

Wavelength Selection in the Noisy Stabilized Kuramoto-Sivashinskii Equation

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INTRODUCTION:

Problem is to find stationary state of driven out of equilibrium system.

- System evolves via Langevin dynamics

$$\frac{\partial \psi(x, t)}{\partial t} = K[\psi] + \eta(x, t) \quad \langle \eta(x, t) \eta(x', t') \rangle = 2\epsilon \delta(x - x') \delta(t - t')$$

Equivalent to Fokker-Planck equation:

$$\frac{\partial P(\psi(x), t)}{\partial t} = -\frac{\delta}{\delta \psi} \left(K[\psi] P - \epsilon \frac{\delta P}{\delta \psi} \right)$$

Stationary state:

$$\epsilon \frac{\delta P}{\delta \psi} = K[\psi] P \Rightarrow P_0(\psi_0 \rightarrow \psi_f) \propto \exp \left(\frac{1}{\epsilon} \int_{\psi_0}^{\psi_f} \mathcal{D}\psi K[\psi] \right)$$

- Assume that free energy potential exists and $K[\psi]$ is derivative of the potential:

$$K[\psi] = -\frac{\delta\mathcal{F}[\psi]}{\delta\psi} \Rightarrow P_0[\psi_0 \rightarrow \psi] \propto \exp\left(-\frac{1}{\epsilon}(\mathcal{F}[\psi] - \mathcal{F}[\psi_0])\right)$$

- Boltzmann form for stationary distribution:

$$P[\psi] = \frac{1}{Z} \exp\left(-\frac{\mathcal{F}[\psi]}{\epsilon}\right) \quad Z = \int \mathcal{D}\psi \exp\left(-\frac{\mathcal{F}[\psi]}{\epsilon}\right)$$

- Ingredients of $K[\psi]$: Set of stationary state solutions to $K[\psi]=0$.
 $K[\psi]$ must be non-linear (insoluble) and not derivative of potential.

Hypothesis: Additive stochastic noise distinguishes between stationary states and selects one or a band whose width vanishes in the thermodynamic limit.

Methods: Numerical simulations and stability analysis of stationary states – most stable state is selected??

Simplest model we could think of is stabilized Kuramoto-Sivashinsky equation in 1D with additive stochastic noise:

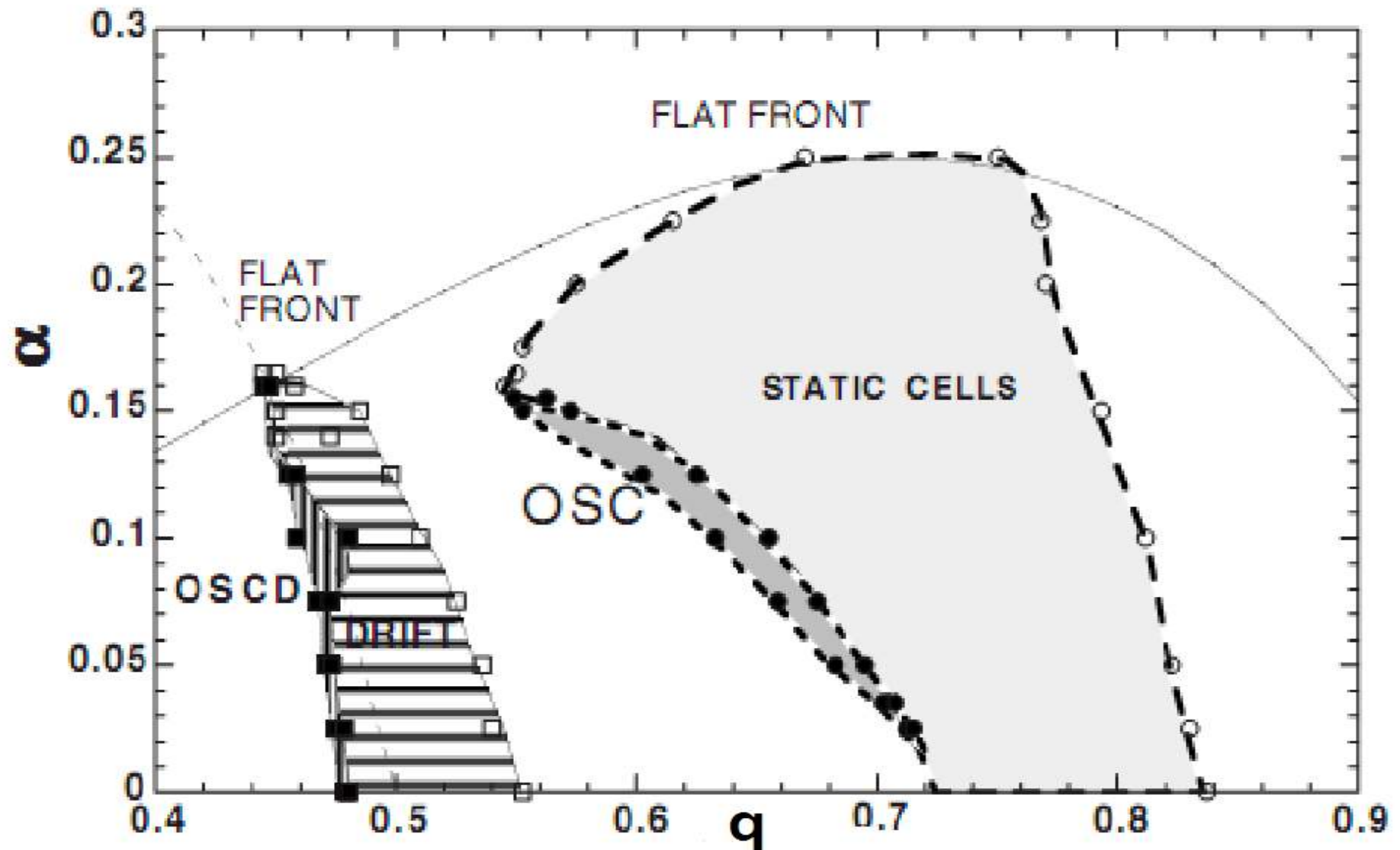
$$\frac{\partial h(x, t)}{\partial t} = - \left(\alpha + \frac{\partial^2}{\partial x^2} + \frac{\partial^4}{\partial x^4} \right) h(x, t) + \left(\frac{\partial h}{\partial x} \right)^2 + \eta(x, t)$$

$$\dot{h}(q, t) = - \left(\alpha - \frac{1}{4} + \left(q^2 - \frac{1}{2} \right)^2 \right) h(q, t) - \int \frac{dk}{2\pi} k(q-k) h(k) h(q-k) + \eta(q, t)$$

Modes are linearly stable if $\alpha > 1/4$ and linearly unstable if $\alpha < 1/4$.

$$1/2 + \sqrt{1/4 - \alpha} \geq q^2 \geq 1/2 - \sqrt{1/4 - \alpha}$$

- Stability diagram of deterministic SKS equation.
- P. Brunet, Phys. Rev. E 76, 017204 (2007)



Numerical simulations:

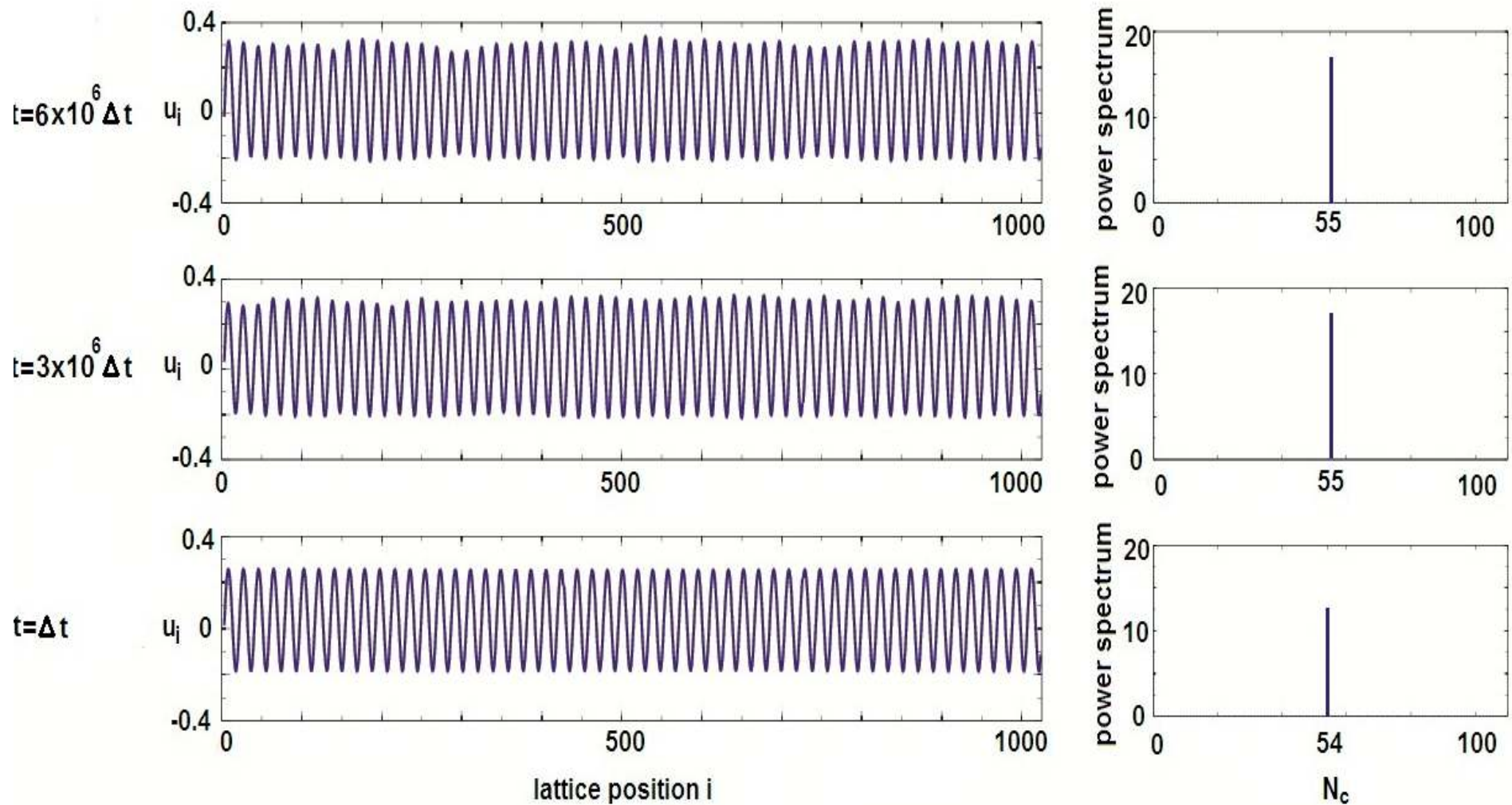
$$h_i^n = h(x_i, t_n) \quad t_n = n\Delta t \quad x_i = i\Delta x \quad 1 \leq i \leq N \quad h_{i+N}^n = h_i^n$$

$$h_i^{n+1} = h_i^n + \Delta t K_i^n[h] + \sqrt{\frac{2\epsilon\Delta t}{\Delta x}} \eta_i^n \quad \langle \eta_i^n \rangle = 0 \quad \langle \eta_i^n \eta_j^m \rangle = \delta_{ij} \delta_{nm}$$

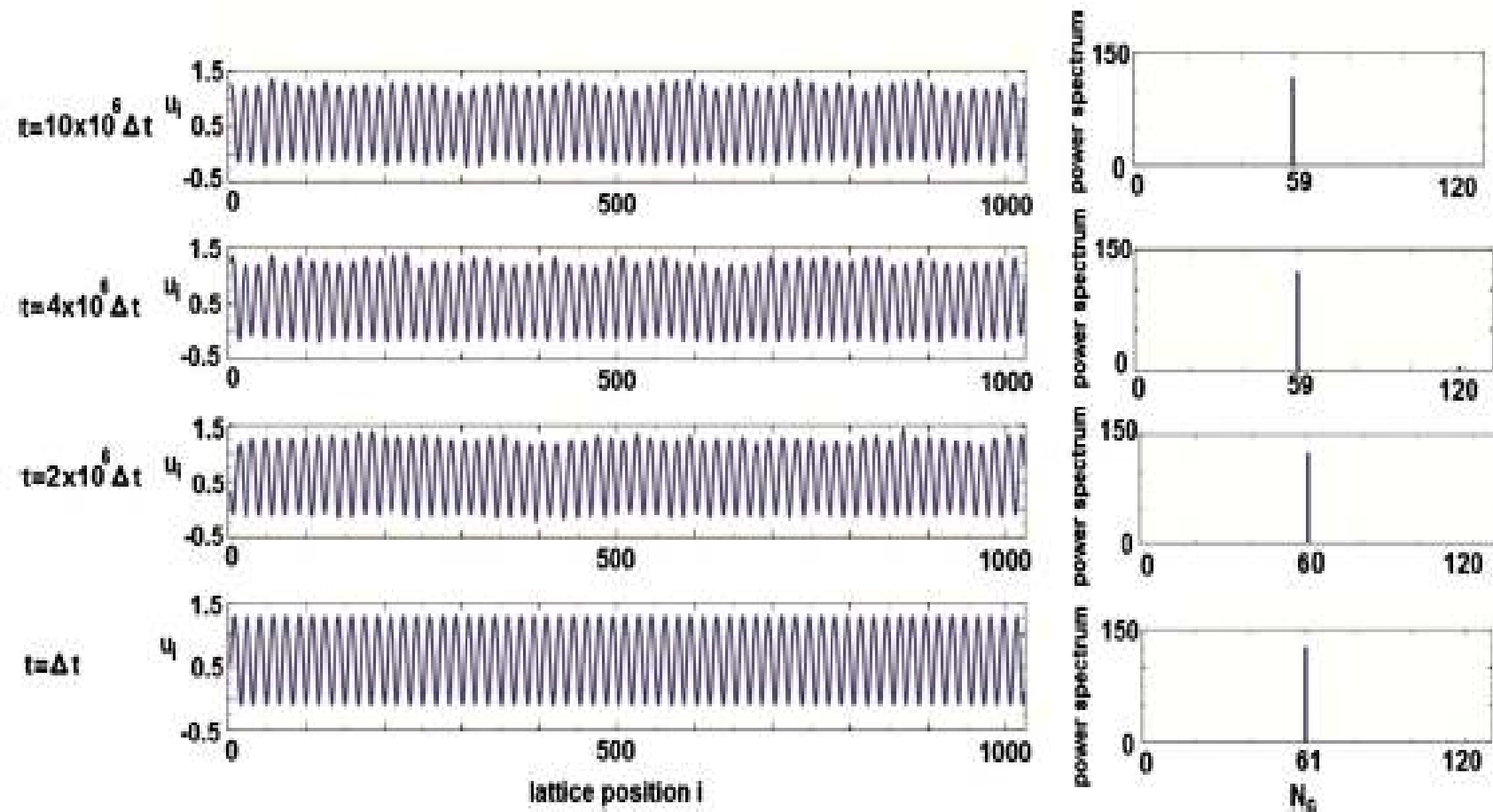
$$K_i^n[h] = -\alpha h_i^n - \frac{1}{(\Delta x)^2} (h_{i+1}^n - 2h_i^n + h_{i-1}^n) - \frac{1}{(\Delta x)^4} (h_{i+2}^n - 4h_{i+1}^n + 6h_i^n - 4h_{i-1}^n + h_{i-2}^n) \\ + \frac{1}{4(\Delta x)^2} (h_{i+1}^n - h_{i-1}^n)^2$$

$$q = \frac{2\pi N_c}{N\Delta x} \quad N = 1024 \quad \Delta x = \frac{1}{2} \quad \Delta t \leq C(\Delta x)^4 \quad \Delta t = 0.006$$

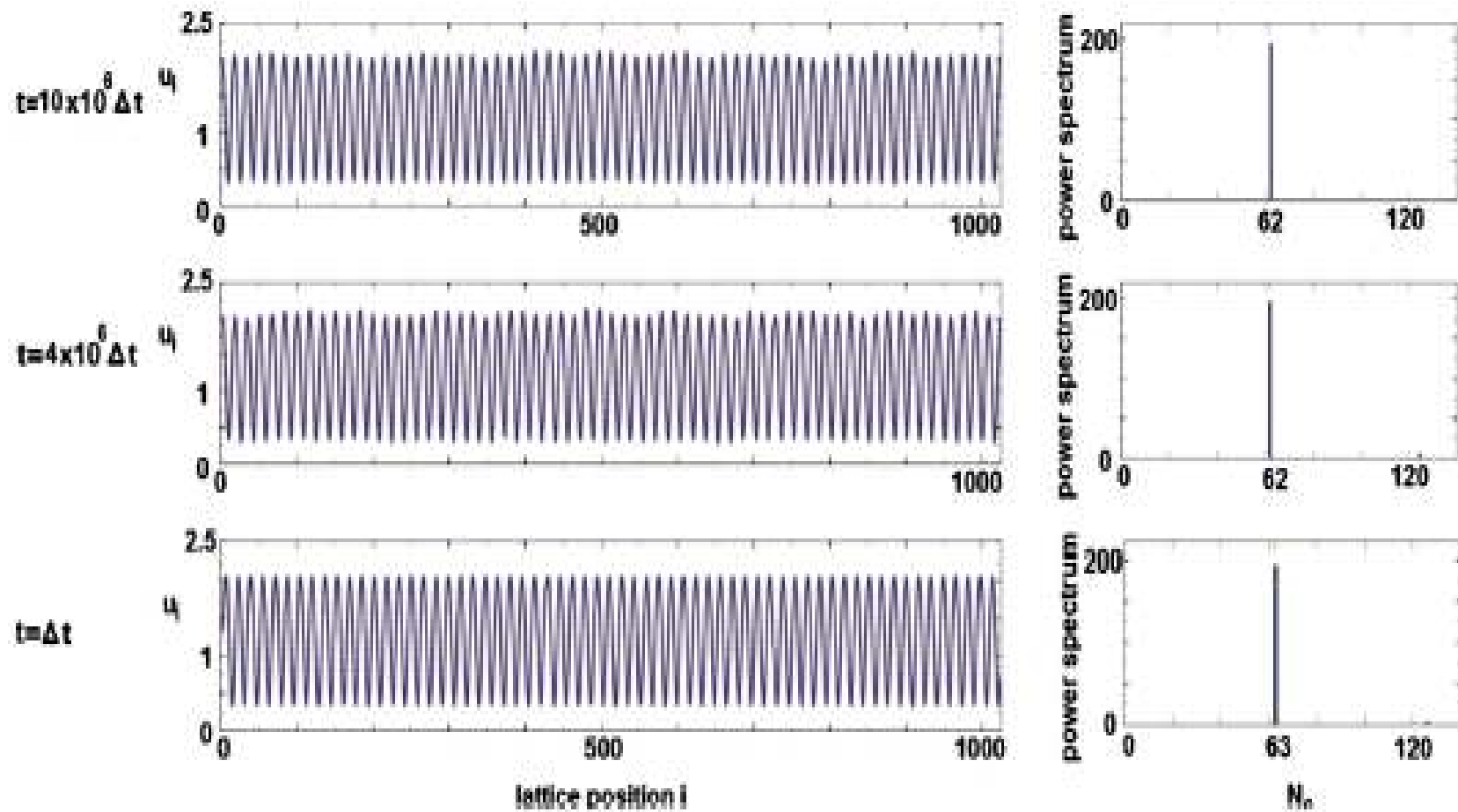
- $\alpha=0.24$ Eckhaus stable band without noise
 $53 < N_c < 61$. Simulation done with noise $\varepsilon=0.00001$.
 Initial state by evolving deterministic equation.



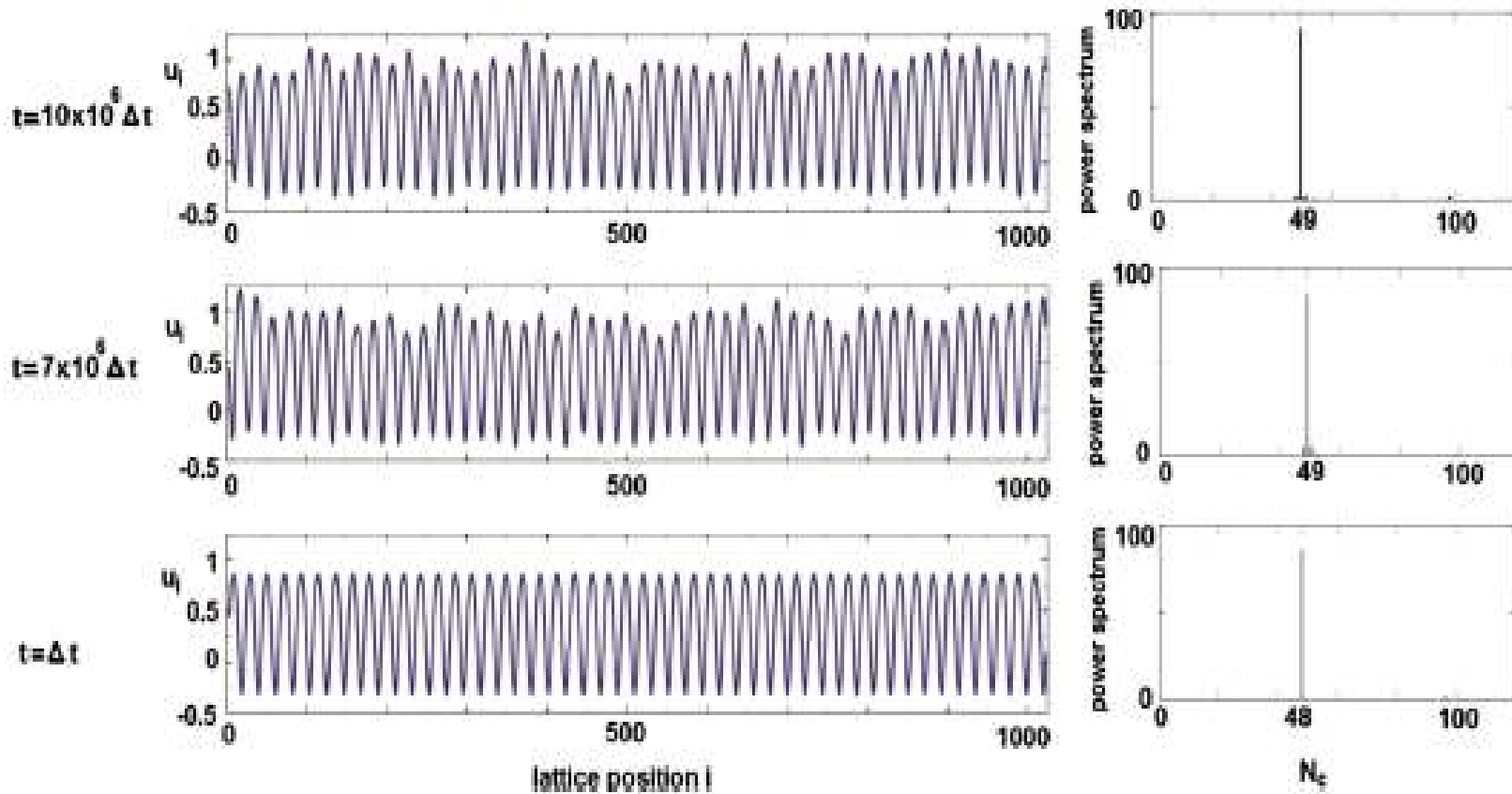
Time evolution of an $N_c=61$ state, $\alpha=0.20$, noise strength $\varepsilon=0.0005$



Evolution of $N_c = 63$ state with noise $\varepsilon = 0.0005$, $\alpha = 0.17$



Evolution of $N_c = 48$ state, $\alpha = 0.17$, $\varepsilon = 0.0005$



Phase diffusion stability: need eigenvalues for perturbations about periodic steady state $H(x)$:

$$h(x, t) = H(x) + v(x, t) \quad H(x + L) = H(x) \quad v(x + L) = v(x)$$

$$-(\alpha + \partial_x^2 + \partial_x^4) H + (\partial_x H)^2 = 0 \quad v_t = [-\alpha - \partial_x^2 - \partial_x^4 + 2(\partial_x H)\partial_x]v \equiv \mathcal{L}v$$

Eigenvalue equation: $[\alpha + (\partial_x + \nu)^2 + (\partial_x + \nu)^4 - 2(\partial_x H)(\partial_x + \nu) + \lambda]v = 0$

$$\lambda^0(\nu) = \lambda^0(0) + \nu\lambda_\nu^0(0) + \frac{1}{2}\nu^2\lambda_{\nu\nu}^0(0) \dots$$

We find eigenvalues with "Auto" software package for system of first order ODE.

$$H \equiv (h, h_x, h_{xx}, h_{xxx}) \quad H_x = LF(H, c)H \quad \text{where} \quad H(x + 1) = H(x)$$

$$F(H, c) = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -\alpha & (h_x - c) & -1 & 0 \end{pmatrix}$$

Linearize about steady state $H(x)$ and expand to $O(v^2)$ and solve resulting eigenvalue problems with Auto:

$$H \rightarrow H + V \quad V = (v, v_x, v_{xx}, v_{xxx}) \quad V_x = L[A + \lambda B - \nu]V$$

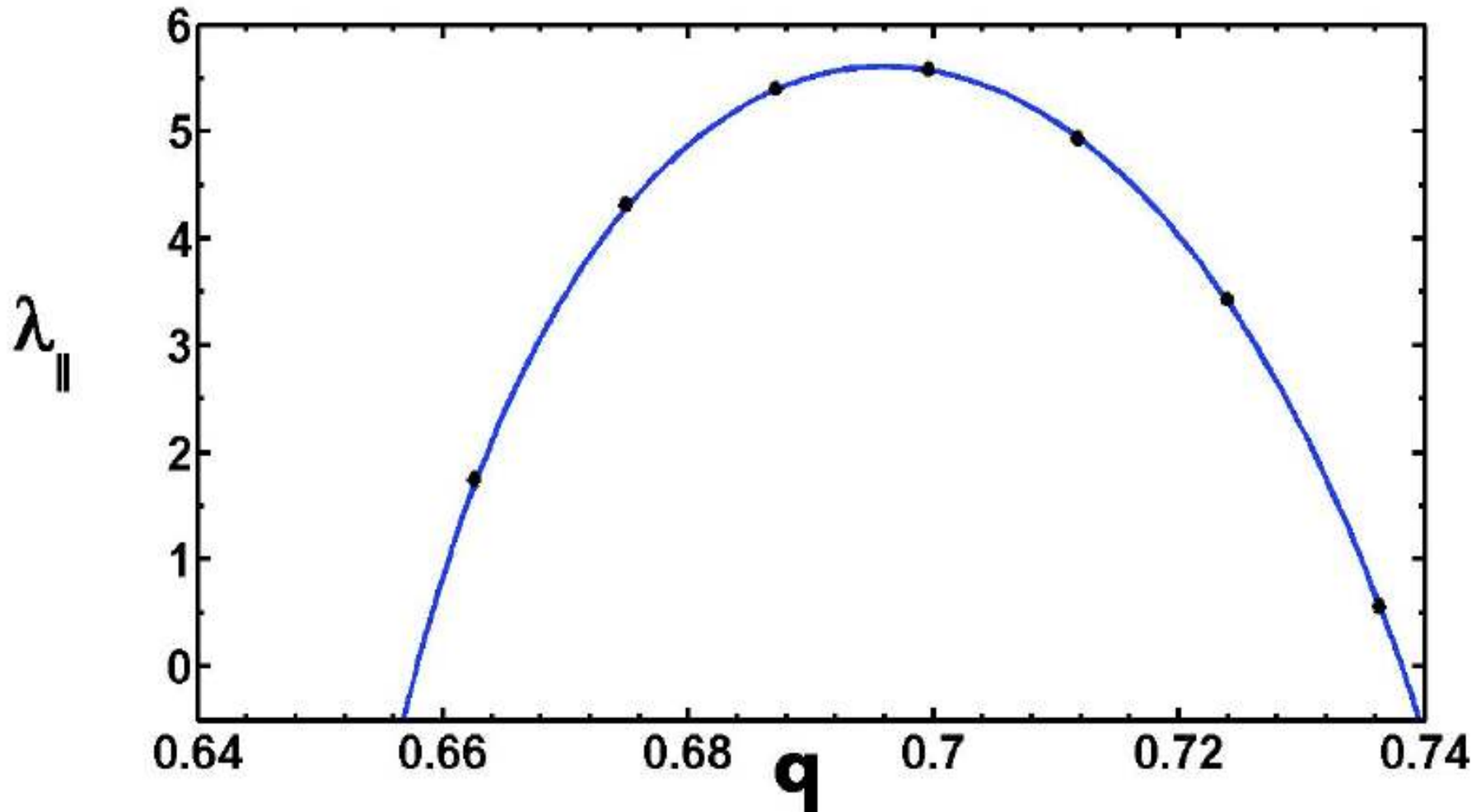
$$V'_1 = L[AV_1 + (\lambda_1 B - 1)V] \quad \lambda_1 \equiv \left. \frac{d\lambda^0}{d\nu} \right|_{\nu=0}$$

$$V'_{\parallel} = L[AV_{\parallel} + 2(\lambda_1 B - 1)V_1 + \lambda_{\parallel} BV] \quad \lambda_{\parallel} \equiv \left. \frac{d^2\lambda^0}{d\nu^2} \right|_{\nu=0}$$

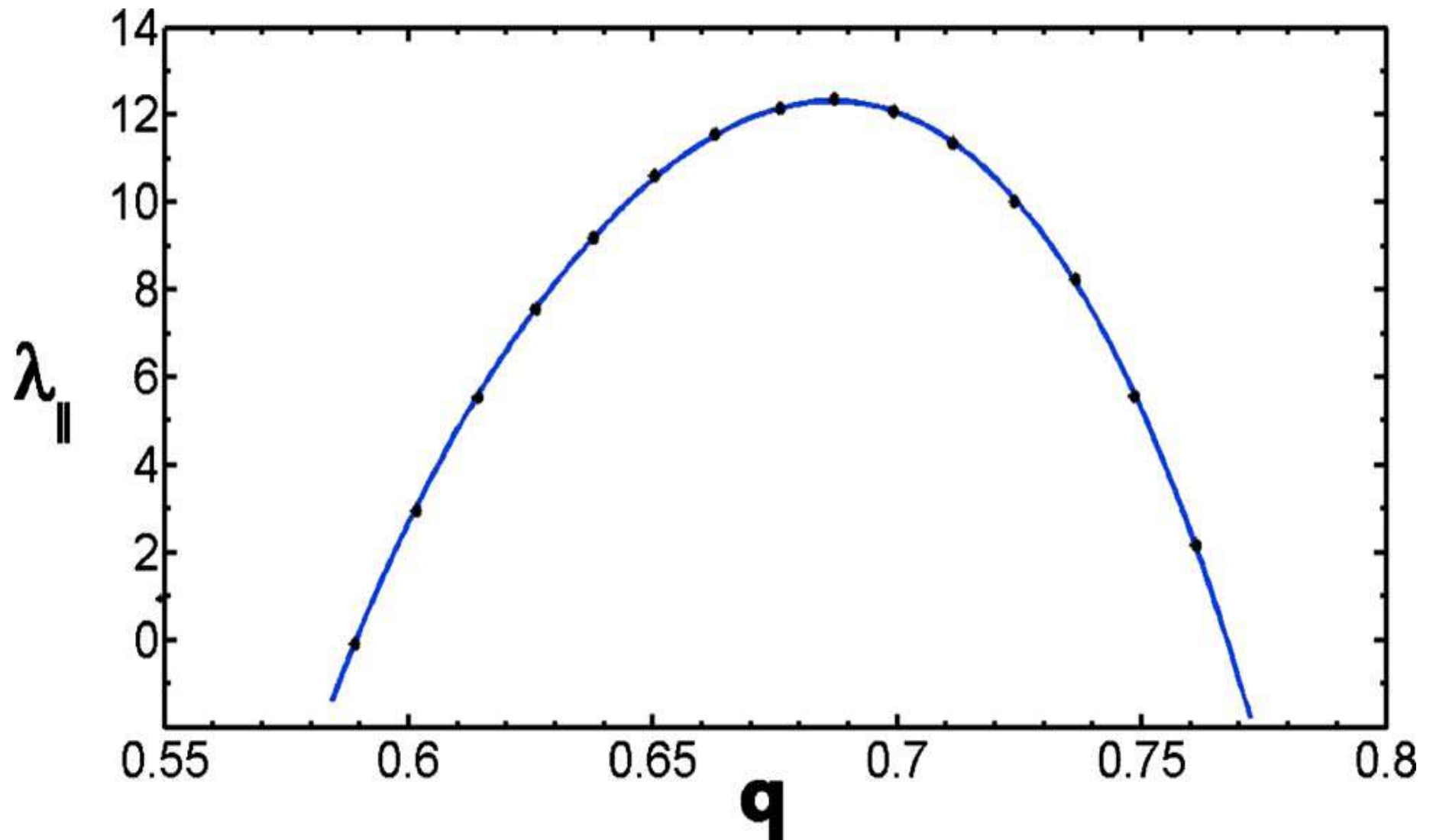
$$A = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -\alpha & (2H_x - c) & -1 & 0 \end{pmatrix}$$

$$B = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}$$

$\lambda_{\parallel}(q)=2D(q)$ for $\alpha=0.24$. Dots are $q=2\pi N_c/N\Delta x$ of simulation with $N=1024$, $\Delta x=0.50$. $\lambda_{\parallel}(\max)=5.5$

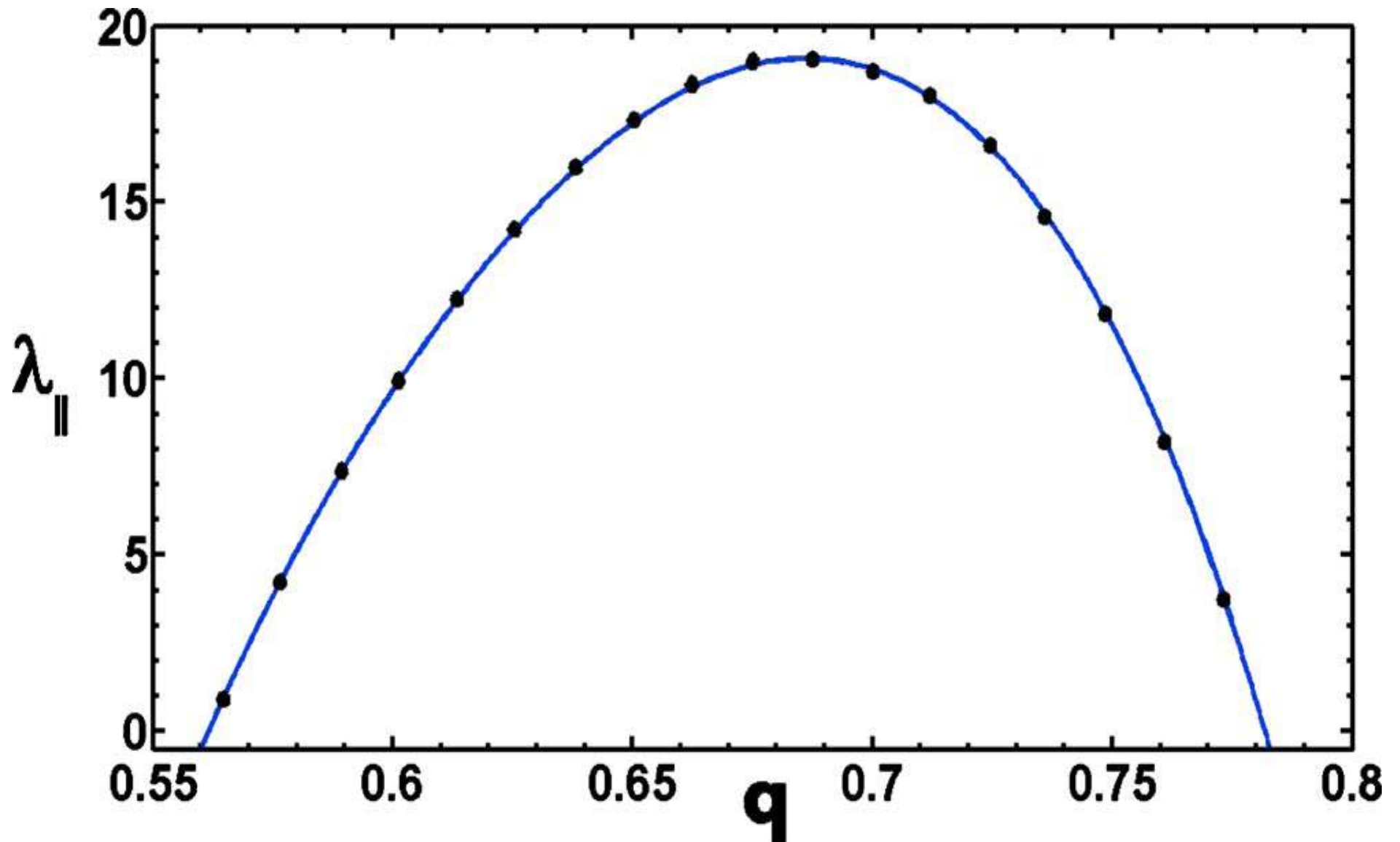


D(q) for $\alpha=0.20$. N $\lambda_{\parallel}(\max)=12.5$

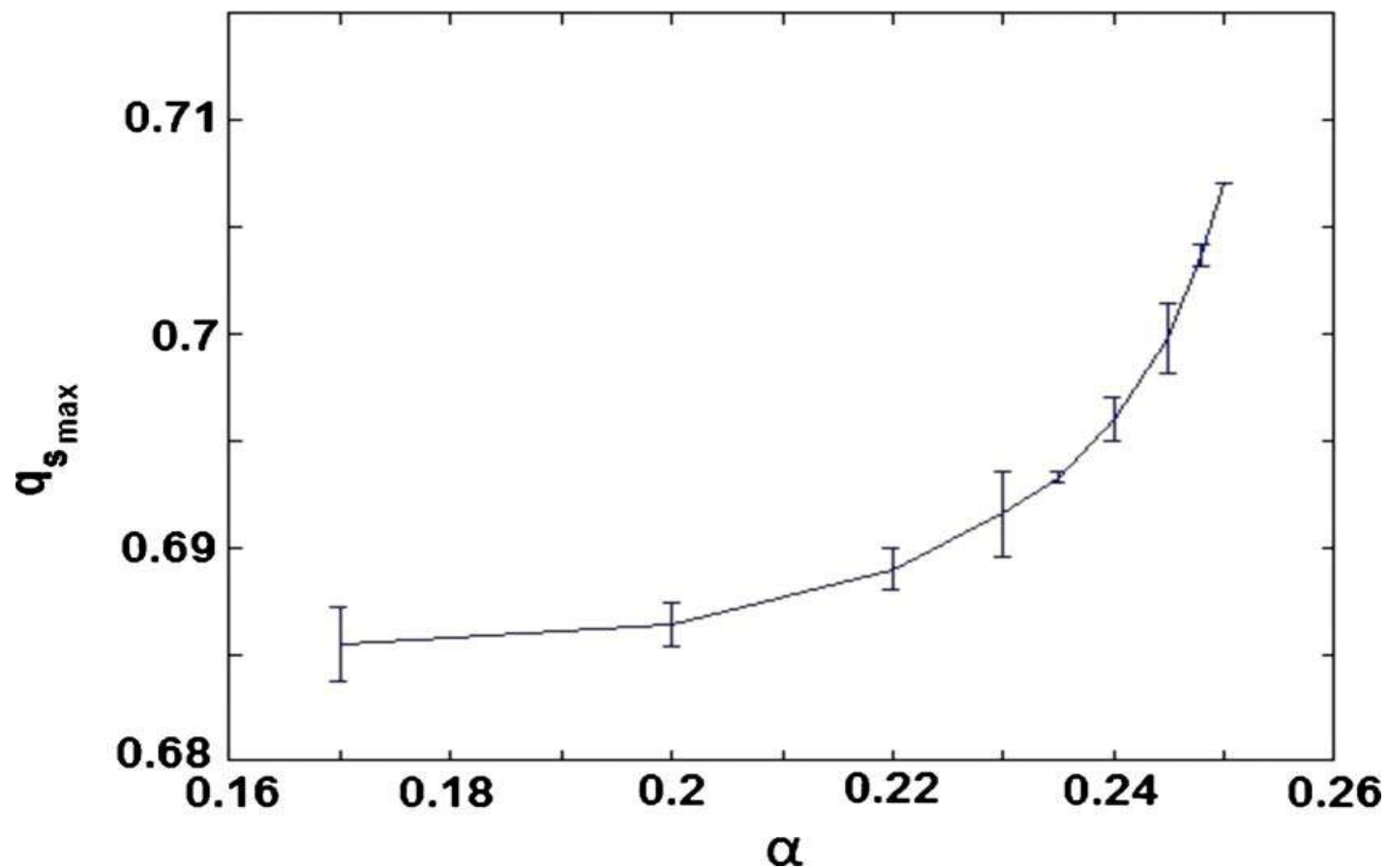


$D(q)$ for $\alpha=0.17$. $\lambda_{\parallel}(\max)=19$

$D(q)$ increases as a decreases. Patterns more stable.



We assume that maximum of $D(q)$ gives selected value since this is the most stable against perturbations. Agrees with simulation results for α close to 0.25.



Conclusions

- We have some evidence that a driven out of equilibrium system (eg. directional solidification, Rayleigh-Benard convection, etc) will eventually come to a unique stationary state.
- Evidence is mostly numerical on 1D stabilized KS equation.
- Analytical evidence is that the most stable state with the largest diffusion coefficient $D(q)$ seems to be selected.
- Clearly, much more work required to really establish this.