

# Chern-Simons Theory and Its Applications

The 10<sup>th</sup> Summer Institute for Theoretical Physics  
Ki-Myeong Lee

# Maxwell Theory

★ Maxwell Theory:

$$\mathcal{L} = -\frac{1}{4g^2} F_{\mu\nu} F^{\mu\nu} - A_\mu J^\mu$$

★ Gauge Transformation and Invariance

$$A_\mu \Rightarrow A_\mu + \partial_\mu \Lambda, \quad \partial_\mu J^\mu = 0$$

★ Gauss Law

$$-\partial_i \left( \frac{1}{g^2} F_{i0} \right) - J^0 = 0$$

★ Charge

$$Q = \int d^3x J^0 = - \int d\Sigma_{S_\infty^2}^i \frac{1}{g^2} F_{i0}$$

★ Degrees of Freedom: In 4-dimension, two helicity  $\pm 1$

★ In 3-dim, one massless scalar  $\frac{1}{g^2} F_{\mu\nu} = \epsilon_{\mu\nu\rho} \partial^\rho \varphi$

# Chern-Simons Term

- ★ Chern-Simons Theory:

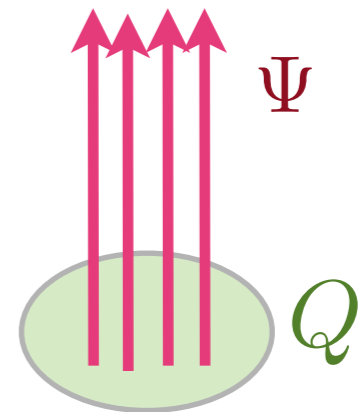
$$\mathcal{L} = \frac{\kappa}{2} \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho - A_\mu J^\mu$$

- ★ Gauge Transformation and Invariance: Total derivative

$$\delta\mathcal{L} = \frac{\kappa}{4} \epsilon^{\mu\nu\rho} \partial_\mu (\Lambda F_{\nu\rho})$$

- ★ Field Equation

$$F_{i0} = -\frac{1}{\kappa} \epsilon_{ij} J^j \quad F_{12} = \frac{1}{\kappa} J^0$$



- ★ Gauss Law: Flux-Charge Composite

$$\Psi = \int d^2x F_{12} = \frac{1}{\kappa} Q$$

- ★ Without Current, the Gauge Field Becomes Pure Gauge  $F_{\mu\nu} = 0$

- ★ Topological Theory: Independent of Metric

$$S = \frac{\kappa}{2} \int d^3 \sqrt{g} \frac{1}{\sqrt{g}} \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho = \frac{\kappa}{2} \int A \wedge dA$$

# Topics to be covered

- ★ Anyons: Particles of Fractional Spin and Statistics
- ★ Maxwell-Chern-Simons -(Higgs)Theory
- ★ Non-Abelian Chern-Simons Theory
- ★ Quantization
- ★ Massive Fermion
- ★ Induced Chern-Simons Term
- ★ Knot Invariants, Jons Polynomial
- ★ Matter Coupling
- ★ Supersymmetry and Superconformal Symmetry
- ★ 2+1 dimensional Gravity and Chern-Simons Theory
- ★ Aharony-Bergman-Jafferis-Maldacena Theory of M2 Branes
- ★ Quantum Hall Effect

# Anyons

- ★ Identical Charged Particles with Conserved Current:

$$J^\mu(x) = e \sum_a (1, \dot{\mathbf{q}}_a) \delta^2(\mathbf{x} - \mathbf{q}_a)$$

- ★ The Lagrangian for these Particles is

$$L = \frac{m}{2} \sum_a \dot{\mathbf{q}}_a^2 + \int d^2x \left[ \frac{k}{2} \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho - A_\mu J^\mu \right]$$

- ★ Field Equation:  $F_{12} = \frac{e}{\kappa} \sum_a \delta(\mathbf{x} - \mathbf{q}_a)$

- ★ Use the relation  $\epsilon_{ij} \partial_i \partial_j \text{Arg}(\mathbf{x}) = 2\pi \delta^2(\mathbf{x})$

- ★ Pure Gauge (Singular)

$$A_\mu = \frac{e}{2\pi\kappa} \partial_\mu \text{Arg} \delta^2(\mathbf{x} - \mathbf{q}_a)$$

- ★ Momentum and Hamiltonian

$$\mathbf{p}_a = m\dot{\mathbf{q}}_a - e\mathbf{A}(\mathbf{x} = \mathbf{q}_a, t)$$

$$\mathcal{H} = \frac{1}{2m} |\mathbf{p} + e\mathbf{A}(\mathbf{q}_a)|^2$$

# Anyon (2)

- ★ Two Body & Center of Mass and Relative Positions

$$\mathbf{R} = \frac{1}{2}(\mathbf{q}_1 + \mathbf{q}_2), \quad \mathbf{r} = \mathbf{q}_1 - \mathbf{q}_2$$

- ★ Ignore Self-Interaction and the Lagrangian becomes

$$\mathcal{L} = \frac{M}{2}\dot{\mathbf{R}}^2 + \frac{\mu}{2}\dot{\mathbf{r}}^2 - \frac{e}{2\pi\kappa}\dot{\mathbf{r}} \cdot \nabla \text{Arg}(\mathbf{r})$$

- ★ The Relative Hamiltonian

$$\mathcal{H} = \frac{1}{2\mu}\left(\mathbf{p} + \frac{e}{2\pi\kappa}\nabla \text{Arg}(\mathbf{r})\right)^2$$

- ★ Quantize and Consider the Energy Eigen States

$$\mathbf{p} = -i\nabla, \quad \mathcal{H}\Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

- ★ Assume the original particles are bosons, and so the wave function is symmetric under the exchange.

$$\Psi(\mathbf{r}) = \Psi(-\mathbf{r}) \quad \Psi(r, \theta) = f(r)e^{2in\theta},$$

# Anyon(3)

- ★ Make a Singular Gauge Transformation

$$\Psi(\mathbf{r}) = e^{-\frac{ie}{2\pi\kappa} \text{Arg}(\mathbf{r})} \Psi_A(\mathbf{r}) \qquad \Psi_A(r, \theta) = f(r) e^{i(2n + \frac{e}{2\pi\kappa})\theta}$$

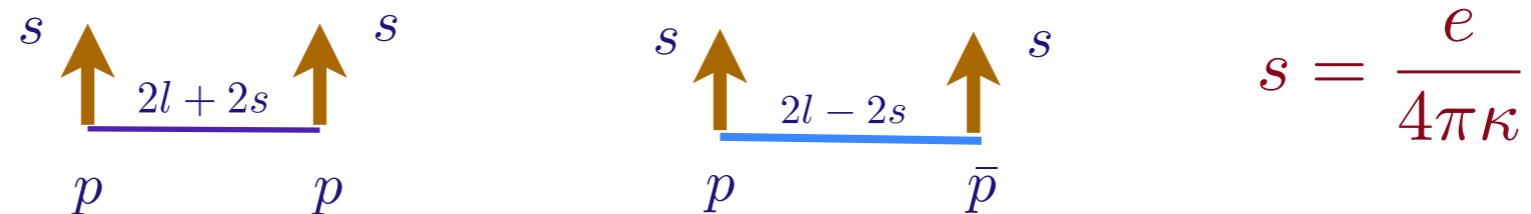
- ★ The Hamiltonian becomes free.

$$\mathcal{H}_{\text{free}} \Psi_A(\mathbf{r}) = E \Psi_A(\mathbf{r}), \quad \mathcal{H}_{\text{free}} = \frac{\mathbf{p}^2}{2\mu} = \frac{1}{2\mu} \left( p_r^2 + \frac{L^2}{r^2} \right)$$

- ★ the Orbital Angular Momentum

$$L_{\text{free}} = -i\partial_\theta = 2n + \frac{e}{2\pi\kappa}$$

- ★ Spin-Statistics: Relativity + Positive Energy (Fiertz, Pauli)



$$s = \frac{e}{4\pi\kappa}$$

- ★ Exchange Statics of Anyons:

$$\Psi_A(r, \theta + \pi) = e^{i2\pi s} \Psi_A(r, \theta)$$

# New Kinds of Particles in 2+1 dim

★ 3-dim Rotation  $SO(3)=SU(2)$ :  $|s,m\rangle$

$$s = 0, \frac{1}{2}, 1, \frac{3}{2}, 2, \frac{5}{2} \dots$$

★ 2-dim Rotation  $U(1)$ :  $|s\rangle$

$$U(\theta)|s\rangle = e^{is\theta}|s\rangle$$



★ First Homotopy Group of Circle

★ Braid Group



# Maxwell-Chern-Simons + Higgs

- ★ 2+1 dim Maxwell Theory = A Single Massless Scalar Theory

$$\mathcal{L}_M = -\frac{1}{4g^2} F_{\mu\nu} F^{\mu\nu} \Leftrightarrow \mathcal{L}_s = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi \quad F_{\mu\nu} = \epsilon_{\mu\nu\rho} \partial^\rho \phi$$

- ★ Maxwell-Chern-Simons Theory

$$\mathcal{L}_{MCS} = -\frac{1}{4g^2} F_{\mu\nu} F^{\mu\nu} - \frac{\kappa}{2\pi} \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho$$

- ★ Field Equation

$$\partial_\sigma F^{\sigma\mu} - \frac{g^2 \kappa}{4\pi} \epsilon^{\mu\nu\rho} F_{\nu\rho} = 0 \quad \Rightarrow \quad (\partial_\mu^2 + m^2) F_{\mu\nu} = 0 \quad \text{with } m = \frac{g^2 \kappa}{2\pi}$$

- ★ Massive Neutral Vector Boson of Mass  $|m|$  and Spin  $\frac{m}{|m|} = \frac{\kappa}{|\kappa|}$

- ★ Maxwell + Higgs :  $\mathcal{L} = -\frac{1}{4g^2} F_{\mu\nu} F^{\mu\nu} + v^2 |\partial_\mu \varphi + A_\mu|^2$

- ★ Parity Even and Massive Vector Boson of Mass  $|gv|$  and Spin  $\pm 1$

- ★ Chern-Simons-Higgs: Mass  $\left| \frac{v^2}{\kappa} \right|$  and Spin  $-\frac{\kappa}{|\kappa|}$

# Non-Abelian Theory

- ★ The Lagrangian

$$\mathcal{L} = \frac{\kappa}{2} \epsilon^{\mu\nu\rho} \text{Tr}(A_\mu \partial_\nu A_\rho + \frac{2i}{3} A_\mu A_\nu A_\rho), \quad D_\mu \Phi = (\partial_\mu + A_\mu) \Phi$$

- ★ Under the gauge transformation  $A_\mu \rightarrow A^g = \bar{g} A_\mu g + \bar{g} \partial_\mu g$

$$\delta \mathcal{L} = \frac{\kappa}{2} \epsilon^{\mu\nu\rho} \partial_\mu \text{Tr}(\partial_\nu g \bar{g} A_\rho) + \frac{\kappa}{6} \epsilon^{\mu\nu\rho} \text{Tr}(\bar{g} \partial_\mu g \bar{g} \partial_\nu g \bar{g} \partial_\rho g)$$

- ★ The third homotopy group of Lie group is interger.  $\pi_3(G) = \mathbb{Z}$

- ★ The winding number is  $N = \frac{i}{24\pi^2} \int d^3x \epsilon^{\mu\nu\rho} \text{Tr}(\bar{g} \partial_\mu g \bar{g} \partial_\nu g \bar{g} \partial_\rho g) \in \mathbb{Z}$

$$S_{CS} \rightarrow S_{CS} + 2\pi N(2\pi\kappa) \quad e^{iS_{CS}} \rightarrow e^{iS_{CS} + 2\pi i N(2\pi\kappa)}$$

- ★ The Path Integral is gauge invariant and well defined if

$$2\pi\kappa = k = \text{integer}$$

- ★ The Chern-Simons Level is the integer  $k$

# Simple Exercise for SU(2) Case

★ SU(2) = Three Sphere

$$g = x_4 + i\sigma_i x_i, \quad \bar{g}g = 1, \det g = 1 \Rightarrow x_\mu \in R, \sum_{\mu} x_\mu^2 = 1$$

★ Near the Pole  $x_0 = 0,$

$$\text{Tr}(\bar{g}dg)^3 = i\text{Tr}(\sigma_i dx_i)^3 = 12dx \wedge dy \wedge dz$$

★ Then as Volume of  $S^3 = 2\pi^2$

$$N = \frac{1}{24\pi^2} 12 \int_{S^3} dx \wedge dy \wedge dz = 1$$

# Quantization of Maxwell-Chern-Simons

★ Lagrangian  $\mathcal{L} = \frac{1}{2g^2} F_{0i}^2 - \frac{1}{2e^2} F_{12}^2 + \frac{\kappa}{2} \epsilon^{ij} \dot{A}_i A_j + \kappa A_0 F_{12}$

★ Momenta and Gauss Law

$$\Pi_i = \frac{1}{e^2} F_{0i} + \frac{\kappa}{2} \epsilon^{ij} A_j \quad \partial_i \Pi_i + \frac{\kappa}{2} F_{12} = 0$$

★ Hamiltonian

$$\mathcal{H} = \Pi_i \dot{A}_i - \mathcal{L} = \frac{e^2}{2} (\Pi_i - \frac{\kappa}{2} \epsilon_{ij} A_j)^2 + \frac{1}{2e^2} F_{12}^2 + A_0 (\partial_i \Pi_i + \frac{\kappa}{2} F_{12})$$

★ Quantization

$$[A_i(\mathbf{x}, t), \Pi_j(\mathbf{y}, t)] = i \delta_{ij} \delta^2 \mathbf{x} - \mathbf{y} \quad \Pi_i(\mathbf{x}) = -i \frac{\delta}{\delta A_i(\mathbf{x})}$$

★ The Wave Function  $\Psi(A_1, A_2) \quad (\partial_i \Pi_i + \frac{\kappa}{2} F_{12}) \Psi = 0$

★ Decompose  $A_i = \partial_i \lambda + A_i^T, \quad \partial_i A_i^T = 0, \quad \lambda = \nabla^{-2} \partial_i A_i$

$$\Psi = e^{-\frac{i\kappa}{2} \int d^2 x \lambda F_{12}} \Phi(A_i^T), \quad \Phi(A_i^T + \partial_i \lambda) = \Phi(A_i^T)$$

# Quantization of Pure Chern-Simons

★ The Commutation Relation  $[A_1(\mathbf{x}), A_2(\mathbf{y})]_{ET} = \frac{i}{k} \delta^2(\mathbf{x} - \mathbf{y})$

★ Wave Function  $\Psi(A_1)$

★ Complex or holomorphic coordinate  $z = x^1 + ix^2$   $A = \frac{1}{2}(A_1 + iA_2)$

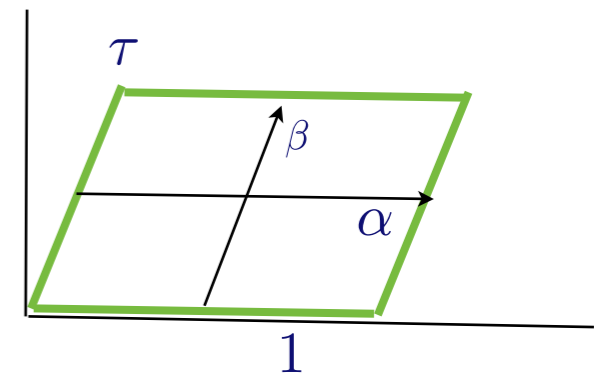
★ Composition  $A_i = \partial_i \omega + \epsilon_{ij} \partial_j \sigma \Rightarrow A = \partial_{\bar{z}} \chi, \chi = \omega - i\sigma$

★ CS Theory on a Complex Riemann Surface  $\Sigma$  of genus  $g$

★ On torus with side 1 and  $\tau$  and area  $\text{Im}(\tau)$

★ Nontrivial cohomology on torus  $\omega(z)$

$$\int_{\alpha} \omega = 1, \int_{\beta} \omega = \tau, \int_{\text{torus}} |\omega|^2 = \text{Im}(\tau)$$



★ The gauge field is the sum of exact and closed 1-form

$$A = \partial_{\bar{z}} \chi(z, \bar{z}, t) + \frac{i\pi a(t)}{\text{Im}(\tau)} \bar{\omega}$$

# Massive (Majorana) Fermions

★ Gamma Matrices  $\gamma^0 = \sigma_2, \gamma^1 = i\sigma_1, \gamma^2 = i\sigma_3$

★ Two Component Spinor  $\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$

★ Lagrangian and Field Equation

$$\mathcal{L} = i\bar{\psi}\gamma^\mu\partial_\mu\psi - m\bar{\psi}\psi \qquad i\gamma^\mu\partial_\mu\psi - m\psi = 0$$

★ Majorana condition  $\psi = \psi^*$

★ Quantization and Elementary Particles

★ Neutral Particles of Mass  $m$  and Spin  $\frac{1}{2}\text{sign}(m)$

★ The Mass Term is both P and CP violating

★ Dirac Fermions of Charge  $\pm 1$

★ Related to Massive Vector Boson of Maxwell-Chern-Simons by Supersymmetry and Quantum Corrections

# Induced Chern-Simons Term

★ Gauged Charged Dirac Particles  $\mathcal{L} = \bar{\psi}(i\gamma^\mu D_\mu - m)\psi$

★ Effective Action  $e^{iS_{eff}(A)} = \int d\psi d\bar{\psi} e^{i\mathcal{L}}$

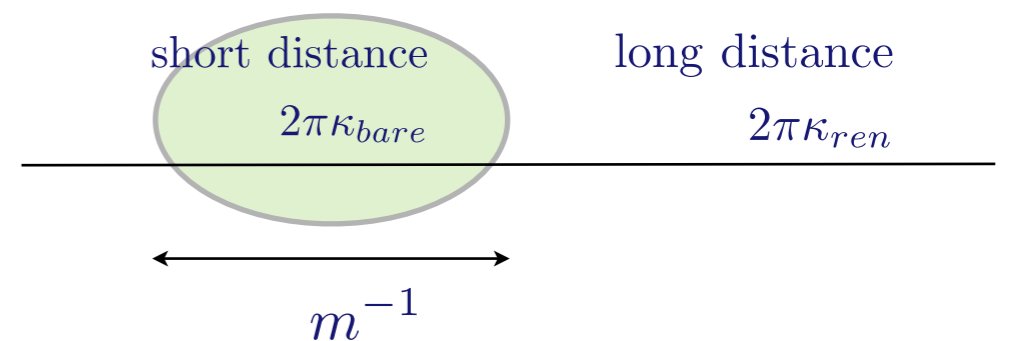
★ Calculation 
$$\begin{aligned} S_{eff}(A) &= N_f \text{tr} \ln(i\cancel{\partial} + A - m) \\ &= N_f \text{tr} \ln(i\cancel{\partial} - m) + N_f \text{tr} \left( \frac{1}{i\cancel{\partial} - m} A \right) \\ &\quad + N_f \text{tr} \left( \frac{1}{i\cancel{\partial} - m} A \frac{1}{i\cancel{\partial} - m} A \right) + \dots \end{aligned}$$

★ Long Wave Length

$$S_{eff}(A) \approx \frac{iN_f}{2} \frac{m}{|m|} \frac{1}{4\pi} \int d^3x \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho$$

★ Renormalization of Chern-Simons

$$2\pi\kappa_{ren} = 2\pi\kappa_{bare} + \frac{N_f}{2} \frac{m}{|m|}$$



# Supersymmetric Yang-Mills Chern-Simons

★ Yang-Mills-Chern-Simons:

$$\mathcal{L} = -\frac{1}{g^2} \text{Tr} F_{\mu\nu} F^{\mu\nu} + \frac{k}{4\pi} \epsilon^{\mu\nu\rho} \text{Tr} (A_\mu \partial_\nu A_\rho - \frac{2i}{3} A_\mu A_\nu A_\rho)$$

★ Mass  $m = \frac{kg^2}{2\pi}$  Spin  $\frac{k}{|k|}$

★ Low Energy Dynamics:  $Energy \ll |m| \Rightarrow$  Pure Chern-Simons Theory

★ Renormalization of k  $k_{ren} = k_{bare} + N$

★ Supersymmetry and Level Shift

spin	1	1/2	0	-1/2	-1	$k_{ren}$
N=0	1					$k+N$
N=1	1	1				$k+N/2$
N=2	1	2	1			$k$
N=3	1	3	3	1		$k$



# SUSY YMCS Theory on Torus (Witten)

- ★ Index Calculation on Two Torus of Size L  $Z = \text{Tr}_{\mathcal{H}} (-1)^F e^{-\beta H}$
- ★ Small Torus Limit and all KK Modes are Ignored.  $g^2 \kappa \ll 1/L$
- ★ Low Energy Dynamics: Flat Diagonal Connections=constant

$$A_i = \sum_{a=1}^r c_i^a T^a, \quad \lambda_{\pm} = \sum_{a=1}^r \eta_{\pm}^a T^a,$$

- ★ Supersymmetric Quantum Mechanics on the Vacuum Moduli Space

$$Q_{\pm} = \int_{T^2} \text{Tr} \lambda_{\mp} \frac{D}{DA_{\bar{z},z}}, \quad \frac{D}{DA_i^a} = \frac{\delta}{\delta A_i^a(x)} + \frac{ik}{4\pi} \epsilon_{ij} A_j^a$$

- ★ Supersymmetry is Broken when  $N = 1 \quad k < N/2, \quad N = 2, 3 \quad k < N$
- ★ Vacuum Degeneracy for Other Cases..

$$I(k) = \frac{(N+t-1)!}{t!(N-1)!}, \quad t = k - \frac{N}{2} \text{ or } t = k - N$$

# Knot and Jones Polynomial (Witten)

- ★ The Partition Function

$$Z = \int [dA_\mu] e^{iS_{CS}} \prod_{i=1}^r W_{R_i}(C_i)$$

- ★ Wilson Loop

$$W_{R_i}(C_i) = \text{Tr}_{R_i} P \exp(i \int dx^\mu A_\mu^a T_{R_i}^a)$$



- ★ Quantize on Space-Time  $\Sigma \times R$

- ★ On Riemann Surface with Marked Points

$$\Sigma, F_{12} = \sum_i T_{R_i} \delta^2(\mathbf{x} - \mathbf{q}_i)$$

- ★ Conformal Blocks, Quantum-Groups,....

# 2+1 Dimensional Gravity (1) (Witten)

★ d-dim Einstein-Hilbert Action  $I = \frac{1}{2} \int d^d x \sqrt{g} R$

★ Vierbein  $e^a = e^a_\mu dx^\mu$  and Spin Connection  $w^a_b = w^a_{b\mu} dx^\mu$

★ Relations  $g_{\mu\nu} = e^a_\mu e^b_\nu \eta_{ab}$   $de^a + w^a_b \wedge e^b = 0$

★ Curvature Tensor  $R^a_{\mu\nu b} = \partial_\mu w^a_{b\nu} - \partial_\nu w^a_{b\mu} + [w_\mu, w_\nu]^a_b$

★ 4-dim  $I_4 = \frac{1}{2} \int d^4 x \epsilon^{\mu\nu\rho} \epsilon_{abcd} e^a_\mu e^b_\nu R^{cd}_{\rho\sigma}$

★ 3-dim with Cosmological Constant

$$I_3 = \frac{1}{2} \int d^3 x \epsilon^{\mu\nu\rho} \epsilon_{abc} \left( e^a_\mu (\partial_\nu w^b_\rho - \partial_\rho w^b_\nu) + [w_\nu, w_\rho]^{bc} + \lambda e^a_\mu e^b_\nu e^c_\rho \right)$$

★ Anti-deSitter Space is the Solution with Isometry SO(2,2)

★ Generators of  $J_a, P_a, a = 1, 2, 3$   $J^{ab} = \epsilon^{abc} J_c$

$$[J_a, J_b] = \epsilon_{abc} J^c, [J_a, P_b] = \epsilon_{abc} P_c, [P_a, P_b] = \lambda \epsilon_{abc} J^c$$

★ Invariant quadratic form

$$\langle J_a, P_b \rangle = \delta_{ab}, \langle J_a, J_b \rangle = 0 = \langle P_a, P_b \rangle$$

# 2+1 Dim Gravity (2)

★ Gauge Connection  $w_\mu^a = \frac{1}{2}\epsilon^{abc}w_{bc\mu}$

$$A_\mu = e_\mu^A P_a + w_\mu^a J_a$$

★ Gauge Transformation = infinitesimal diffeomorphism and Lorentz transformation modulo Einstein equation.

$$\delta A_\mu = D_\mu u = \partial_\mu u + [A_\mu u]$$

★ Zero Field Strength  $F_{\mu\nu} = 0$  is identical to Einstein Eq..

★ Note  $SO(2, 2) = SL(2, R) \times SL(2, R)$

★ Decompose  $J_a^\pm = \frac{1}{2}(J_a \pm \frac{1}{\sqrt{\lambda}}P_a)$   $A_\mu^{a\pm} = w_\mu^a \pm \sqrt{\lambda}e_\mu^a$

★ Action  $I^\pm = \int \epsilon^{\mu\nu\rho} (2A_\mu^{a\pm} \partial_\nu A_\rho^{a\pm} + \frac{2}{3}\epsilon_{abc} A_\mu^{a\pm} A_\nu^{b\pm} A_\rho^{c\pm})$

★ Einstein + lambda =  $\frac{1}{2}(I^+ - I^-)$

★ 3-dim Graviational Chern-Simons =  $\frac{1}{2}(I^+ + I^-) \sim \int w \wedge dw$

# Supersymmetric CS Theory(I)

★ N=2 Chern-Simons Higgs  $V, \Phi, \Sigma = \bar{D}^\alpha D_\alpha V$

$$S_{Ab}^{\mathcal{N}=2} = \int d^3x \int d^4\theta \left( \frac{k}{4\pi} V \Sigma + \bar{\Phi} e^V \Phi \right)$$

★ Nonabelian  $S^{\mathcal{N}=2} = \int d^3x \int d^4\theta \left\{ \frac{k}{2\pi} \int_0^1 dt \text{Tr} [V \bar{D}^\alpha (e^{-tV} D_\alpha e^{tV})] + \bar{\Phi} e^V \Phi \right\}$

★ Pure Chern-Simons  $S_{CS}^{\mathcal{N}=2} = \frac{k}{4\pi} \int \text{Tr}(A \wedge dA + \frac{2}{3} A^3 - \bar{\chi} \chi + 2D\sigma)$

★ Matter  $S_{matter} = \int d^4\theta \sum_{i=1}^{N_f} \bar{\Phi}^i e^V \Phi^i = \int \sum_{i=1}^{N_f} (D_\mu \bar{\phi}^i D^\mu \phi^i + i \bar{\psi}^i \gamma^\mu D_\mu \psi^i - \bar{\phi}^i \sigma^2 \phi^i + \bar{\phi}^i D \phi^i - \bar{\psi}^i \sigma \psi^i + i \bar{\phi}^i \bar{\chi} \psi^i - i \bar{\psi}^i \chi \phi^i)$ .

★ FI Term:  $S_{FI} = \int d^3x \int d^4\theta \frac{k}{4\pi} (-2v^2) V = - \int d^3x \frac{k}{2\pi} v^2 D$

# Supersymmetric CS Theory (2)

★ D-Term Eq  $\sigma = \frac{2\pi}{k} \left( \sum_f e_f |\phi_f|^2 - v^2 \right)$

★ Potential  $U = \sum_f e_f^2 |\phi_f|^2 \sigma^2$

★ N=2 Superconformal Theory when  $v=0$

★ N=3 Theory: N=2+ Adjoint Chiral multiplet  $\Phi$ , hypermultiplet  $Q_i, \tilde{Q}_i$

★ N=3 Lagrangian

$$S^{\mathcal{N}=2} = S_{CS}^{\mathcal{N}=2} + \int d^4\theta (\bar{Q} e^V Q + \tilde{Q} e^{-V} \bar{\tilde{Q}}) + \left[ \int d^2\theta \left( -\frac{k}{4\pi} \text{Tr} \Phi^2 + \tilde{Q} \Phi Q \right) + c.c. \right]$$

★ Integration over Adjoint Chiral  $\Rightarrow$  Superpotential

$$W = \frac{2\pi}{k} (\tilde{Q} T^a Q) (\tilde{Q} T^a Q).$$

★ N=3  $SO(3)_R$  R-Symmetry, Superconformal Symmetry,

★ Susy preserving FI Term, Mass-Term breaks Conformal Symmetry

# Maxwell-Matter

★ The Lagrangian

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + |D_\mu\phi|^2 - \frac{1}{4}(|\phi|^2 - v^2)^2$$

★ Energy Density  $E = \frac{1}{2}F_{12}^2 + |D_i\phi|^2 + U$

★ BPS Bound

$$E = \frac{1}{2}(F_{12} - (v^2 - |\phi|^2))^2 + |(D_1 + iD_2)\phi|^2 + v^2 F_{12}$$

★ Energy is Bounded  $\mathcal{H} = \int d^2x E \geq v^2\Psi = v^2 \int d^2x F_{12}$

★ BPS Equation  $(D_1 + iD_2)\phi = 0, F_{12} = \frac{1}{2}(v^2 - |\phi|^2)$

★ Vortices:

$$\phi = f(r)e^{in\theta}, \quad -\partial_i^2 \ln f^2 = (v^2 - f^2) + 4\pi n\delta^2(x)$$

# Abelian CS Matter Theory (1)

★ Lagrangian  $\mathcal{L} = \frac{k}{4\pi} \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho + D_\mu \bar{\phi} D^\mu \phi - U(|\phi|)$

★ Gauss Law  $\frac{k}{2\pi} F_{12} - i(\phi D_0 \bar{\phi} - D_0 \phi \bar{\phi}) = 0$

★ Integration  $\frac{k}{2\pi} \Psi_{flux} = Q = \int d^2x (i\phi\pi_\phi - i\bar{\phi}\pi_{\bar{\phi}})$

★ Angular Momentum  $J = - \int d^2x \epsilon_{ij} (D_0 \bar{\phi} D_j \phi + D_j \bar{\phi} D_0 \phi)$

★ Q-Balls in Symmetric Phase and Vortices in Broken Phase



# Abelian CS Matter(2)

- ★ Potential (N=2 Supersymmetry)

$$U(\phi) = \frac{4\pi^2}{k^2} |\phi|^2 (|\phi|^2 - v^2)^2$$



- ★ Energy Density + Gauss Law

$$E = |D_0\phi \mp \frac{2\pi i}{k} (|\phi|^2 - v^2)|^2 + |(D_1 \mp iD_2)\phi|^2 \pm v^2 F_{12}$$

- ★ BPS Bound

$$\mathcal{H} = \int d^2x E \geq v^2 |\Psi|$$

- ★ Solutions in Symmetric Phase: Q-balls and Q-balls with vortices

$$J = \frac{Q^2}{2k}$$

- ★ Solutions in Broken Phase: Vortices

$$\Psi = 2\pi n, \quad Q = nk, \quad J = -\frac{Q^2}{2k}$$

# Nonrelativistic Limit of Symmetric Phase

★ Jackiw-Pi Model  $\mathcal{L}_{\text{JP}} = \frac{\kappa}{2} \epsilon^{\mu\nu\rho} A_\mu \partial_\nu A_\rho + i\psi^* D_0 \psi - \frac{1}{2m} |\vec{D}\psi|^2 + \frac{g}{2} |\psi|^4$   $g = \frac{1}{m\kappa}$

★ Gauss Law  $F_{12} = \frac{1}{\kappa} |\psi|^2$

★ Energy  $\mathcal{E} = \int d^2x \frac{1}{2m} |(D_1 - iD_2)\psi|^2$

★ BPS Equation=Liouville Equation  $\nabla^2 \ln \rho = -\frac{2}{\rho}$   
 $\rho = |\psi|^2$

★ Solution Found Explicitly  $\rho = \frac{4\kappa N^2}{r_0^2} \frac{\left(\frac{r}{r_0}\right)^{2(N-1)}}{\left(1 + \left(\frac{r}{r_0}\right)^{2N}\right)^2}$

★ Nonrelativistic Conformal Symmetry or Schrodinger Symmetry

# On M2,D3,M5 Branes

- ★ Theory on D3 Branes: N=4 Supersymmetric Yang-Mills Theory, SCFT, dual to

$$AdS_5 \times S^5$$

- ★ Theory on M2 Branes: Strong Coupling Limit of N=8 Supersymmetric Yang-Mills Theory, dual to

$$AdS_4 \times S^7$$

- ★ Theory on M5 Branes: 6-dim (2,0) SCFT Theory with SO(5) R-symmetry dual to

$$AdS_7 \times S^4$$

- ★ Abelian (2,0) Theory  $B_{\mu\nu}, \Psi, \Phi_I, I = 1, \dots, 5, , H_{\mu\nu\rho} = \tilde{H}_{\mu\nu\rho}$

- ★ 2+1 Conformal Field Theory

$$\mathcal{L} \sim AdA + (D\phi)^2 + \bar{\psi}D\psi + \bar{\psi}\psi|\phi|^2 + |\phi|^6$$

# Higher Super Symmetric CS Theories

- ★ Supersymmetric Algebra: Central Term

$$\{Q, Q\} \sim P + Z, \quad [R, Q] \sim Q, \quad [Z, Q] = 0 = [Z, R]$$

- ★ Non-Central Term for  $d < 4$

$$\{Q, Q\} \sim P + R, \quad [R, Q] \sim Q,$$

- ★ Yang-Mills Chern-Simons (N=4 is possible.)

- ★ Bagger-Lambert-Gustavsson Theory:  $SU(2)_k \times SU(2)_{-k} = SO(4)$ , N=8  $SO(8)$  R-symmetry

$$X_I^a, \Psi_A^a, \quad a = 1, 2, 3, 4, I = 1, 2, \dots, 8, A = 1, \dots, 8_c$$



$$\begin{aligned} \mathcal{L} = & -\frac{1}{2} D_\mu X_I^a D^\mu X_I^a + \frac{i}{2} \bar{\Psi}^a \Gamma^\mu D_\mu \Psi^a \\ & - \frac{i}{4\kappa} f^{abcd} \bar{\Psi}^a \Gamma_{IJ} \Psi^b X_I^c X_J^d - \frac{1}{12\kappa^2} \sum_{a,I,J,K} \left( f^{abcd} X_I^b X_J^c X_K^d \right)^2 \\ & + \frac{\kappa}{2} \epsilon^{\mu\nu\rho} \left( f^{abcd} A_\mu^{ab} \partial_\nu A_\rho^{cd} + \frac{2}{3} f^{acde} f^{bcfg} A_\mu^{ab} A_\nu^{de} A_\rho^{fg} \right), \end{aligned}$$

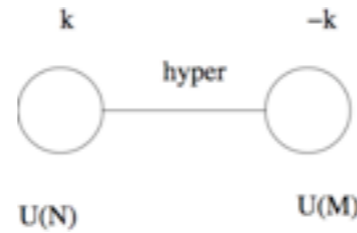
# N=4 Supersymmetric Theory

- ★ Gaiotto-Witten: Gauge Field + Hypermultiplet

$$\mathcal{L} = \frac{\varepsilon^{\mu\nu\lambda}}{4\pi} \left( k_{mn} A_\mu^m \partial_\nu A_\lambda^n + \frac{1}{3} f_{mnp} A_\mu^m A_\nu^n A_\lambda^p \right) + \frac{1}{2} \omega_{AB} \left( -\varepsilon^{\alpha\beta} Dq_\alpha^A Dq_\beta^B + i\varepsilon^{\dot{\alpha}\dot{\beta}} \psi_{\dot{\alpha}}^A \not{D} \psi_{\dot{\beta}}^B \right) - i\pi k_{mn} \varepsilon^{\alpha\beta} \varepsilon^{\dot{\gamma}\dot{\delta}} J_{\alpha\dot{\gamma}}^m J_{\beta\dot{\delta}}^n - \frac{\pi^2}{6} f_{mnp} (\mu^m)^\alpha_\beta (\mu^n)^\beta_\gamma (\mu^p)^\gamma_\alpha, \quad (j_{\alpha\dot{\gamma}}^m \equiv q_\alpha^A t_{AC}^m \psi_{\dot{\gamma}}^C). \quad (2.15)$$

- ★ Gauge Group U(N) x U(M), O(N) x Sp(2M)

- ★ Linear Quiver



- ★ Adding Twisted Hyper (HLLLP) includes BLG

$$\mathcal{L} = \frac{\varepsilon^{\mu\nu\lambda}}{4\pi} \left( k_{mn} A_\mu^m \partial_\nu A_\lambda^n + \frac{1}{3} f_{mnp} A_\mu^m A_\nu^n A_\lambda^p \right) + \frac{1}{2} \omega_{AB} \left( -\varepsilon^{\alpha\beta} Dq_\alpha^A Dq_\beta^B + i\varepsilon^{\dot{\alpha}\dot{\beta}} \psi_{\dot{\alpha}}^A \not{D} \psi_{\dot{\beta}}^B \right) + \frac{1}{2} \tilde{\omega}_{AB} \left( -\varepsilon^{\dot{\alpha}\dot{\beta}} D\tilde{q}_\alpha^A D\tilde{q}_\beta^B + i\varepsilon^{\alpha\beta} \tilde{\psi}_\alpha^A \not{D} \tilde{\psi}_\beta^B \right) - i\pi k_{mn} \varepsilon^{\alpha\beta} \varepsilon^{\dot{\gamma}\dot{\delta}} J_{\alpha\dot{\gamma}}^m J_{\beta\dot{\delta}}^n - i\pi k_{mn} \varepsilon^{\dot{\alpha}\dot{\beta}} \varepsilon^{\gamma\delta} \tilde{J}_{\dot{\alpha}\gamma}^m \tilde{J}_{\dot{\beta}\delta}^n + 4\pi i k_{mn} \varepsilon^{\alpha\gamma} \varepsilon^{\beta\delta} J_{\alpha\dot{\beta}}^m \tilde{J}_{\delta\dot{\gamma}}^n + i\pi k_{mn} \left( \varepsilon^{\dot{\alpha}\dot{\gamma}} \varepsilon^{\beta\delta} \tilde{\mu}_{\dot{\alpha}\beta}^m \psi_{\dot{\gamma}}^A t_{AB}^n \psi_{\delta}^B + \varepsilon^{\alpha\gamma} \varepsilon^{\beta\delta} \mu_{\alpha\beta}^m \tilde{\psi}_\gamma^A t_{AB}^n \tilde{\psi}_\delta^B \right) - \frac{\pi^2}{6} f_{mnp} (\mu^m)^\alpha_\beta (\mu^n)^\beta_\gamma (\mu^p)^\gamma_\alpha - \frac{\pi^2}{6} f_{mnp} (\tilde{\mu}^m)^\alpha_\beta (\tilde{\mu}^n)^\beta_\gamma (\tilde{\mu}^p)^\gamma_\alpha + \pi^2 (\tilde{\mu}^{mn})^\dot{\gamma}_\gamma (\mu_m)^\alpha_\beta (\mu_n)^\beta_\alpha + \pi^2 (\mu^{mn})^\gamma_\gamma (\tilde{\mu}_m)^\dot{\alpha}_\dot{\beta} (\tilde{\mu}_n)^\dot{\beta}_\dot{\alpha}, \quad (2.31)$$

# ABJM Model

★  $U(N)_k \times U(N)_{-k}$  Chern-Simons Theory N=6 Theory

★ Matter  $A_\mu, \tilde{A}_\mu, Z_\alpha, \Psi^\alpha$

★ U(1)-Global, SU(4) R-Symmetry,  $\xi_{\alpha\beta} = (\xi^{\alpha\beta})^* = \frac{1}{2}\epsilon_{\alpha\beta\gamma\delta}\xi^{\gamma\delta}$ ,

★ Lagrangian  $\mathcal{L}_{\text{Yukawa}} = \frac{2\pi i}{k} \text{Tr} \left( \bar{Z}^\alpha Z_\alpha \bar{\Psi}_\beta \Psi^\beta - Z_\alpha \bar{Z}^\alpha \Psi^\beta \bar{\Psi}_\beta + 2Z_\alpha \bar{Z}^\beta \Psi^\alpha \bar{\Psi}_\beta - 2\bar{Z}^\alpha Z_\beta \bar{\Psi}_\alpha \Psi^\beta \right. \\ \left. + \epsilon_{\alpha\beta\gamma\delta} \bar{Z}^\alpha \Psi^\beta \bar{Z}^\gamma \Psi^\delta - \epsilon^{\alpha\beta\gamma\delta} Z_\alpha \bar{\Psi}_\beta Z_\gamma \bar{\Psi}_\delta \right).$  (2.3)

$\mathcal{L}_{\text{potential}} = -U = -\frac{4\pi^2}{3k^2} \text{Tr} \left( 6Z_\alpha \bar{Z}^\alpha Z_\beta \bar{Z}^\beta Z_\gamma \bar{Z}^\gamma - 4Z_\alpha \bar{Z}^\beta Z_\gamma \bar{Z}^\alpha Z_\beta \bar{Z}^\gamma \right. \\ \left. - Z_\alpha \bar{Z}^\alpha Z_\beta \bar{Z}^\beta Z_\gamma \bar{Z}^\gamma - Z_\alpha \bar{Z}^\beta Z_\beta \bar{Z}^\gamma Z_\gamma \bar{Z}^\alpha \right).$  (2.4)

$\mathcal{L}_{\text{CS}} + \mathcal{L}_{\text{kin}} = \frac{k}{4\pi} \epsilon^{\mu\nu\rho} \text{Tr} \left( A_\mu \partial_\nu A_\rho - \frac{2i}{3} A_\mu A_\nu A_\rho - \tilde{A}_\mu \partial_\nu \tilde{A}_\rho + \frac{2i}{3} \tilde{A}_\mu \tilde{A}_\nu \tilde{A}_\rho \right) \\ - \text{Tr} \left( D_\mu \bar{Z}^\alpha D^\mu Z_\alpha + i \bar{\Psi}_\alpha \gamma^\mu D_\mu \Psi^\alpha \right),$

★ Vacuum Moduli Space  $C^4/Z_k, (C^4/Z_k)^N/S_N$

★  $k=1,2 \Rightarrow N=8$  Enhancement

# $U(1)_k \times U(1)_{-k}$ Case

★ Lagrangian

$$b_\mu = A_\mu - \tilde{A}_\mu, \quad c_\mu = \frac{1}{2}(A_\mu + \tilde{A}_\mu)$$

$$\mathcal{L} = \frac{k}{2\pi} \epsilon^{\mu\nu\rho} b_\mu \partial_\nu c_\rho + |(\partial_\mu - ib_\mu)Z_\alpha|^2$$

★ Auxiliary Field

$$\mathcal{L} = \frac{k}{2\pi} \epsilon^{\mu\nu\rho} (b_\mu + \frac{\sigma}{k}) f_{\nu\rho} + |(\partial_\mu - ib_\mu)Z_\alpha|^2$$

★ Monopole or flux change

$$\epsilon^{\mu\nu\rho} \partial_\mu f_{\nu\rho} = 2\pi \delta^3(x)$$

★ Periodic

$$\sigma \sim \sigma + 2\pi$$

★ Integration over f:

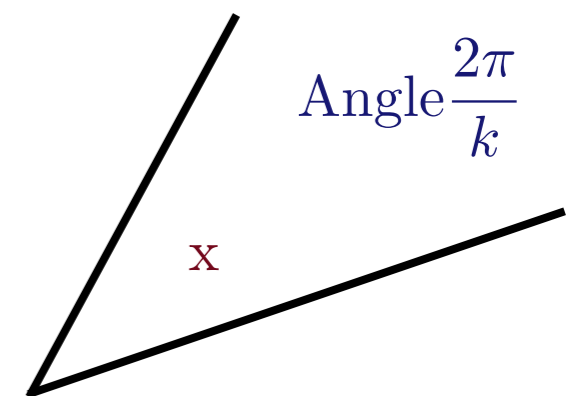
$$b_\mu = \frac{1}{k} \partial_\mu \sigma$$

★ Gauge Invariant

$$\hat{Z}_\alpha = e^{-\frac{i}{k}\sigma} Z_\alpha \quad \hat{Z}_\alpha \sim e^{2\pi i/k} \hat{Z}_\alpha$$

★ Moduli Space

$$C^4 / Z_k$$



# Chern-Simons Terms Everywhere

- ★ 5-dim Yang-Mills
- ★ 11-dim Supergravity action (Parity Even)
- ★ Mixed Chern-Simons Term
- ★ Noncommutative Plane and Fluid Mechanics
- ★ Anomaly and Axions in 4-dim
- ★ Quantum Hall Effect
- ★ More to appear