

The current distribution of the pushing asymmetric simple exclusion process

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Introduction and Motivation

- Central Limit Theorem states the the limiting distribution for the scaled fluctuation is Gaussian. ($1/2$ power law of exponent)
- KPZ(Kadar-Parisi-Zhang) predicted scale exponents $1/3$ for fluctuation (and $2/3$ for spatial correlation) for a large class of models.
- What are limiting distributions for KPZ universality class ?

Tracy-Widom distribution

GUE ensemble

The Gaussian Unitary Ensemble (GUE) is defined as a sequence \mathbb{P}_N of Gaussian probability measure on $N \times N$ Hermitian matrices of the form

$$d\mathbb{P}_N(M) = \frac{1}{Z_N} e^{\text{tr}(M^2)} dM$$

where Z_N is a normalizing constant.

Tracy-Widom distribution; 1994

Distribution of the largest eigenvalue in a Gaussian unitary random matrix ensemble

$$\lim_{N \rightarrow \infty} \mathbb{P}_N \left(\frac{\sqrt{2N} \lambda_1(M) - 2N}{N^{1/3}} \leq s \right) := TW_2(s)$$

Fluctuation of the length of the longest increasing subsequence of a random permutation (Baik, Deift, Johansson; 1999)

$$\lim_{n \rightarrow \infty} \mathbb{P}\left(\frac{L_n - 2\sqrt{n}}{n^{1/6}} \leq s\right) = TW_2(s)$$

Fluctuation of the current of the TASEP on \mathbb{Z} with step initial condition (the corner growth model) (Johansson;2000)

$F_0(t)$: the total number of particles that have crossed from 0 and 1 during the time $[0, t]$

$$\lim_{t \rightarrow \infty} \mathbb{P}\left(\frac{F_0(t) - Jt}{Vt^{1/3}} \leq s\right) = 1 - TW_2(-s)$$

where J and V are some constants

Techniques to obtain the Tracy-Widom distribution

- Combinatorics : Robinson-Schensted-Knuth (RSK) type bijection
 - Longest increasing subsequence of a random permutation
 - Discrete TASEP
- (Coordinate) Bethe ansatz for continuous-time ASEP on \mathbb{Z}

- (Schütz;1997) Transition probability $P_Y(X; t)$ of N -particle (continuous-time) TASEP
- It is a determinant of which entries are hypergeometric functions.
- Entries have *special* properties.
- $\mathbb{P}_Y(x_m(t) = x)$ is obtained by summing $P_Y(X; t)$ over all possible configurations
- Asymptotic study of $\mathbb{P}_Y(x_m(t) = x)$ for step initial condition rediscovers Johansson's result (Nagao, Sasamoto;2004, Rákos, Schütz;2005)
- How about ASEP ?

Tracy and Widom's works on ASEP;2008,2009

- Exact expression for the transition probability of N -particle system : the sum of $N!$ N -dimensional contour integrals
- $\mathbb{P}_Y(x_m(t) \leq x)$ to use some special identities
- Obtained a closed form of $\mathbb{P}_Y(x_m(t) \leq x)$ when Y is step initial condition and expressed it as a Fredholm determinant
- Asymptotic analysis of the Fredholm determinant $\rightarrow 1/3$ law

$$\lim_{t \rightarrow \infty} \mathbb{P}\left(\frac{\mathcal{T}([-vt], t) - a_1 t}{a_2 t^{1/3}}\right) = 1 - TW_2(-s)$$

where $\mathcal{T}([-vt], t)$ is the number of particles of which positions are less than or equal to $[-vt]$ at t , and $0 \leq v < 1$, a_1, a_2 are some constants.

More results on particle models

- TASEP with flat initial condition : TW_1 (Borodin, Ferrari, Sasamoto; 2007)
- (one-sided) pushing asymmetric simple exclusion process (PushASEP) : TW_1, TW_2 (Borodin, Ferrari; 2008)
- Depending on the geometry of initial conditions
- Based on determinantal form of transition probabilities
- Distribution of the Hopt-Cole solution of the KPZ equation with narrow edge initial condition based on Tracy-Widom's ASEP results

Dynamics of N -particle two-sided PushASEP

- Each site is occupied by at most one particle or empty.
- Each particle is equipped with $2N$ Poisson clocks with rates pr_n and ql_n with $p + q = 1$ and $n = 1, \dots, N$ and all clocks are independent.
- If (1) the clock with rate $pr_n(ql_n)$ of the particle at x rings, (2) $x + 1, \dots, x + n - 1$ ($x - 1, \dots, x - (n - 1)$) are occupied and (3) $x + n$ ($x - n$) is empty, then the particle at x jump to $x + n$ ($x - n$).
- Otherwise, nothing happens and the clock resumes.

Bethe ansatz applicability (Sasamoto, Wadati;1998)

Constraint on r_n and l_n ;

$$\frac{1 - \left(\frac{\lambda}{\mu}\right)^n}{1 - \frac{\lambda}{\mu}} = \frac{1}{r_n}$$
$$\frac{1 - \left(\frac{\mu}{\lambda}\right)^n}{1 - \frac{\mu}{\lambda}} = \frac{1}{l_n}$$

with $\lambda + \mu = 1$.

$\lambda \rightarrow 0$; pushing dynamics on the right with constant rate p and TASEP dynamics on the left with rate q , that is, one-sided PushASEP.

Transition probability of the two-sided PushASEP (L;2012)

Let $\lambda + \mu = 1$ ($1/2 < \mu \leq 1$) and \mathcal{C}_{R_i} ($i = 1, \dots, N$) be a circle oriented counterclockwise, centered at 0 with radius R_i . Assume that $1 < R_1 < \dots < R_N < c$ where

$$c = \begin{cases} \frac{\mu}{\lambda} & \text{if } \lambda \neq 0 \\ \infty & \text{if } \lambda = 0. \end{cases}$$

The transition probability of the two-sided PushASEP is

$$P_Y(X; t) = \sum_{\sigma \in \mathbb{S}_N} \left(\frac{1}{2\pi i} \right)^N \int_{\mathcal{C}_{R_N}} \cdots \int_{\mathcal{C}_{R_1}} A_{\sigma}^{\dagger} \prod_i^N \left(\xi_{\sigma(i)}^{x_i - y_{\sigma(i)} - 1} e^{\varepsilon(\xi_i)t} \right) d\xi_1 \cdots d\xi_N$$

where

$$A_{\sigma}^{\dagger} = \prod_{\substack{i < j, \\ \sigma(i) > \sigma(j)}} \frac{\xi_{\sigma(i)}}{\xi_{\sigma(j)}} \cdot S_{\sigma(i)\sigma(j)}.$$

and

$$S_{\beta\alpha} := -\frac{\mu + \lambda \xi_{\alpha} \xi_{\beta} - \xi_{\alpha}}{\mu + \lambda \xi_{\alpha} \xi_{\beta} - \xi_{\beta}} \quad (\mu + \lambda = 1)$$

Difference between S -matrices of the ASEP and the two-sided PushASEP

For the Bethe ansatz solution

$$\sum_{\sigma \in \mathbb{S}_N} A_\sigma \prod_i \xi_{\sigma(i)}^{x_i}$$

- ASEP

$$S_{\beta\alpha} := -\frac{p + q\xi_\alpha\xi_\beta - \xi_\beta}{p + q\xi_\alpha\xi_\beta - \xi_\alpha}$$

where $\alpha < \beta$.

- Two-sided PushASEP

$$S_{\beta\alpha} := -\frac{\xi_\beta}{\xi_\alpha} \cdot \frac{\mu + \lambda\xi_\alpha\xi_\beta - \xi_\alpha}{\mu + \lambda\xi_\alpha\xi_\beta - \xi_\beta}$$

where $\alpha < \beta$.

- This difference requires new proof.

Contours of the ASEP and the two-sided PushASEP

- ASEP: \mathcal{C}_r sufficiently small $r < 1$ and same for all variables
- Two-sided PushASEP: \mathcal{C}_{R_i} with $1 < R_1 < \dots < R_N < \frac{\mu}{\lambda}$ for variable ξ_i

Main part of the proof

Show the initial condition; $P_Y(X;0) = \delta_Y(X)$. In the sum $\sum_{\sigma \in \mathbb{S}_N}$

- the term for $\sigma = \textit{identity}$ contributes to $\delta_Y(X)$
- all other terms for $\sigma \neq \textit{identity}$ give 0.
- mathematical induction

$$\mathbb{P}_Y(x_m(t) = x)$$

- Probability that the m th leftmost particle is at x at time t

$$\mathbb{P}_Y(x_m(t) = x) = \sum_{\text{all possible } X} P_Y(X; t)$$

- $x_{m+1} = x + z_1, x_{m+2} = x + z_1 + z_2, \dots, x_N = x + z_1 + \dots + z_{N-m}, z_i \in \mathbb{N}$
- $x_{m-1} = x - v_1, x_{m-2} = x - v_1 - v_2, \dots, x_1 = x - v_1 - \dots - v_{m-1}, v_i \in \mathbb{N}$

$$\mathbb{P}_Y(x_m(t) = x) = \sum_{\text{all } v_i, z_i=1}^{\infty} P_Y(X; t)$$

- The sum is a multiple geometric series which can be shown to be convergent.

Main result $\mathbb{P}_Y(x_m(t) = x)$ (L;2012)

Let $S = \{s_1, \dots, s_k\} \subset \{1, \dots, N\}$ with $s_i < s_{i+1}$ and

$$I(\xi; s_1, \dots, s_k) = \prod_{i < j} \frac{\xi_{s_i} - \xi_{s_j}}{\mu + \lambda \xi_{s_i} \xi_{s_j} - \xi_{s_j}} \cdot \frac{1}{\prod_{s \in S} (\xi_s - 1)} \cdot \left(\prod_{s \in S} \xi_s - 1 \right)$$

Then

$$\begin{aligned} \mathbb{P}_Y(x_m(t) = x) &= \sum_{|S| \geq m} c_S \left(\frac{1}{2\pi i} \right)^k \int_{\mathcal{C}_{R_{s_k}}} \cdots \int_{\mathcal{C}_{R_{s_1}}} I(\xi; s_1, \dots, s_k) \\ &\quad \prod_{s \in S} \xi_s^{x - (y_s - s) - 1} e^{\varepsilon(\xi_s)t} d\xi_{s_1} \cdots d\xi_{s_k}. \end{aligned}$$

where

$$\varepsilon(\xi_s) = \frac{p}{\xi_s} + q\xi_s - 1$$

Coefficient c_S

Notation

$$[N] = \frac{\mu^N - \lambda^N}{\mu - \lambda}$$

and

$$[N]! = [N][N-1] \cdots [1], \quad \begin{bmatrix} N \\ m \end{bmatrix} = \frac{[N]!}{[m]![N-m]!}$$

with $[0]! = 1$

$$c_S = (-1)^{|S|+m} (\mu\lambda)^{m(m-1)/2} \begin{bmatrix} |S| - 1 \\ |S| - m \end{bmatrix} \frac{\lambda^{\Sigma[S]-m|S|}}{\mu^{\Sigma[S]-|S|(|S|+1)/2}}$$

where $\Sigma[S]$ is the sum of all elements in S .

- Surprisingly, c_S for the two-sided PushASEP is in the same form as the c_S for the ASEP. (p, q instead of λ, μ)

Sketch of the proof

- (1) Find $\mathbb{P}(x_N(t) = x)$; easy part using the same technique as the ASEP's $\mathbb{P}(x_1(t) = x)$
 - (2) Find $\mathbb{P}(x_1(t) = x)$; hard part
 - (3) Use $\mathbb{P}(x_N(t) = x)$ and $\mathbb{P}(x_1(t) = x)$ to find general $\mathbb{P}(x_m(t) = x)$
- Interestingly, no new identities required (ASEP's identities appeared in the PushASEP as well.)

More results to report and ongoing works

- Possible to obtain Fredholm determinant representation of $\mathbb{P}(x_m(t) = x)$ for step (Bernoulli) initial condition by the same technique as the Tracy-Widom's. (But slightly different operator so easy to apply the ASEP's technique.)
- Asymptotic analysis of the Fredholm determinant representation for the Tracy-Widom law (ongoing)
- Transition probability of the Bethe ansatz solvable zero range process on \mathbb{Z} (We found new identities which is a good signal in this direction of the work.)
- Ultimately, we want to generalize the ASEP's result to the Bethe ansatz solvable AZRP (ongoing).

Thank you.