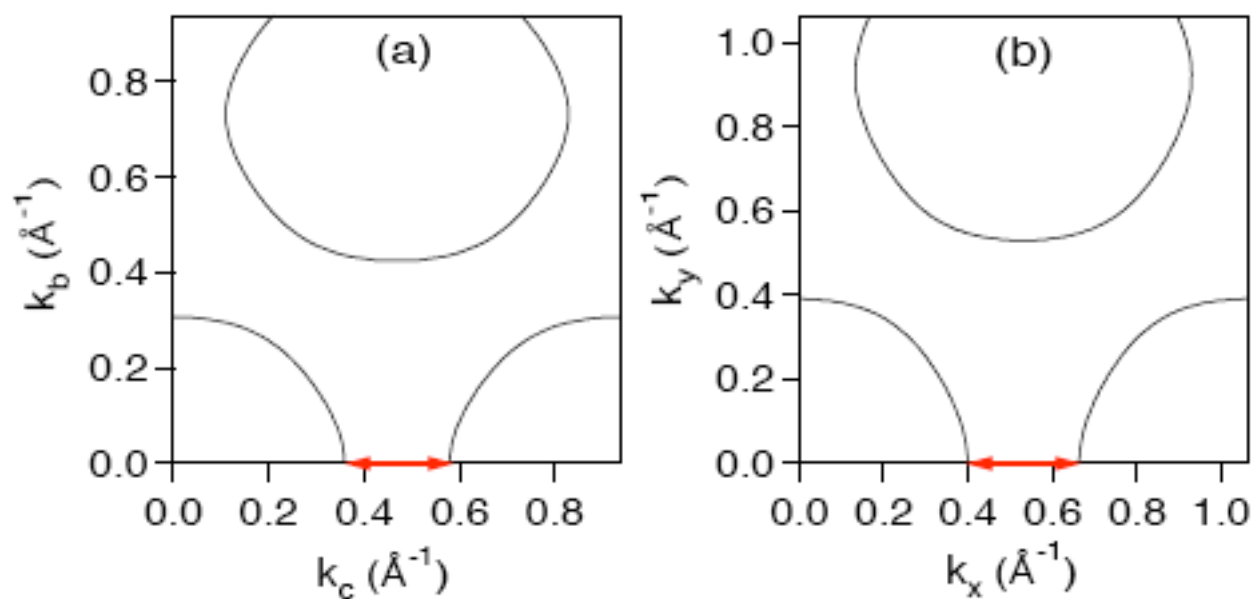


FIG. 2: (Color online) Spinon Fermi surface for (a)  $\kappa$ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$  ( $b=8.59\text{\AA}$ ,  $c=13.40\text{\AA}$ ) and (b) ZnCu $_3$ (OH) $_6$ Cl $_2$  ( $a=6.84\text{\AA}$ ). Spanning vectors are indicated by arrows.



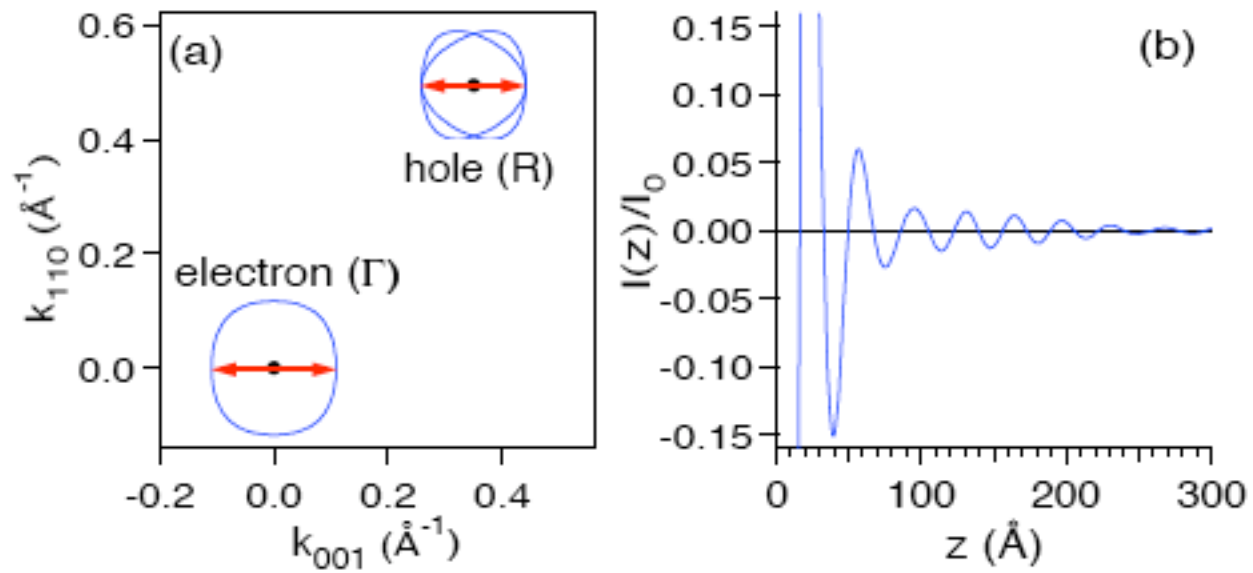


FIG. 3: (Color online) (a) Spinon Fermi surface for  $\text{Na}_4\text{Ir}_3\text{O}_8$  ( $a=8.985\text{\AA}$ ). Spanning vectors are indicated by arrows. (b) Calculated oscillatory response from Eq. 7.

$$I(z) = I_0 \left(\frac{d}{z}\right)^2 \sum_n \frac{m_n^*}{m} \sin(2k_{F_n} z)$$

Cuprates are Mott insulators characterized by  
a half filled band

So are iridates!

Can iridates be doped, and if so, will they  
become high temperature superconductors?